
Incompressible Flow Over Airfoils

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● *We will focus on how to obtain **airfoil properties**.*

- ✱ Circulation theory

- ✱ Source Panel Method

- ✱ Design and Performance

● *Aerodynamic consideration of wings:*

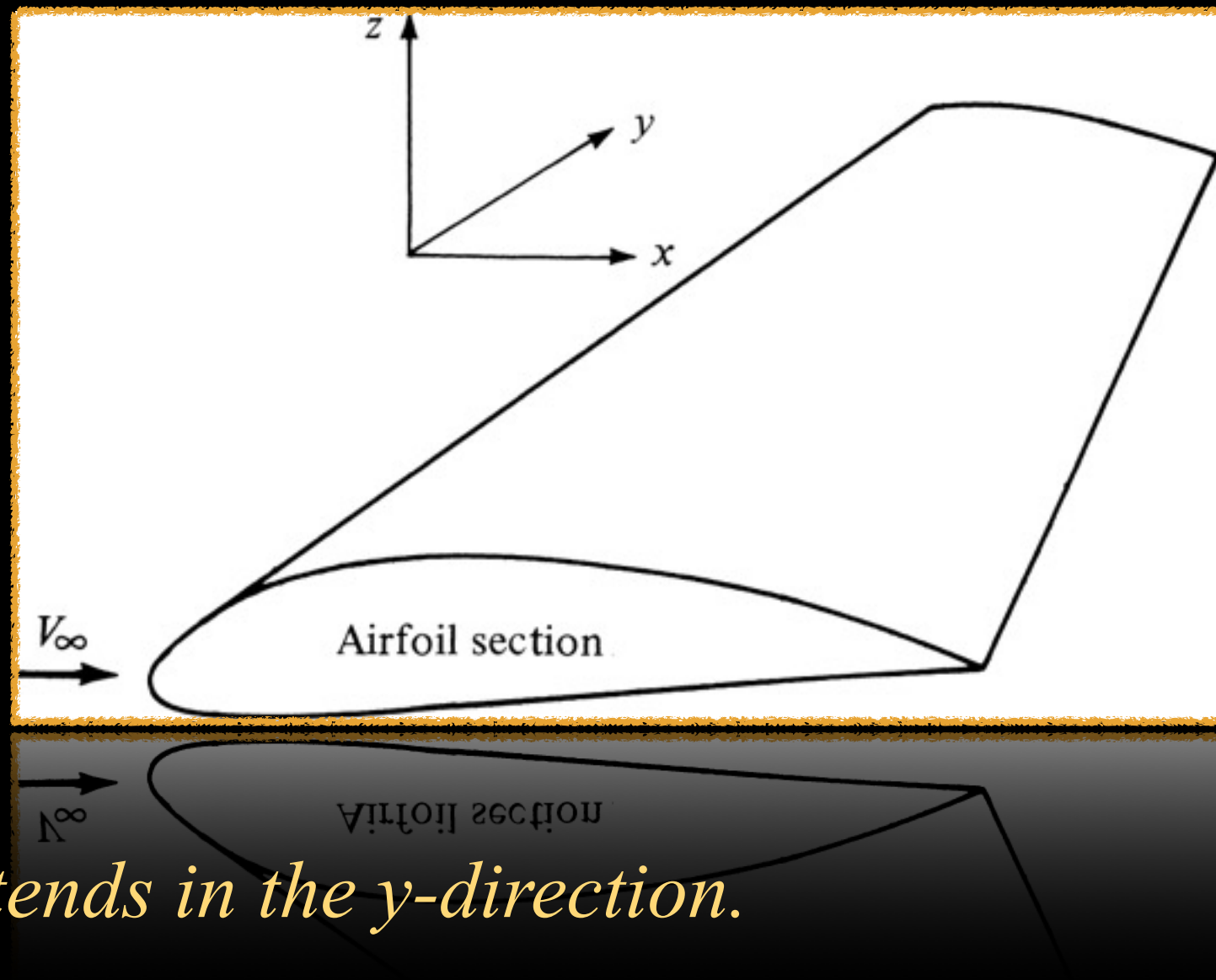
- ✱ Section of a wing - **airfoil**

- ✱ The complete **finite wing**

● *In this chapter, we will deal with airfoils.*

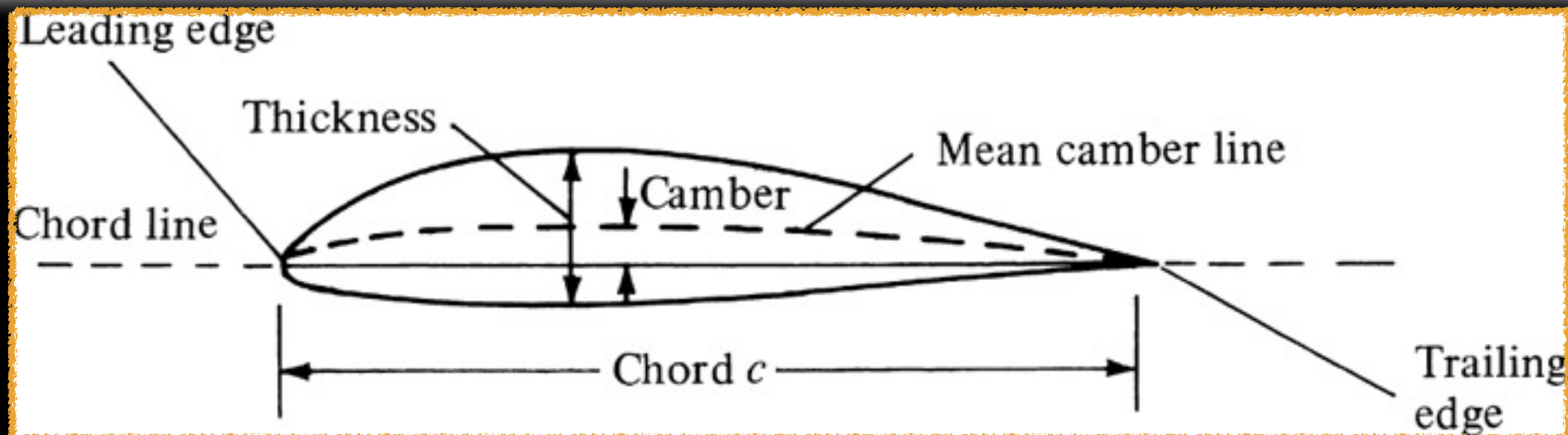
● *In the next chapter, we will deal with finite wings*

What is an Airfoil?



- *The wing extends in the y -direction.*
- *Any section of the wing cut by a plane parallel to xz -plane is called an **airfoil**.*

Airfoil Nomenclature

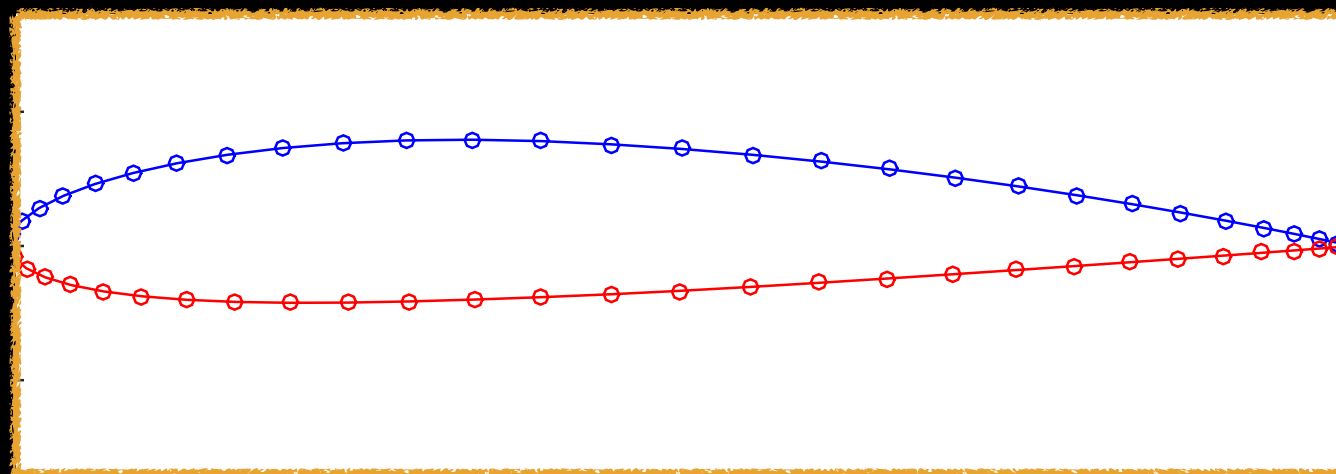


- *Mean camber line* - locus of points halfway between upper and lower surfaces
- *Chord* - Straight line connecting the leading and trailing edges.
- *Camber* - Maximum distance between the mean camber line and the chord line.
- *Thickness* - Distance between the upper and lower surfaces measured perpendicular to the chord line.
- *Leading-edge* is generally circular with a radius of $0.02c$.

