Structural, mechanical and osmotic properties of injectable hyaluronan-based composite hydrogels

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The osmotic and scattering properties of hyaluronan-based composite hydrogels composed of stiff biopolymer chains (carboxymethylated thiolated hyaluranon (CMHA-S)) crosslinked by a flexible polymer (polyethylene glycol diacrylate (PEGDA)) are investigated and analyzed in terms of the scaling theory. The total pre-gel polymer weight concentration is varied between 0.5 wt.% and 3.2 wt.%, while the mole ratio between the reactive PEG chain ends and the thiolated HA moieties is changed between 0.15 and 1.0. The shear modulus $G$ of the fully-swollen gels exhibits a stronger dependence on pre-gel concentration than on the crosslink density. Osmotic deswelling measurements reveal that the osmotic mixing pressure depends on the weight ratio CMHA-S/PEGDA, and is practically unaffected by the pre-gel concentration. Small-angle neutron scattering observations indicate that the thermodynamic properties of these composite gels are governed by total polymer concentration, i.e., specific interactions between the two polymeric components do not play a significant role.

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1. Introduction

Native hyaluronan (HA) is an ultrahigh molecular weight charged linear polysaccharide that is virtually invisible to the body’s immune system. It serves many important physiological roles, such as a lubricant within the synovial fluid and as an osmotic agent within the extracellular matrix [1]. Because of its unique biocompatibility and structure, HA and its derivatives are already used in a number of clinical applications, and many more have been proposed, including tissue engineering, wound healing, surgical adhesion protection, drug delivery, and soft tissue augmentation [2,3]. For many of these in vivo biomedical applications, it is useful to be able to crosslink native HA into gels. This can be accomplished by photoinitiation of a crosslinking reaction after injection of chemically-modified HA with a photoinitiator [4], or by premixing and injecting modified reactive HA molecules with a crosslinker designed to induce HA gelation after a suitable induction period (5–10 min). Here we focus on the latter approach, as developed by Shu and co-workers [5]. These authors found that linear HA chains of moderate molecular weight (≈200 kDa) can be used to formulate aqueous solutions that have relatively low viscosity as needed for injection and yet have a suitable induction period for gelation (≈10 min). The HA chains are made reactive by thiolating carboxylic acid groups in HA, and then forming crosslinks between thiolated carboxylic acid groups using polyethylene glycol (PEG) with reactive acrylate groups at each chain end. In addition to being useful for tissue engineering applications [5] and vocal cord augmentation [6], this system is also scientifically interesting because the gel is a complex composite containing approximately equal amounts of stiff charged carboxymethylated thiolated hyaluronan (CMHA-S) chains (persistence length $\approx$4–15 nm [7]) and flexible uncharged polyethylene glycol diacrylate (PEGDA) chains (persistence length $\approx$0.4 nm [8]). An important unanswered question is whether or not simple network models developed for flexible polymer gels [9] can describe the mechanical and swelling behavior of these complex systems. To address this question, we measure the elastic modulus $G$ and the osmotic modulus $K_T$ for nine different gel formulations obtained by varying the mole ratio between reactive crosslinker chain ends and thiols between 0.15 and 1.0, and varying the pre-gel polymer concentration between 0.5 wt.% and 3.2 wt.%. We also use SANS technique to determine the structure of these complex gels over a wide range of length scales. The values of $G$ and $K_T$ strongly affect the range of potential tissue...
engineering applications. For example, it is now well documented that a scaffold designed to support tissue growth of a certain type should have an elastic modulus value that approximately matches that of the native tissue [10]. This may range from 10^2 Pa for nerve tissue to 10^9 Pa for connective tissue. The osmotic modulus, which is a measure of the pressure required to squeeze water out of a gel, is particularly important for HA-containing hydrogels designed to replace tissues that cushion mechanical forces within the body. As far as we know, the osmotic modulus value has not been reported before for injectable CMHA-S based gels. In addition, there have been no previous reports of the influence of the pre-gel CMHA-S concentration, which we find has a stronger effect on the elastic shear modulus than does the crosslinker mole ratio.

