Incompressible Flow Over Finite wings

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Introduction



Beechcraft Baron-58 aircraft
Airfoil - NACA 23015

- $\bigstar Angle of attack = 4 deg.$
- ***** $c_l = 0.54$ and $c_d = 0.0068$



Are the lift and drag coefficients for the Beechcraft wing same as the airfoil?
NO!!

Downwash and Induced Drag

 Why are the aerodynamic characteristics of finite wings different from its airfoil sections?
 Flow over an airfoil is two-dimensional.

In contrast, flow over a finite wing is three-dimensional.





Downwash and Induced Drag







Trailing vortices at each wing tip.

- The two vortices drag the surrounding air inducing a velocity component in the downward direction downwash.
- The downwash combines with the local freestream to create a local relative wind.

Downwash and Induced Drag



The downwash has two important effects:

***** The effective angle of attack is reduced.

***** Induced drag is created due to the tilting of the local lift vector.

* The total drag = friction drag + pressure drag + induced drag.

The Vortex Filament Theorem

Establish a rational aerodynamic theory for a finite wing.



The curved filament induces a flow field in the surrounding space.
Circulation taken about any closed path enclosing the filament is constant.

Consider a segment *dl*. It induces a velocity at *P* equal to:

$$dV = \frac{\Gamma}{4\pi} \frac{dl \times r}{|r|^3}$$

Biot-Savart Law

The Vortex Filament Theorem

When a number of vortex filaments are used in conjunction with a uniform free stream, it is possible to synthesize the flow over a finite wing.



Biot-Savart Law to a straight vortex filament of infinite length

Velocity induced at P by the entire vortex filament is given by:

$$V = \int_{-\infty}^{\infty} \frac{\Gamma}{4\pi} \frac{dl \times r}{|r|^3} = \frac{\Gamma}{2\pi h}$$

The above result is the same as that for a point vortex in a 2D flow.

The Vortex Filament Theorem



Velocity induced at P by the vortex filament is given by: $V = \int_{A}^{\infty} \frac{\Gamma}{4\pi} \frac{dl \times r}{|r|^{3}} = \frac{\Gamma}{4\pi h}$

- Helmholtz's Vortex Theorems (basic principles of vortex behavior)
 The strength of a vortex filament is constant along its length.
 - A vortex filament cannot end in a fluid; it must extent to the boundaries of the fluid or form a closed path.

Lift Distribution on a Finite Wing





Consider a given spanwise location y_1 , where the local chord is c.

The lift per unit span can vary along the span.

Different spanwise locations can have different angles of attack (geometric twist).



Lift Distribution on a Finite Wing



Consider a given spanwise location y_1 , where the local chord is c.

- Wings can also have different airfoil section spanwise (aerodynamic twist).
- Pressure equalization occurs at y = -b/2and b/2, and consequently there is no lift at these locations.
- Our objective is to estimate this lift distribution, total lift and induced drag for the finite wing.

Solution The theory is useful for predicting the aerodynamic characteristics of finite wings.



The finite wing is replaced with a bound vortex.

Due to Helmholtz's theorem, a vortex filament cannot end in the fluid.

* Therefore, assume the vortex filament continues as two free vortices trailing downstream from the wing tips to infinity.

Bound vortex + Trailing vortices —> Horseshoe Vortex.

Consider the downwash induced along the bound vortex by the horseshoe vortex.



The velocity at any point along the bound vortex induced by the trailing vortex is: $w(y) = -\frac{\Gamma}{4\pi(b/2+y)} - \frac{\Gamma}{4\pi(b/2-y)}$ left vortex right vortex $w(y) = -\frac{\Gamma b}{4\pi[(b/2)^2 - y^2]}$

The downwash approaching infinity value at the tips is disconcerting.

Instead of representing the wing by a single horseshoe vortex, superimpose using a large number of horseshoe vortices.
 Each horseshoe with a different length of the bound vortex.

* All bound vortices coincident along a single line - Lifting Line.



The series of trailing vortices represents pairs of vortices.
 * Each pair is associated with a given horseshoe vortex.

The strength of each trailing vortex is equal to the change in circulation along the lifting line.

Let us extrapolate to the case where an infinite number of horseshoe vortices are superimposed along the lifting line.
Each horseshoe has vanishingly small strength.



The finite number of trailing vortices in the earlier case have become a continuous vortex sheet.

The total strength of the sheet integrated across the span of the wing is zero (because of pairs of trailing vortices of equal but opposite strengths).

Solution O(x) = C(x) + C(x) and C(x) = C(x) + C(x) + C(x). Substituting the segment of the lifting line dy.

***** The circulation, and the change in circulation at y are respectively:

 $\Gamma(y)$ and $d\Gamma = (d\Gamma/dy)dy$



In turn, the strength of the trailing vortex at y must equal the change in circulation along the lifting line.