NACA Airfoils

● NACA 4-digit series (e.g. NACA 2412)

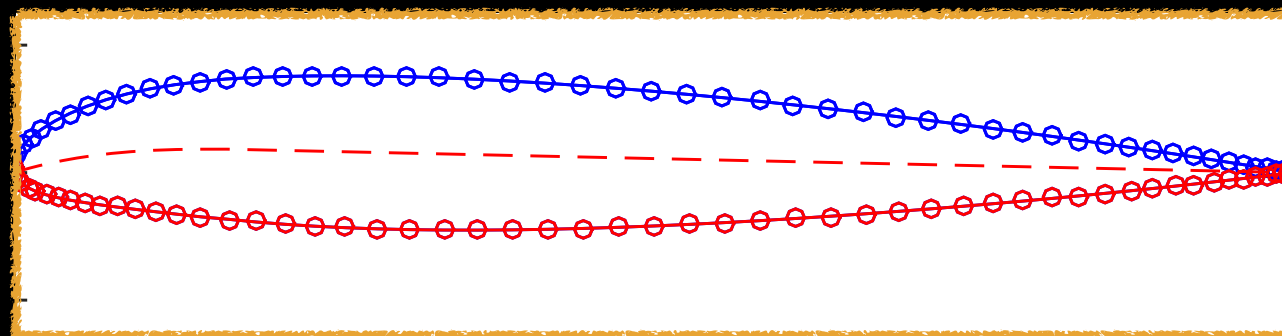
- ✱ *1st digit: maximum camber in hundredths of chord ($0.02c$ or 2%).*
- ✱ *2nd digit: location of maximum camber from the leading edge along the chord in tenths of chord ($0.4c$ or 40%).*
- ✱ *Last two digits: maximum thickness in hundredths of chord ($0.12c$ or 12%).*



NACA Airfoils

● NACA 5-digit series (e.g. NACA 23012)

- ✱ *1st digit x (3/2) gives design lift coefficient in tenths of chord (0.3).*
- ✱ *(Next two digits)/2 gives: location of maximum camber from the leading edge along the chord in hundredths of chord (0.15c or 15%).*
- ✱ *Last two digits: maximum thickness in hundredths of chord (0.12c or 12%).*

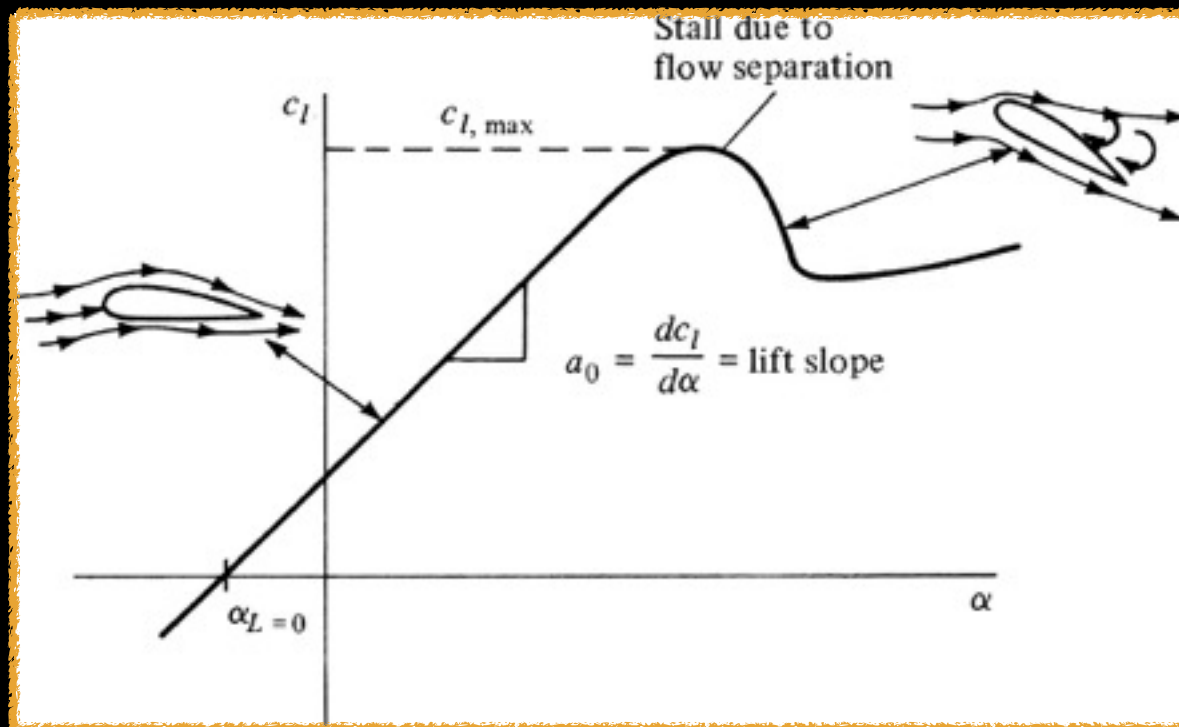


● *NACA 6-digit series (e.g. NACA 65-218)*

- ✱ *Very widely used laminar flow airfoils*
- ✱ *1st digit: represents the series.*
- ✱ *2nd digit: Gives location of minimum pressure in tenths of the chord from the leading edge ($0.5c$).*
- ✱ *3rd digit: Design lift coefficient in tenths of the chord (0.2).*
- ✱ *Last two digits: maximum thickness in hundredths of chord ($0.18c$ or 18%).*

Airfoil Characteristics

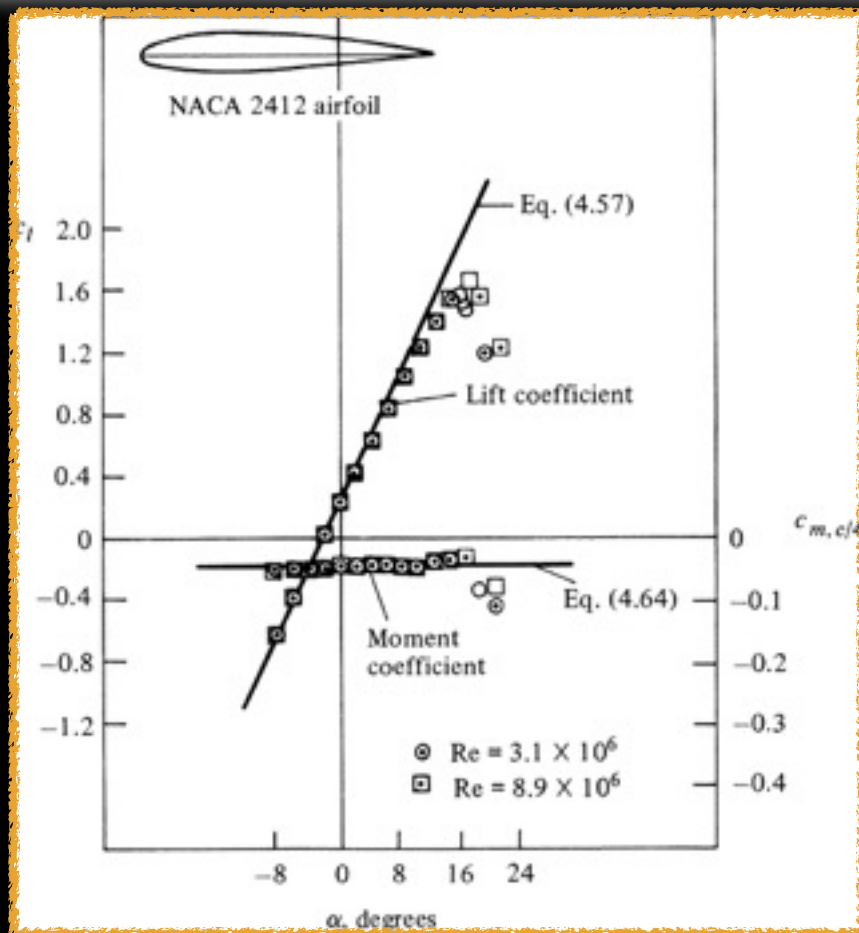
- Airfoil characteristics are typically lift, drag and moment coefficients.



- * Lift slope: lift coefficient varies linearly angle of angle of attack.
- * Zero-lift angle of attack: the value of angle of attack when lift is zero.

- Using inviscid flow theory, we can predict lift slope and zero angle of attack. But the maximum lift coefficient can only be calculated using viscous flow theory.