2. Theory

Unconfined neutral hydrogels swell until the total change in free energy, \( \Delta G_{\text{tot}} \), reaches a minimum or, equivalently, until the chemical potential of each mobile species becomes equal in the coexisting phases. The mixing of water with the polymer chains makes a negative contribution to the free energy \( \Delta G_{\text{mix}} \), while the stretching of the hydrogel network makes a positive contribution \( \Delta G_{\text{el}} \). Assuming these terms are independent, we can write [11–13]:

\[
\Delta G_{\text{tot}} = \Delta G_{\text{mix}} + \Delta G_{\text{el}}
\]

In the case of charged polymers such as HA the ions also contribute to \( \Delta G_{\text{tot}} \). However, several experimental studies on highly swollen polyelectrolyte hydrogels (e.g., polyacrylic acid gels, DNA gels) [13,14] as well as molecular dynamics simulations [15] indicate that in the presence of large amount of added salt (e.g., in buffer) their osmotic behavior can be reasonable approximated by the simple expression derived for neutral polymer gels.

In an osmotic swelling experiment the measurable quantities involve derivatives of the free energy, i.e.

\[
\Pi_{\text{tot}} = \frac{1}{V_1} \left( \frac{\partial \Delta G_{\text{tot}}}{\partial n_1} \right) = \Pi_{\text{mix}} + \Pi_{\text{el}}
\]

where \( \Pi_{\text{tot}} \) is the swelling pressure of the gel, \( \Pi_{\text{mix}} \) and \( \Pi_{\text{el}} \) are the mixing and elastic contributions of \( \Pi_{\text{tot}} \), respectively. \( V_1 \) is the molar volume of solvent (water), and \( n_1 \) is the number of moles of water. At swelling equilibrium with the pure solvent \( \Pi_{\text{el}} \) must equal zero. One can study nonzero values of the total swelling pressure by equilibrating the hydrogel with an osmotic stressing agent of known osmotic pressure [13,14].

The osmotic modulus \( K_T \) is defined by:

\[
K_T = C \left( \frac{\partial \Pi_{\text{tot}}}{\partial c} \right)_T = \left( \frac{\partial \Pi_{\text{tot}}}{\partial n_1} \right)_T
\]

where \( C \) is the weight fraction of polymer in the hydrogel.

3. Experimental methods

3.1. Gel preparation

HA biopolymer with a backbone molecular weight of approximately 200 kDa and modified with thiol groups to make it crosslinkable was prepared by Glycosan Biosystems, Inc., Salt Lake City, Utah following a similar reaction procedure as in Ref. [5]. In brief, the repeat group of HA is a disaccharide containing D-glucuronic acid and N-acetyl-D-glucosamine, with one pendant carboxylic acid group in D-glucuronic acid. In order to increase the number of sites available for crosslinking, an additional carboxylic acid group was attached to each glucosamine unit. The primary hydroxyl group in glucosamine was deprotonated by pH adjustment (pH > pKa), and then reacted with chloroacetic acid to obtain the additional carboxylic acid group. Both carboxylic acid groups in the modified HA repeat unit were then available for thiolation via reaction with 3-thiopropanoyl hydrizide, with the degree of thiolation measured using a modified Ellman method [5]. The resulting HA derivative, denoted as CMHA-S, has the repeat unit shown in Fig. 1(a). The crosslinker molecule, polyethylene glycol diacylate (PEGDA) consists of a linear PEG backbone having a molecular weight of 3500 Da with one acrylate group attached to each chain end [5,16]. PEGDA rapidly crosslinks CHMA-S (within 10 min) following the reaction scheme shown in Fig. 1(b). Even in the absence of PEGDA, CMHA-S crosslinking will occur in the presence of oxygen due to disulfide linkage formation (Fig. 1(c), [16,17]). However, the rate of the disulfide crosslinking reaction is slow compared to the PEGDA crosslinking reaction [16].

Cylindrical hydrogel specimens were cast at room temperature in a Teflon™ mold of thickness 1 cm and diameter 1 cm. Aqueous solutions containing the PEGDA and CMHA-S were mixed in an adjustable-volume pipette for 30 s to give the pre-gel polymer concentrations listed in Table 1 and then rapidly injected into the mold. For each pre-gel CMHA-S concentration value, the weight ratio PEGDA/CMHA-S was varied between 0.25 and 1.66 giving a reactive functional group mole ratio (acrylate to thiolated carboxylic acid) that varied between 0.15 and 1.0. Since the thiol–acylate reaction rate is strongly pH sensitive, the crosslinking reaction was carried out at pH = 7.4. To ensure complete crosslinking the samples were kept in closed containers for two weeks before measurements.