The velocity dw induced at y_0 by the entire semi-infinite trailing vortex located at y is:

$$dw = -\frac{(d\Gamma/dy)dy}{4\pi(y_0 - y)}$$

The total velocity induced at y₀ by the entire trailing vortex sheet is:

$$w(y_0) = -\frac{1}{4\pi} \int_{-b/2}^{b/2} \frac{(d\Gamma/dy)dy}{(y_0 - y)}$$

However, our central problem remains to be solved: we want to calculate $\Gamma(y)$ for a given finite wing.



$$\alpha_i(y_0) = \tan^{-1}\left(\frac{-w(y_0)}{V_{\infty}}\right)$$

For small angles, $\alpha_i(y_0) = \left(\frac{-w(y_0)}{V_{\infty}}\right)$

The geometric angle of attack is given by:

$$\alpha(y_0) = \frac{\Gamma(y_0)}{\pi V_{\infty} c(y_0)} + \alpha_{L=0}(y_0) + \frac{1}{4\pi V_{\infty}} \int_{-b/2}^{b/2} \frac{(d\Gamma/dy)dy}{y_0 - y}$$

Prandtl's Lifting Line Theory Equation

Consider a circulation distribution given by:

$$\Gamma(y) = \Gamma_0 \sqrt{1 - \left(\frac{2y}{b}\right)^2}$$

The circulation, and hence lift, goes to zero at the wing tips.

* What are the aerodynamic properties of a finite wing with such an elliptic lift distribution?

 $w(\theta_0) = -\frac{\Gamma_0}{2h}$

Constant downwash over the span

Constant induced angle of attack over the span $\alpha_i = -\frac{\Gamma_0}{2bV_{\infty}} = \frac{C_L}{\pi(AR)}$

Induced drag proportional to square $C_{D,i} = \frac{C_L^2}{\pi(AR)}$ of lift coefficient



- The dependence of induced drag on lift is not surprising.
 Induced drag is a consequence of the presence of wing-tip vortices produced by the pressure difference between lower and upper wing surfaces.
 - Lift is also produced by this same pressure difference and hence the dependence.
- Induced drag is the price for generation of lift. An aircraft cannot generate lift for free.
 - Power required for an aircraft to overcome induced drag is the power required to generate the lift of the aircraft.
 - Induced drag is high at take-off, landing, and about 25% of total drag at cruising speeds.



- Induced drag is inversely proportional to the aspect ratio.
 To reduce induced drag, we want a finite wing with the highest possible aspect ratio.
 - Design of high aspect ratio wings with sufficient structural strength is difficult.
 - Therefore aspect ratio is a compromise between aerodynamic and structural requirements.

 AR = b²/S



- Another property of an elliptic lift distribution:
 Consider a wing with no geometric twist (angle of attack is constant along the span).
 - Also, assume no aerodynamic twist (zero-lift angle of attack is constant along the span).

The local section lift coefficient is $c_l = a_0(\alpha_{eff} - \alpha_{L=0})$

The lift per unit span is then given by $L'(y) = q_{\infty}cc_l$

Solving for the chord, we have $c(y) = \frac{L'(y)}{q_{\infty}c_l}$

- The above equations show that for an elliptic lift distribution:
 The chord must vary elliptically along the span.
 - ***** The wing planform is elliptical.

Prandtl's Lifting Line Theory - General Lift Dist.

Consider the transformation: $y = -\frac{b}{2}cos\theta$ The elliptic lift distribution is then given by:
 $\Gamma(\theta) = \Gamma_0 sin\theta = 2bV_\infty \sum_{i=1}^{N} A_n sin(n\theta)$

So the angle of attack evaluated at a given spanwise location is: $\alpha(\theta_0) = \frac{2b}{\pi c(\theta_0)} \sum_{1}^{N} A_n \sin(n\theta_0) + \alpha_{L=0}(\theta_0) + \sum_{1}^{N} nA_n \frac{\sin(n\theta_0)}{\sin\theta_0}$

The only unknowns in the above equation are the Fourier coefficients.

* They can be evaluated by considering 'n' spanwise locations.

Prandtl's Lifting Line Theory - General Lift Dist.

Solution Now that $\Gamma(\theta)$ is known, the lift coefficient can be calculated: $C_L = A_1 \pi(AR)$

Solution The induced drag coefficient is then calculated to be:

$$C_{D,i} = \frac{C_L^2}{\pi(AR)} (1+\delta), \text{ where}$$
$$\delta = \sum_{2}^{N} n(A_n/A_1)^2$$

Prandtl's Lifting Line Theory - General Lift Dist.



AR typically varies from 6 to 22 for standard subsonic aircraft.
 The primary design factor for minimizing induced drag is to make the aspect ratio as large as possible.

The total drag is: profile drag + induced drag

$$C_D = c_d + \frac{C_L^2}{\pi e A R}$$

Consider two wings with different aspect ratios but same airfoil cross-section.

* It is possible to scale the data of a wing with one aspect ratio to correspond to the case of another aspect ratio.

$$C_{D,1} = C_{D,2} + \frac{C_L^2}{\pi e} \left[\frac{1}{5} - \frac{1}{AR_2} \right]$$

Effect of Aspect ratio

- There are two primary differences between airfoil and finite wing properties.
 - * One: finite wing generates induced drag.

***** *Two*: *The lift slope of a finite wing is less than that of an infinite wing.*



$$a = \frac{a_0}{1 + (a_0/\pi AR)(1+\tau)}$$

tis a function of Fourier coefficients and typically range between 0.05 and 0.25.