Lift and Moment Coefficients



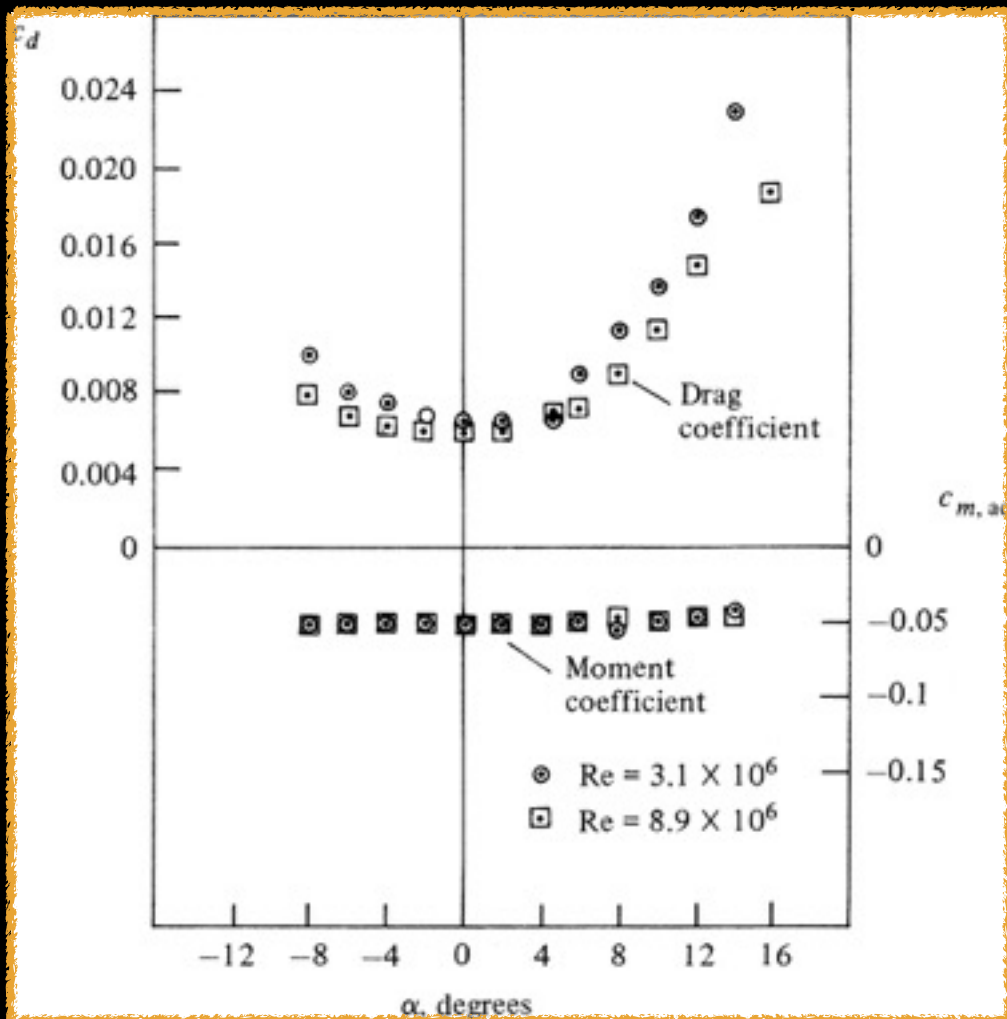
- * *These coefficients are taken about the quarter-chord location.*
- * *Note the difference for the two Re cases wrt lift-slope and max. lift coefficient.*
- * *Max. lift coefficient depends on Re (viscous effects).*

🏠 *Moment coefficient does not vary for this airfoil with angle of attack.*

✅ *Center of pressure: location where the resultant of a distributed load acts.*

✅ *Aerodynamic Center: location where the pitching moment is relatively constant with angle of attack.*

Profile Drag



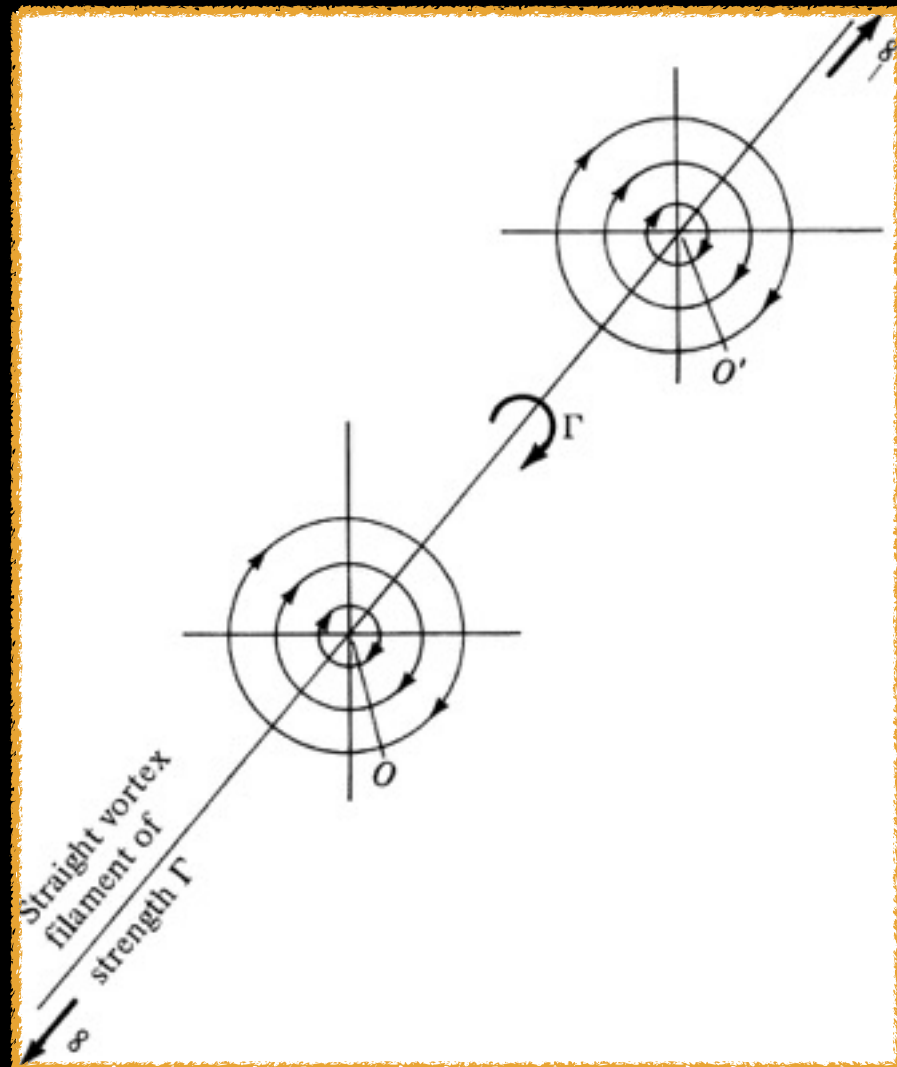
- * *Profile drag: friction drag + pressure drag*