3.2. Osmotic swelling measurements

The reference environmental solution (infinite bath) was chosen to be phosphate-buffered saline (PBS) solution at 25 °C with ionic strength 0.15 M and pH = 7.4. Hydrogels of known dry mass \( m_d \) were equilibrated with this reference solution, at which point gel swelling pressure \( \Pi_{\text{tot}} \) must equal zero. Deswelling of the gels was achieved by enclosing them in a semipermeable membrane (dialysis bag, seamless cellulose tubing; cut off molecular weight: 12 kDa, Sigma Chemical Co., St. Louis, MO). Known concentrations

\[
\begin{array}{c}
\text{Fig. 1. Crosslinking reaction scheme showing (a) repeat unit of CMHA-S, (b) rapid crosslinking of CMHA-S by PEGDA, and (c) slow disulfide crosslinking of CMHA-S that occurs in the presence of oxygen.}
\end{array}
\]
of an osmotic deswelling agent [poly(vinyl pyrrolidone)], PVP, $M_n = 29$ kDa, were added to the environmental solution, and equilibrated with the gel for 8–10 days. The semipermeable membrane prevented penetration of PVP into the gel. At equilibrium, the swelling pressure $P_{\text{tot}}$ of the gel inside the dialysis bag is equal to the known osmotic pressure of the PVP solution outside \([18]\). Periodically the mass of the gel $m_g$ was measured, and used to calculate the total polymer weight fraction $C = m_g/m_0$, with hydrogel swelling ratio $Q = 1/C$. Equilibrium was achieved when no further changes in either gel swelling degree or solution composition were detectable. The gel samples were dried at 95 °C. First, the mass of the salt-containing gel was measured. Then the gel was soaked in large excess of distilled water to remove the salt. The swelling ratio was calculated by correcting the measured mass with the known amount of salt.

Reversibility was checked by transferring gels into PVP solutions at different osmotic pressure values.

### 3.3. Elastic (shear) modulus measurements

Shear modulus data of fully-gelled materials were obtained from uniaxial compression measurements using a TA.XT21 HR Texture Analyser (Stable Micro Systems, UK). As discussed in a standard rubber testing handbook \([19]\), torsional measurements of the shear modulus $G$ of fully-gelled polymers will be erroneous if sample "slip" occurs at the interface between the solid polymer and the steel rheometer plates, because then the true value of the shear strain exerted on the sample will be unknown. Hence the standard rubber testing handbook recommends lubricating the sample interface to ensure that sample slip does occur, and then determining $G$ in the absence of "barrel" distortion by measuring the force needed to apply compressive strains (uniaxial) on the sample \([19]\).

Equilibrated gel samples were rapidly transferred from the dialysis bag into the TA XT21 HR apparatus, which measures the compressive deformation (precision: $\pm 0.001$ mm) as a function of an applied force (precision: $\pm 0.01$ N). Stress-strain isotherms were determined at 25 °C in about 3–5 min with no detectable change in gel weight. The absence of sample volume change and barrel distortion was confirmed, and stress-strain isotherms were found to fit the Mooney–Rivlin relation \([12,20,21]\)

$$\sigma = C_1 (\Lambda - \Lambda^{-2}) + C_2 (1 - \Lambda^{-3}) \quad (4)$$

where $\sigma$ is the nominal stress (related to the undeformed cross-section of the gel), $\Lambda$ is the deformation ratio ($\Lambda = L/L_0$, $L$ and $L_0$ are the lengths of the deformed and undeformed specimen, respectively), and $C_1$ an $C_2$ are constants. The stress-strain data were determined in the range of deformation ratio $0.7 < \Lambda < 1.0$. The value of $C_2$ proved to be negligibly small for the gel systems studied.

In this situation, the constant $C_1$ can be identified with the shear modulus ($G$) of the swollen network.

The swelling and mechanical measurements were carried out at $25 \pm 0.1$ °C. Repeated measurements showed a mean change in the osmotic swelling pressure and elastic modulus less than 2–3%.

### 3.4. Small-angle neutron scattering

SANS measurements were made on gels on the NG3 30 m instrument \([13]\) at the National Institute of Standards and Technology (NIST, Gaithersburg MD). Gel samples were prepared in solutions of heavy water in 2 mm thick sample cells. The sample cell consisted of 1 mm thick quartz windows separated by a 2 mm thick spacer. The beam diameter was 20 mm. The measurements were made at three sample-detector distances, 1.3 m, 4 m and 13.1 m, with incident wavelength 8 Å. The explored wave vector range was $0.003 \text{ Å}^{-1} \leq q \leq 0.2 \text{ Å}^{-1}$, and counting times from 20 min to two hours were used. The total counts of neutron at the detector varied in the range 0.5–3 million. After radial averaging, corrections for detector response and cell window scattering were applied. The neutron scattering intensities were calibrated using NIST absolute intensity standards. The incoherent background was subtracted following the procedure described in reference \([22]\). All experiments were carried out at $25 \pm 0.1$ °C.

### 4. Results and discussion

#### 4.1. Mechanical measurements and osmotic pressure observations

In Fig. 2, the shear modulus $G$ is plotted against the total weight fraction of polymer in the fully-swollen hydrogel for each of the nine compositions synthesized (Table 1). Each hydrogel has been equilibrated against the reference environmental solution (PBS buffer), hence the swelling pressure $P_{\text{tot}}$ is zero. As shown in the insets to the figure, the value of $G$ increases with increase in mole ratio of crosslinker (at fixed CMHA-S pre-gel concentration), and with increase in CMHA-S pre-gel concentration (at fixed crosslinker mole ratio). In inset A, the data lies below a line of slope one, indicating that some of the reactions between acrylate and thiolated HA moieties are either incomplete or do not produce cross-links that contribute to the elastic modulus, i.e., the efficiency of the crosslinking process decreases with increasing crosslinker mole ratio. In inset B, the data lies above the continuous lines that have a slope of two, indicating that the dependence of $G$ on CMHA-S pre-gel concentration is stronger than quadratic. The origin of the strong dependence of $G$ on pre-gel CMHA-S concentration is most likely trapped interchain physical entanglements. Physical entanglements are primarily detected via rheological measurements. For example, solutions of high molecular weight polymer chains exhibit gel-like behavior even in the absence of crosslinking reactions, which is indirect evidence for the presence of physical entanglements. At long times, these physical entanglements are eliminated in un-crosslinked systems by diffusion of the polymer chains along their contour lengths. However, when covalent crosslinking occurs in a polymer solution that is already physically entangled, then the physical entanglements become trapped and make a permanent contribution to $G$. Experimental evidence for the importance of pre-gel polymer concentration and trapped interchain physical entanglements on the $G$ value of acrylamide-based hydrogels has been reported by Baker et al. \([23]\), and for thiolated chitosan gels crosslinked by PEGDA by Teng et al. \([17]\). The intermediate time plateau shear modulus of a physically-entangled gel is usually observed to scale as approximately the second power of the polymer concentration \([24]\). The theoretical lines shown in inset B have slope of 2.0, which is the value based on the simple

<table>
<thead>
<tr>
<th>Composition of hyaluronan hydrogels studied.</th>
<th>$C_1$ (wt.%)</th>
<th>$C_2$ (wt.%)</th>
<th>$Q^1$</th>
<th>$Q^2$ (Pa)</th>
</tr>
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</tr>
</tbody>
</table>

$^a$ Measured at zero swelling pressure.
idea that physical entanglements are random point contacts between two chains. This simple theory appears to underestimate the observed concentration dependence. According to de Gennes [25,26] the concentration dependence of the plateau modulus for neutral polymers in good solvent conditions is given by $G \sim C^{2/3}$, while in a theta solvent $G \sim C^{1/3}$ as discussed by Colby and Rubinstein [27,28]. The more advanced scaling theories do a better job of representing the experimental data in inset B.

In Fig. 2, the $G$ values given were measured for each hydrogel in equilibrium with PBS buffer, i.e., $\Pi_{\text{tot}} = 0$. One observes that under these conditions, all of the fully-swollen hydrogels contain at least 95 wt.% water. However, by placing a hydrogel on one side of a dialysis membrane, and by placing an osmotic stressing agent on the other side of the membrane, one can force the hydrogel to adopt a smaller swelling ratio at which the swelling pressure $\Pi_{\text{tot}}$ is greater than zero.

The results of the osmotic deswelling experiments are shown in Fig. 3 where the swelling pressure $\Pi_{\text{tot}}$ is plotted against the total polymer weight fraction $C$ for each of the nine different samples. For each hydrogel, it takes at least 60 kPa of pressure to reduce the water content to 85 wt.%. It is also apparent that the behavior of gels prepared with the same crosslinker ratio is practically independent of the initial polymer concentration.