- * *Profile drag is dependent on Re .*

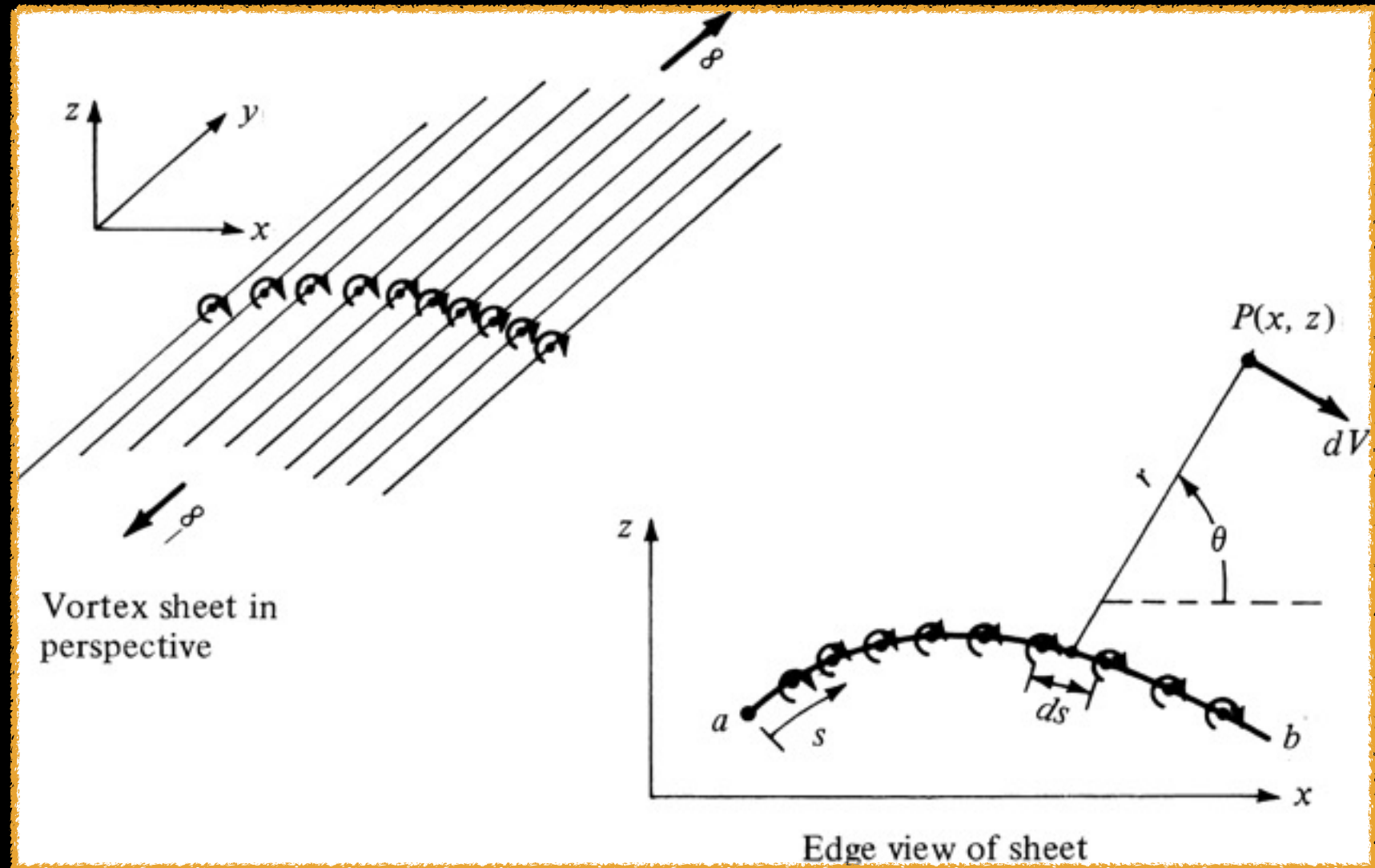
- * *Moments coefficients for moments taken about the aerodynamic center very weakly dependent (nearly independent) on the angle of attack for NACA 2412.*

 *Examples!!*

Theoretical Solutions for Low-Speed Airfoils

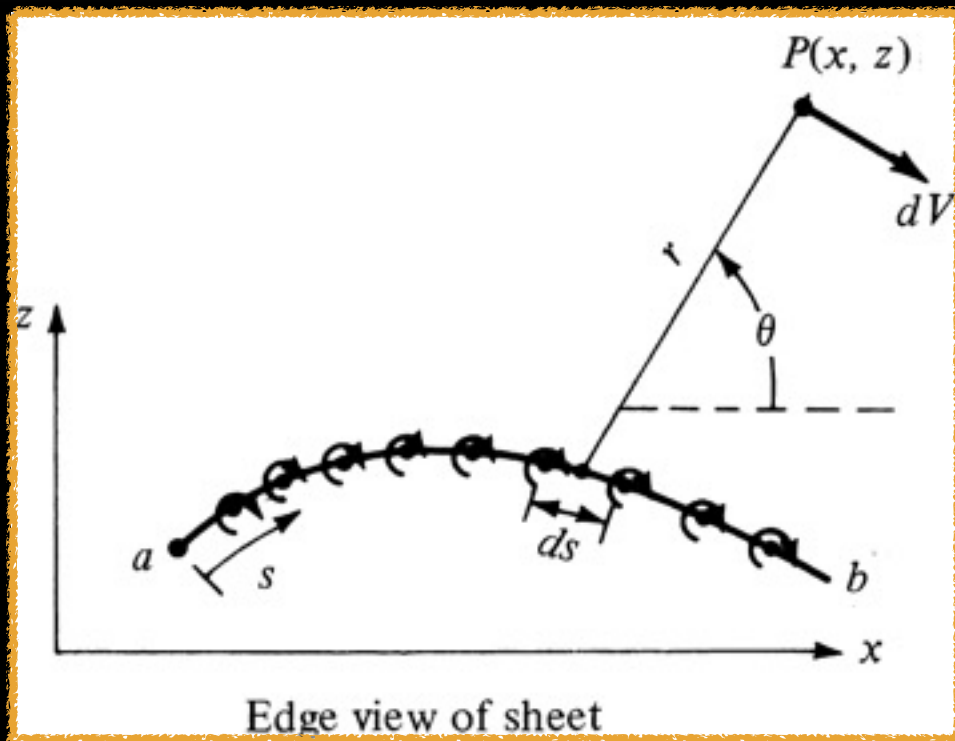


Vortex Filament



Vortex Sheet

Theoretical Solutions for Low-Speed Airfoils



induced velocity $dV = -\frac{\gamma ds}{2\pi r}$

velocity potential $d\Phi = -\frac{\gamma ds}{2\pi} \theta$

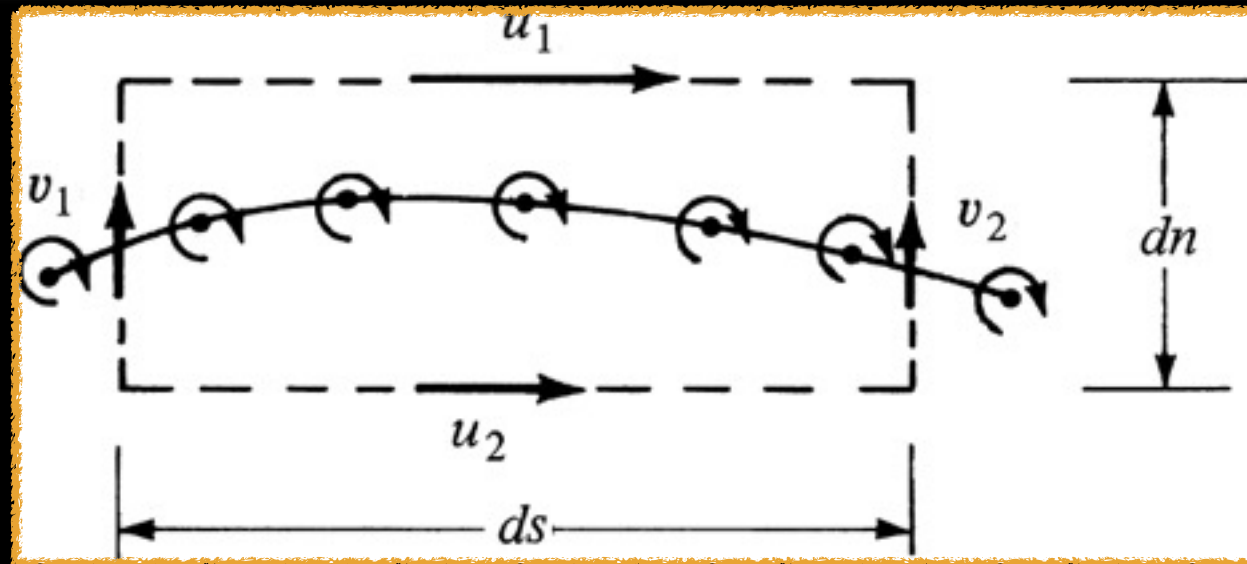
- Velocity potential at P due to entire vortex sheet is:

$$\Phi(x, z) = -\frac{1}{2\pi} \int_a^b \theta \gamma ds$$

- Circulation around entire vortex sheet (sum of strengths of elemental vortices) is:

$$\Gamma = \int_a^b \gamma ds$$

Circulation around an Airfoil



$$\Gamma = -(v_2 dn - u_1 ds - v_1 dn + u_2 ds) = \gamma ds$$

$$\Gamma = (u_1 - u_2)ds + (v_1 - v_2)dn = \gamma ds$$

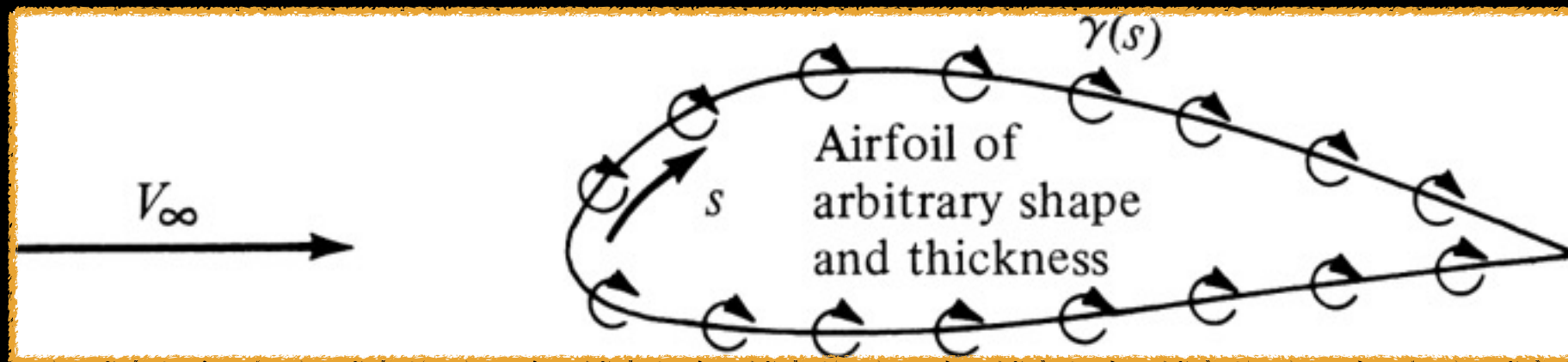
$$\text{assuming } dn \rightarrow 0, \gamma ds = (u_1 - u_2)ds$$

$$\Rightarrow \Gamma = u_1 - u_2$$

■ The local jump in tangential velocity across a vortex sheet is equal to the local sheet strength.

Circulation around an Airfoil

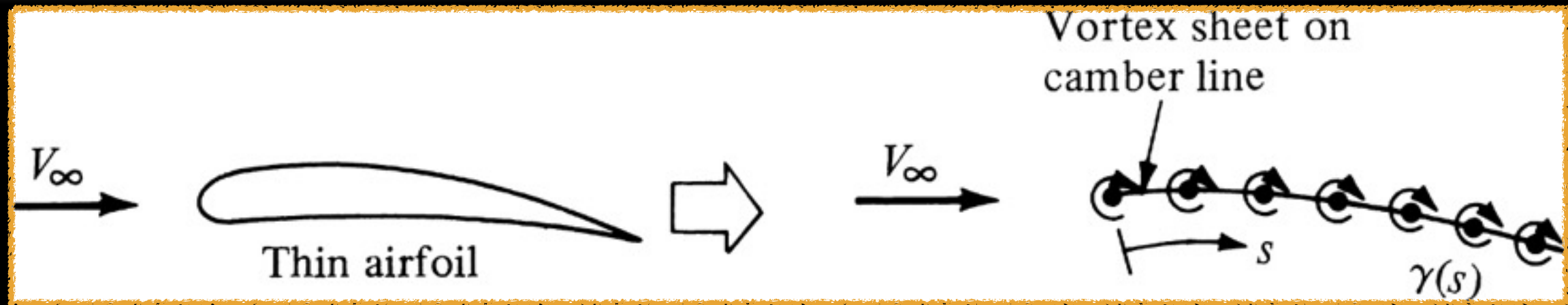
- *The concept of vortex sheet is instrumental in the analysis of low-speed characteristics of an airfoil.*



- * *Replace airfoil surface with a vortex sheet of variable strength.*
- * *Calculate strength of elemental vortices as a function of 's' such that when vortex induced velocity is added to freestream, the streamline would represent the airfoil surface.*
- * *The circulation around the airfoil is then given by:*

$$\Gamma = \int \gamma ds \Rightarrow L' = \rho_\infty V_\infty \Gamma \quad (\text{Kutta-Joukowski theorem})$$

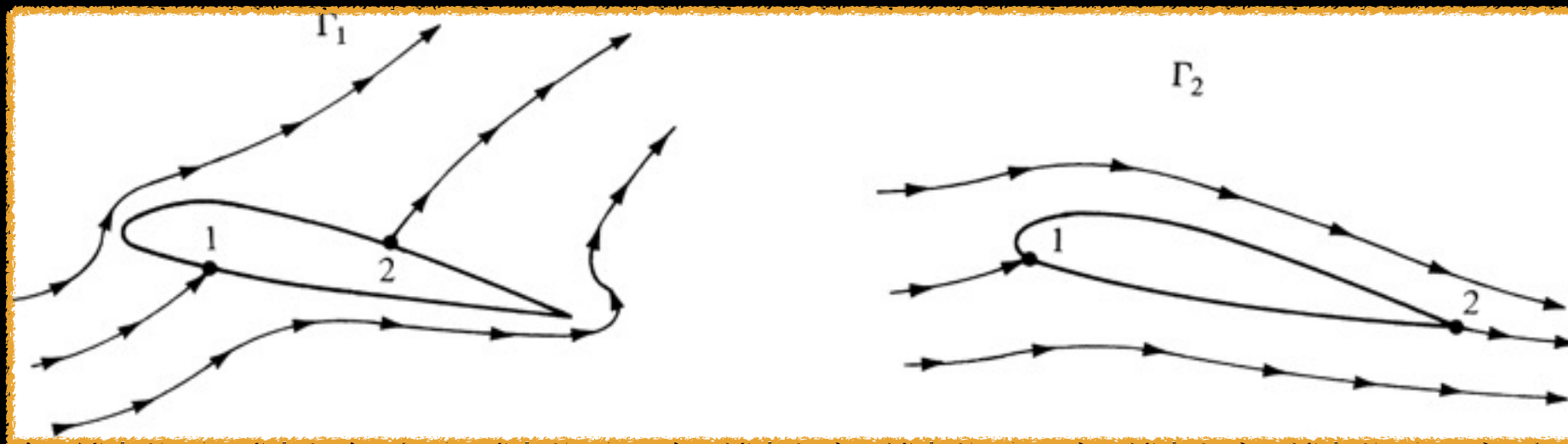
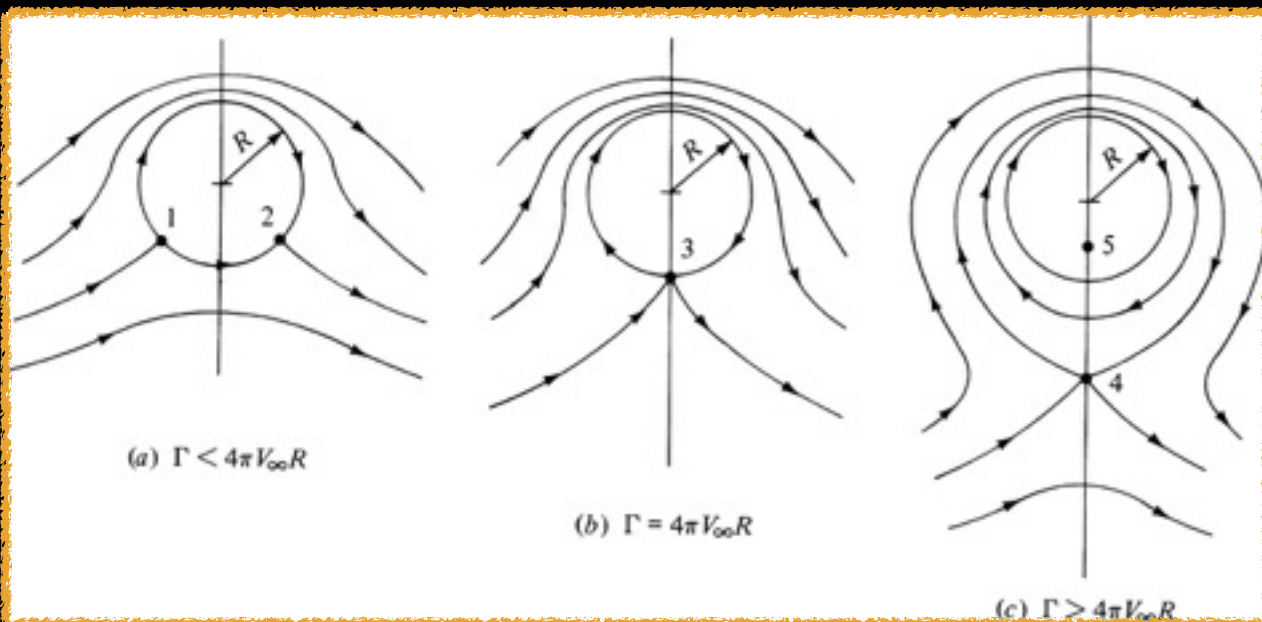
Thin Airfoil Approximation



- * *Imagine the airfoil is made very thin (top and bottom surfaces coincide).*
- * *Airfoil can be represented with a single vortex sheet distributed over the camber line.*
- * *The strength of the vortex sheet can be calculated such that when the induced velocity is added to the free stream velocity, the camber line become a streamline of the flow.*