According to Equation (2), $\Pi_{\text{tot}}$ is the sum of a mixing contribution $\Pi_{\text{mix}}$ and an elastic contribution $\Pi_{\text{el}}$. The elastic contribution is usually taken to be the negative of the shear modulus $G$ [11], which was independently measured as a function of $C$ (see Experimental). Fig. 4 shows the variation of the shear modulus as a function of the total polymer concentration for the hydrogels.

It can be seen that the concentration dependence of $G$ is satisfactorily described by the scaling law

$$G = BC^n$$

where $B$ is a constant depending on both the CMHA-S/PEGDA weight ratio (crosslink density) and the CMHA-S concentration in the pre-gel solution. The best-fit values of the pre-factor $B$ and the exponent $n$ are listed in Table 2. The exponent $n$ is close to one-third as predicted by the Flory–Rehner theory [11]. The experimental values of $n$ are consistent with previous observations made on different synthetic and biopolymer gels [13,29–31].

The mixing contribution to the osmotic swelling pressure $\Pi_{\text{mix}}$ was obtained by adding the measured $G$ and $\Pi_{\text{tot}}$ values (see Fig. 3).
Equation (2)). In Fig. 5, the mixing pressure $P_{\text{mix}}$ is plotted against the total weight fraction of the polymer for all samples. In semi-dilute polymer solutions, the mixing pressure is expected to obey a scaling relation [9]

$$P_{\text{mix}} = A C^m$$

(6)

where $A$ and $m$ are constants. For flexible chains $m \approx 2.3$ (good solvent condition) and $m = 3$ (theta condition). The best-fit values of the pre-factor $A$ and the exponent $m$ are given in Table 3.

The data shown in Fig. 5 and Table 3 illustrate that

(i) the mixing pressure is primarily governed by the polymer/crosslinker (CMHA-S/PEGDA) ratio, and

(ii) at constant polymer/crosslinker ratio $P_{\text{mix}}$ is independent of the polymer concentration in the pre-gel solution.

Equation (3) is the defining equation of the osmotic modulus $K_T$, which is a measure of the resistance of the gel to deswelling under pressure. This quantity can be calculated using Equation (3) in conjunction with the best-fit values of the parameters appearing in Equations (5) and (6).

Combining Equations (5) and (6) we have

$$P_{\text{tot}} = P_{\text{mix}} + P_{\text{el}} = P_{\text{mix}} - G = A C^m - B C^n$$

(7)

At equilibrium with the pure solvent (or salt solution) $P_{\text{tot}} = 0$ and $C = C_e$ (where $C_e$ is the polymer concentration of the fully-swollen gel). Thus, $B = A C^m e^{-n}$ and we get

$$P_{\text{tot}} = A C^m - A C^m C^n$$

(8)

From Equations (7) and (8) the osmotic modulus is given as

$$K_T = C(\partial P_{\text{tot}}/\partial C)_T = m A C^m - n A C^m C^n$$

(9)

The osmotic modulus values obtained from Equation (9) are plotted against total polymer weight fraction in Fig. 6.

Comparing Figs. 4 and 6, one observes that $K_T$ increases much faster with increasing polymer concentration than $G$.

In what follows we estimate the ratio of the osmotic modulus to the shear modulus from simple scaling considerations. In fully-swollen gels ($P_{\text{tot}} = 0$) at $C = C_e$ we have

$$K_T^e = C(\partial P_{\text{tot}}/\partial C)_T = A C^m (m - n)$$

(10)

Since the shear modulus of fully-swollen gels $G_e$ obeys the scaling relation [9]

$$G_e = A C^m$$

(11)

from Equations (10) and (11) we get

Table 2
Scaling parameters for $G$.

<table>
<thead>
<tr>
<th>Hydrogel designation</th>
<th>$B$ (kPa)</th>
<th>$n$</th>
</tr>
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<tbody>
<tr>
<td>1/1</td>
<td>0.40</td>
<td>0.32</td>
</tr>
<tr>
<td>1/2</td>
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<td>0.34</td>
</tr>
<tr>
<td>2/1</td>
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<td>0.36</td>
</tr>
<tr>
<td>2/2</td>
<td>6.93</td>
<td>0.33</td>
</tr>
<tr>
<td>2/3</td>
<td>10.71</td>
<td>0.34</td>
</tr>
<tr>
<td>3/1</td>
<td>4.84</td>
<td>0.29</td>
</tr>
<tr>
<td>3/2</td>
<td>11.18</td>
<td>0.33</td>
</tr>
<tr>
<td>3/3</td>
<td>17.15</td>
<td>0.32</td>
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</table>

Table 3
Scaling parameters for $P_{\text{mix}}$.

<table>
<thead>
<tr>
<th>Hydrogel designation</th>
<th>$A$ (kPa)</th>
<th>$m$</th>
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</thead>
<tbody>
<tr>
<td>1/1</td>
<td>3830</td>
<td>1.95</td>
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<tr>
<td>1/2</td>
<td>3553</td>
<td>1.99</td>
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</tr>
<tr>
<td>3/3</td>
<td>3252</td>
<td>2.06</td>
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Fig. 5. Osmotic mixing pressure vs. total weight fraction of polymer in the hydrogel, for each of the nine hydrogels synthesized. The designation of each hydrogel, as defined in Table 1, is given in the legend. The dashed curves are drawn as a guide to the eyes and connect gels of the same crosslinker mole ratio. The experimental error of these measurements is comparable to the size of the symbols.

Fig. 6. Osmotic modulus vs. total weight fraction of polymer in the hydrogel, for each of the nine hydrogels synthesized. The designation of each hydrogel, as defined in Table 1, is given in the legend. The dashed curves are drawn as a guide to the eyes and connect gels of the same crosslinker mole ratio. The experimental error of these measurements is comparable to the size of the symbols.
where $K_f$ is the value of the osmotic modulus at $C = C_e$.

In Fig. 7 is plotted $K_f/G_0$ as a function of $C_e$ for each of the nine hydrogels. The data points are close to 1.7 (continuous line), which is consistent with Equation (12) using the values of $m$ and $n$ listed in Tables 2 and 3. The power law exponent for the dependence of $\Pi_{\text{aux}}$ on $C$ ($m \approx 2$, see Table 3) is much greater than the power law exponent for the dependence of $G$ on $C$ ($n \approx 0.3$, see Table 2), which explains why $K_f$ is much greater than $G$ in the deswollen gels.

4.2. Small-angle neutron scattering measurements

To obtain information on the spatial organization of the polymer molecules in these composite gels at the nanoscale we made SANS measurements. Fig. 8 shows the SANS spectra for a set of CMHA-S/PEGDA gels prepared at different pre-gel concentrations with constant CMHA-S/PEGDA ratio. The SANS measurements were made on fully-swollen gels. As expected the scattered intensity increases with increasing polymer concentration, it is the greatest for the gel prepared at the highest total polymer content. All the spectra exhibit the same characteristic features: at low $q$ the intensity $I(q)$ decreases and displays a power law scattering regime followed by a shoulder at higher values of $q$. The upturn in $I(q)$ at low $q$ ($q < 0.01 \text{ Å}^{-1}$) indicates cluster formation generally observed in polyelectrolyte solutions [32–34]. The size of these objects is greater than 1000 Å, i.e., the resolution of the SANS measurements. The shoulder becomes more pronounced as the polymer concentration increases. In the high $q$-region the scattering intensity is governed by the local geometry of the polymer molecules.

The SANS spectra were analyzed using Equation (13), which reproduces the main characteristic features of the scattering curves

\[
I(q) = \frac{A_1}{1 + a q^2 + \frac{A_2}{q^2}},
\]

where $\xi$ is the polymer–polymer correlation length, and $A_1, A_2$ and $s$ are constants. The intensities $A_1$ and $A_2$ are proportional to the average scattering contrast between the polymer and the solvent. In Equation (13) the first term (Lorentzian form factor) is primarily governed by the thermodynamic concentration fluctuations, while the second term arises from large-scale static inhomogeneities frozen-in by the crosslinks. These large objects are not expected to make significant contribution to the thermodynamic properties of the system [26,35,36].

The parameters obtained from the least-squares fit of Equation (13) to the SANS data (continuous lines in Fig. 8) are listed in Table 4. The inset in Fig. 8 shows that the correlation length $\xi$ decreases with increasing polymer concentration as

\[
\xi \propto C^{-0.65}
\]

Scaling theory [9] predicts for the osmotic pressure of semi-dilute polymer solutions

\[
\Pi \propto \frac{kT}{\xi^2}
\]

where $k$ is the Boltzmann constant and $T$ is the absolute temperature. Combination of (14) and (15) yields the concentration dependence of the osmotic pressure $\Pi \propto C^{1.95}$, which is in reasonable agreement with the results of the macroscopic osmotic observations (Table 3).