The Kutta Condition

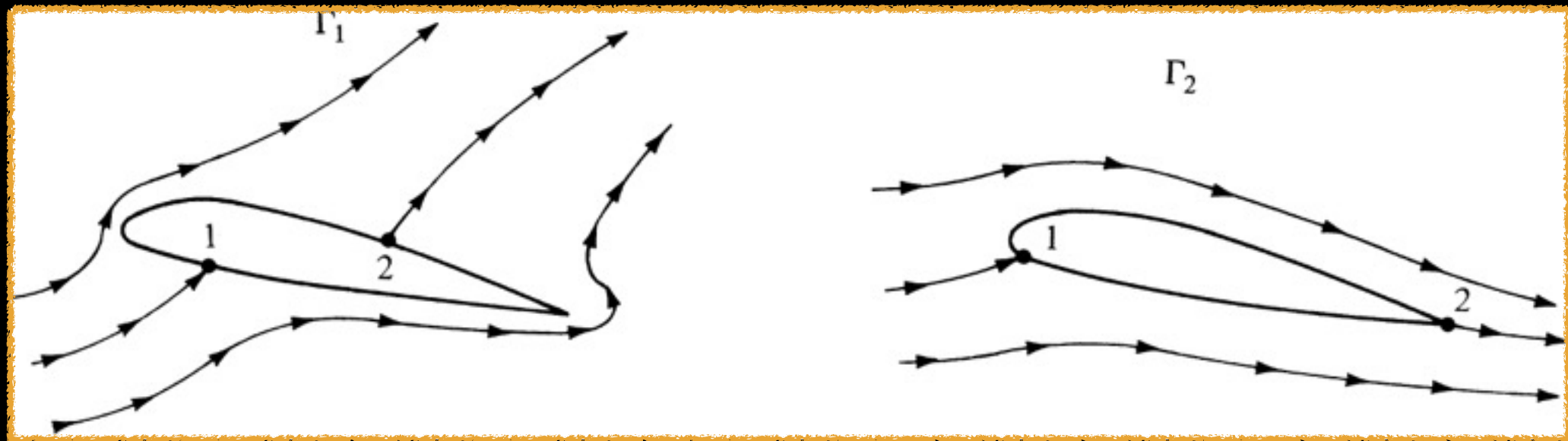
* *Infinite number of potential flow solutions are possible depending on the choice of circulation magnitude.*



* *Similarly, for an airfoil, infinite number of potential flow solutions are possible.*

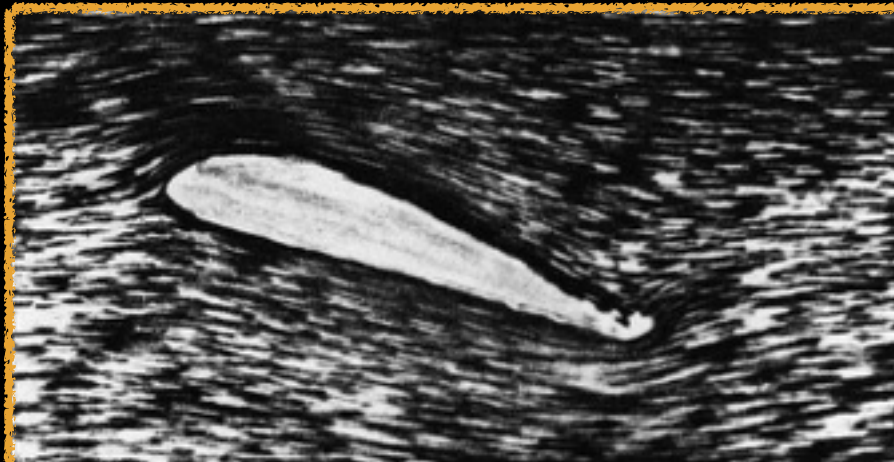
* *So, which one down pick?*

The Kutta Condition



- * *We know, a given angle of attack produces a single value of lift.*
- * *So, which ‘gamma’ does nature choose?*
- * *It has to fix the value of ‘gamma’.*
- * *To find out, let us look at an airfoil set into motion from a state of rest.*

The Kutta Condition

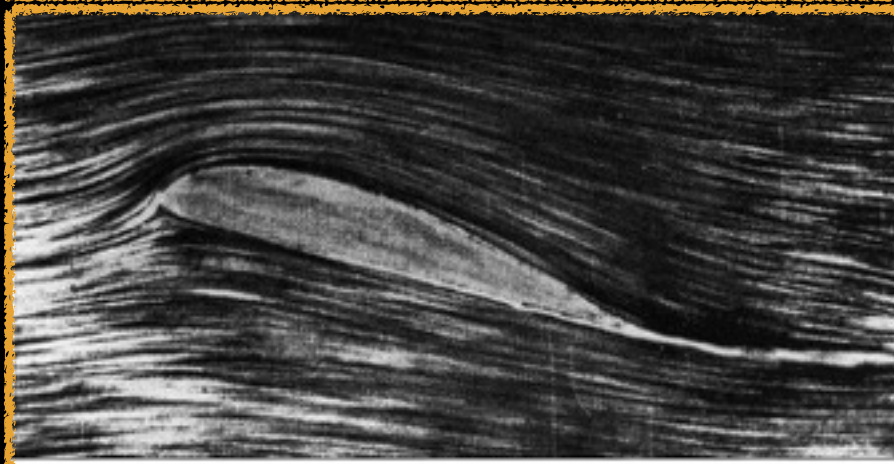


(a)



(b)

Source: Prandtl and Tietjens, Reference 8

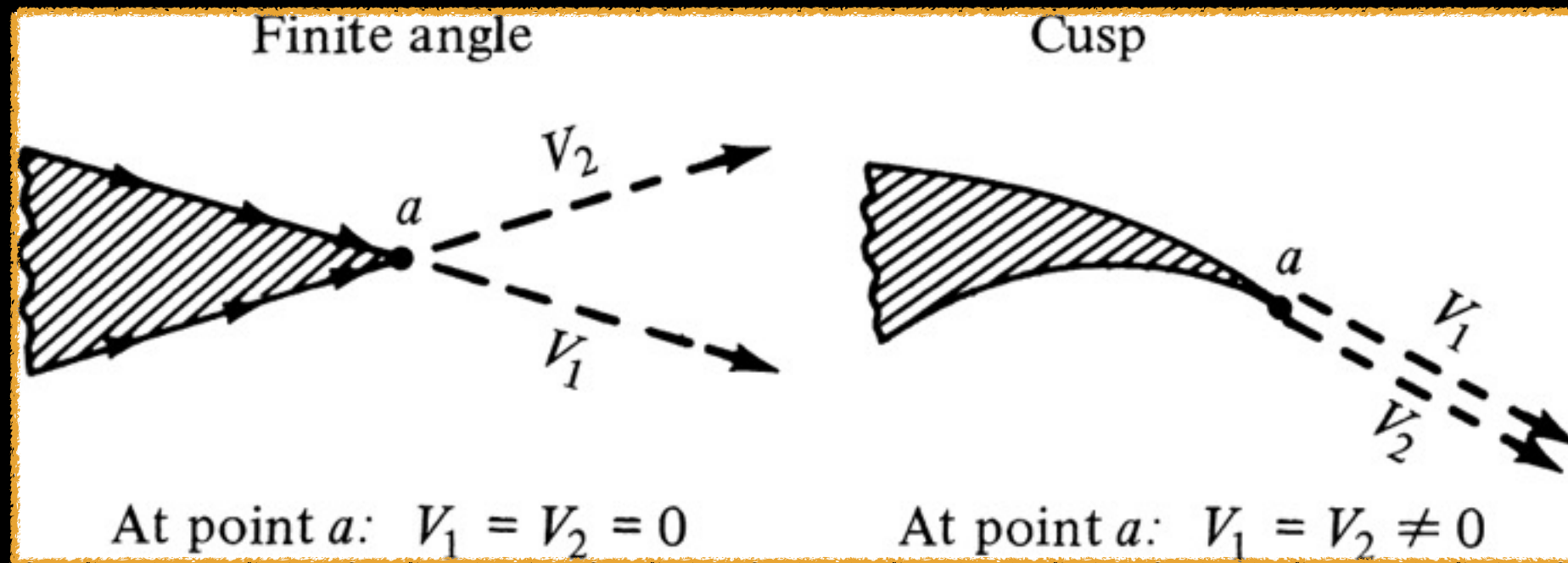


(c)

Source: Prandtl and Tietjens, Reference 8

- * *In (a), flow has just started.*
- * *It tries to curl around the TE.*
- * *Velocity becomes infinitely large at the sharp corner. Realistically impossible.*
- * *Such flow is not tolerated very long by nature.*
- * *So, flow leaves the top and bottom surfaces of the airfoil smoothly.*
- * *Therefore, nature adopts that value of circulation which results in a smooth flow at the trailing edge - Kutta condition.*

The Kutta Condition



- * *If the trailing edge angle is finite, then it is a stagnation point.*
- * *If the trailing edge is a cusp, velocities leaving the top and bottom are finite and equal in magnitude and direction.*
- * *In terms of vortex sheet, at the trailing edge:*

$$\gamma = \gamma(a) = V_1 - V_2 = 0$$

Lift Without Friction?

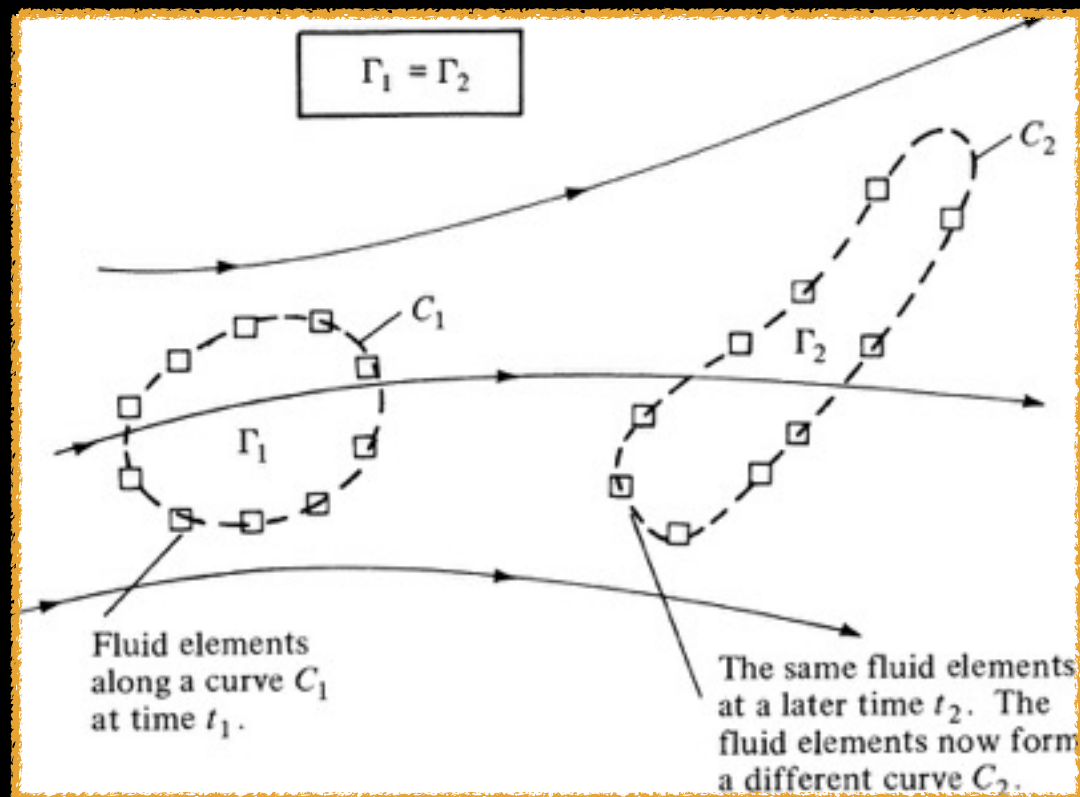
- * *We know that lift on a airfoil is primarily due to surface pressure distribution (acts normal) and not due to shear stress (tangential).*
- * *However, in a perfectly inviscid world, an airfoil would not produce lift. Sounds contradictory!!*
- * *In reality, nature enforces the Kutta condition - i.e. the viscous boundary layer remains attached all the way to the TE.*



- * *Lift, which is created by surface pressure distribution (inviscid phenomenon) cannot exist in an inviscid world.*

Kelvin's Circulation Theorem

* *How does nature enforce the Kutta condition? How does it generate this circulation for a given airfoil?*



$$\Gamma_1 = - \int_{C_1} \mathbf{V} \cdot d\mathbf{s}; \quad \Gamma_2 = - \int_{C_2} \mathbf{V} \cdot d\mathbf{s}$$

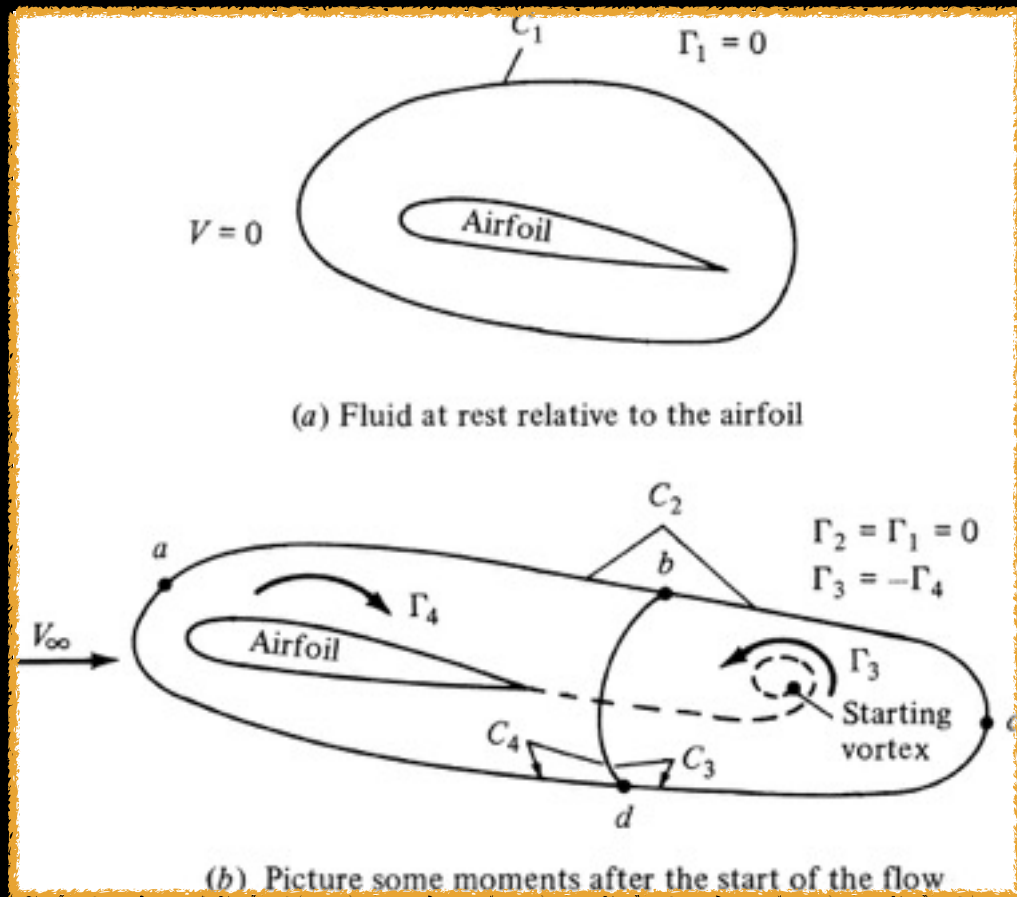
$$\Gamma_1 = \Gamma_2$$

$$\frac{d\Gamma}{Dt} = 0$$

* *Kelvin's Circulation Theorem: Time rate of change of circulation around a closed contour consisting of the same fluid elements is zero.*

Kelvin's Circulation Theorem

* So, how is circulation generated around an airfoil?



* Let $V = 0$, so circulation = 0, around C_1 .

* As the flow is started over the airfoil, large velocity gradients at the trailing edge generate vorticity that rolls up downstream - *starting vortex*.

* Due to Kelvin's theorem, the starting vortex has to induce equal and opposite circulation on the airfoil.

* Ideally, the starting vortex remains forever downstream.
Realistically, it dissipates due to viscous action.