In the low $q$-region the slope $s$ is practically the same for all the three gels; its value ($\approx 2.4$) is characteristic of randomly branched structures.

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**Table 4** Parameters from fits of Equation (13) to the SANS spectra of PEGDA/CMHA-S gels.

<table>
<thead>
<tr>
<th>Sample</th>
<th>$A_1$ (cm$^{-1}$)</th>
<th>$A_2$ (cm$^{-1}$)</th>
<th>$\xi$ (Å)</th>
<th>$s$</th>
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<td>0.009 ± 0.002</td>
<td>0.080 ± 0.01</td>
<td>38 ± 4</td>
<td>2.48 ± 0.1</td>
</tr>
<tr>
<td>2/3</td>
<td>0.011 ± 0.002</td>
<td>0.155 ± 0.02</td>
<td>21 ± 2</td>
<td>2.44 ± 0.1</td>
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<tr>
<td>3/3</td>
<td>0.014 ± 0.002</td>
<td>0.164 ± 0.02</td>
<td>15 ± 2</td>
<td>2.42 ± 0.1</td>
</tr>
<tr>
<td>3/2</td>
<td>0.012 ± 0.002</td>
<td>0.144 ± 0.01</td>
<td>16 ± 2</td>
<td>2.29 ± 0.1</td>
</tr>
<tr>
<td>3/1</td>
<td>0.010 ± 0.002</td>
<td>0.111 ± 0.01</td>
<td>15 ± 2</td>
<td>2.17 ± 0.1</td>
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</tbody>
</table>

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Fig. 7. Dependence of the ratio $K_f/G_0$ on the equilibrium polymer concentration $C_e$ for each of the nine hydrogels. Inset shows the variation of the osmotic modulus $K_f$ and the shear modulus $G_0$ as a function of $C_e$.

Fig. 8. SANS spectra for three hydrogels in the fully-swollen state with the same crosslinker mole ratio. The designation of each hydrogel, as defined in Table 1, is given in the figure. Inset shows the variation of the correlation length $\xi$ and the exponent $s$ with polymer weight fraction.
To estimate the effect of the chemical composition on the structure of CMHA-S/PEGDA hydrogels in Fig. 9 are shown the SANS spectra of three gels prepared with increasing amount of crosslinker (PEGDA).

The SANS measurements were made at the same total polymer concentration [0.044 (w/w)]. The increase in the scattering intensity with increasing PEGDA content reflects the change in the scattering contrast. However, the shape of the SANS curves hardly varies with the chemical composition. This result indicates that in the present gel systems specific interactions between the PEGDA and CMHA-S components do not play a significant role. Use of Equation (13) to describe the scattering curves yields a satisfactory fit to the data at each PEGDA/CMHA-S ratio (continuous lines through the SANS data points). The fitting parameters are listed in Table 4. The correlation length $\xi$ is practically unaffected by the chemical composition of the network implying that the thermodynamic properties of these composite gels are governed by total polymer concentration. The slope of the $I(q)$ plots in the power law regime weakly increases with increasing PEGDA/CMHA-S ratio (see inset). This result suggests that crosslinking slightly affects the static structure of the network.

5. Conclusions

Macro-scale (mechanical and osmotic) and micro-scale (SANS) measurements have been performed on nine different injectable CMHA-S hydrogels that were synthesized by varying crosslinker mole ratio and pre-gel CMHA-S concentration. These gels are potentially useful for biomedical applications (e.g., tissue engineering, soft tissue augmentation) and also scientifically interesting because the gels are complex composites containing roughly equal amounts of stiff charged CMHA-S chains and flexible uncharged PEGDA chains. Nonetheless, scaling theory developed for simple uncharged gels has been shown to describe the mechanical, osmotic, and SANS results at physiological ionic strength (0.15 M). SANS measurements reveal that the nanoscale structure of these composite gels is primarily governed by the total polymer content and only weakly affected by CMHA-S/PEGDA ratio. In addition, scaling theory correctly predicts that the osmotic modulus equals about twice the shear modulus $G$ for fully-swollen gels, and exhibits much faster increase with increase in polymer concentration that occurs when water is expelled from a given gel due to the presence of an external osmotic stressing agent.

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