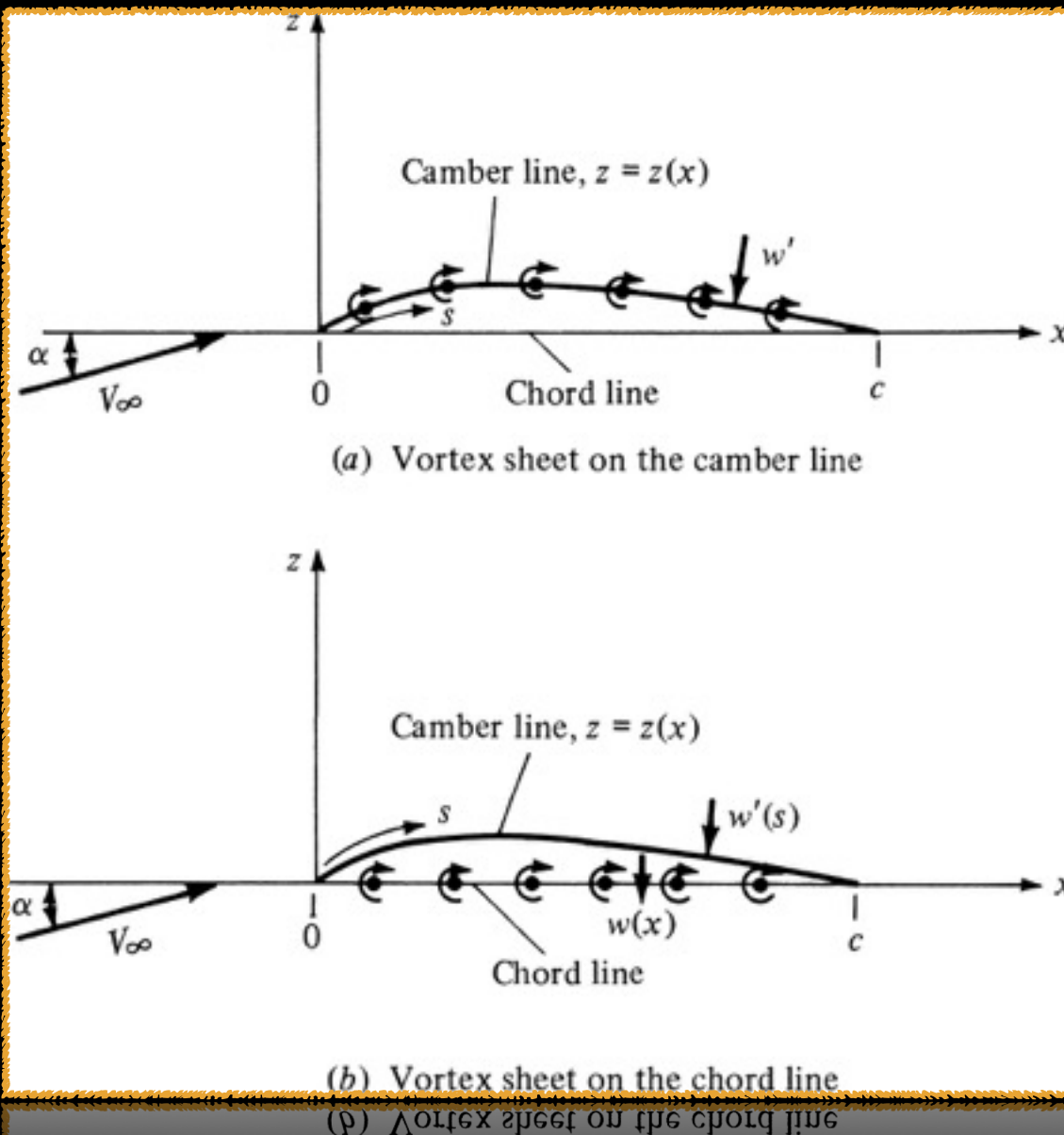
* Example!!

The Thin Airfoil Theory

Thin Airfoil

- ✱ *A vortex sheet placed along the camber line*
- ✱ *Our purpose is to calculate the variation in vorticity such that the camber line becomes a streamline.*
- ✱ *Kutta condition at the trailing edge should be satisfied.*
- ✱ *Once variation of vorticity that satisfies these conditions is found, total circulation is calculated by integrating vorticity from LE to TE.*
- ✱ *Lift can then be calculated using the Kutta-Joukowski theorem.*

The Thin Airfoil Theory



Assumptions

- ✱ *Thin airfoil, i.e. camber line is close to the chord line.*
- ✱ *Vortex sheet falls approximately on the chord line*

$$\gamma = \gamma(x)$$

$$\gamma(c) = 0$$



The strength of the vortex sheet on the chord line must be determined such that the camber line (not the chord line) is a streamline.

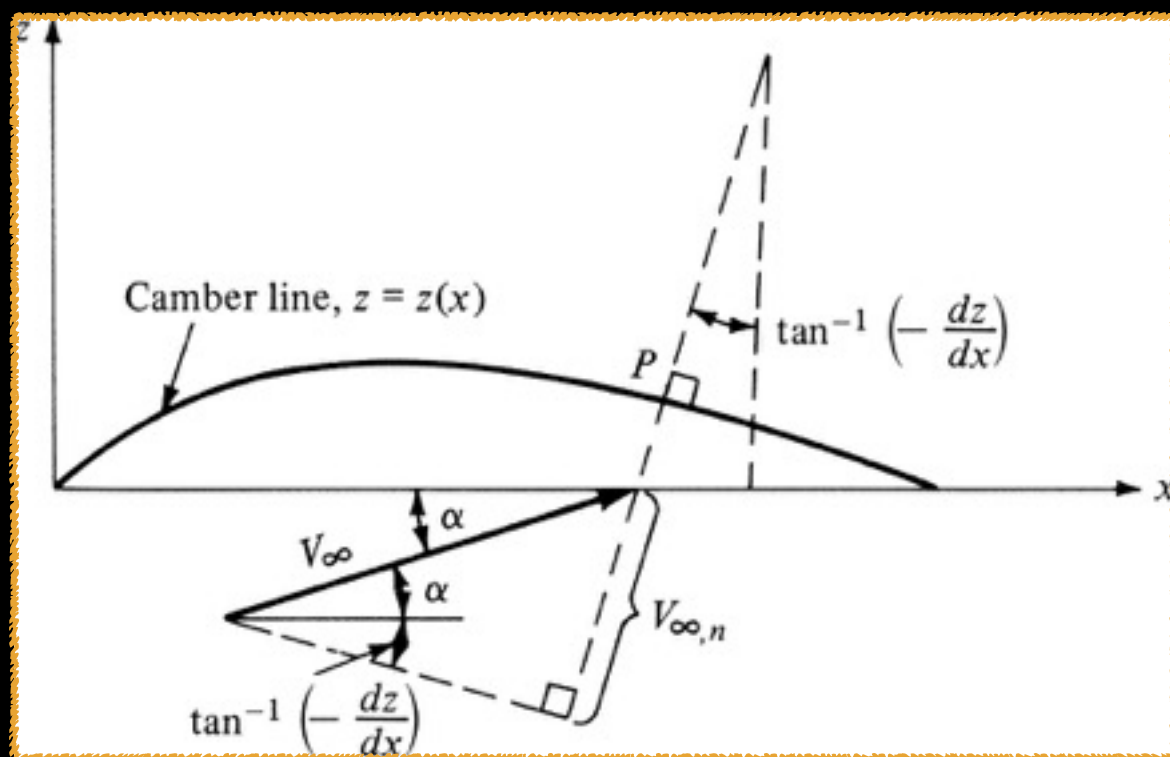
The Thin Airfoil Theory

For the camber line to be a streamline, the component of velocity normal to the camber line must be zero.

$$V_{\infty,n} + w'(s) = 0$$

$$V_{\infty,n} = V_{\infty} \sin \left[\alpha + \tan^{-1} \left(-\frac{dz}{dx} \right) \right]$$

For small angles of attack,
 $\sin \theta = \tan \theta = \theta$



$$V_{\infty,n} = V_{\infty} \left[\alpha - \left(\frac{dz}{dx} \right) \right]$$

The Thin Airfoil Theory

$$\int_0^c \frac{\gamma(\xi) d\xi}{(x - \xi)} = V_\infty \left[\alpha - \left(\frac{dz}{dx} \right) \right] \quad \text{fundamental equation of thin airfoil theory}$$

● In the above equation,

- ✱ A vortex sheet placed along the camber line
- ✱ Our purpose is to calculate the variation in vortex strength such that the camber line becomes a streamline.
- ✱ The central challenge is to calculate the vortex strength variation subject to the Kutta condition, i.e.

$$\gamma(c) = 0$$

The Thin Airfoil Theory

● Consider a symmetric airfoil

✱ No camber.

✱ Camber line is coincident with the chord line ($dz/dx = 0$). Therefore,

$$\frac{1}{2\pi} \int_0^c \frac{\gamma(\xi) d\xi}{x - \xi} = V_\infty \alpha$$

$$\xi \rightarrow \theta : \quad \xi = \frac{c}{2}(1 - \cos\theta)$$

$$x \rightarrow \theta_0 : \quad x = \frac{c}{2}(1 - \cos\theta_0)$$

solving, $\gamma(\theta) = 2\alpha V_\infty \frac{1 + \cos\theta}{\sin\theta}$

The Thin Airfoil Theory

- Now, total circulation around the airfoil is given by:

$$\Gamma = \int_0^c \gamma(\xi) d\xi = \frac{c}{2} \int_0^\pi \gamma(\theta) \sin\theta d\theta$$

- Simplifying,

$$\Gamma = \pi\alpha c V_\infty$$

- Lift can now be calculated using the Kutta-Jouski theorem:

$$L = \rho_\infty V_\infty \Gamma = \pi\alpha c \rho_\infty V_\infty^2$$

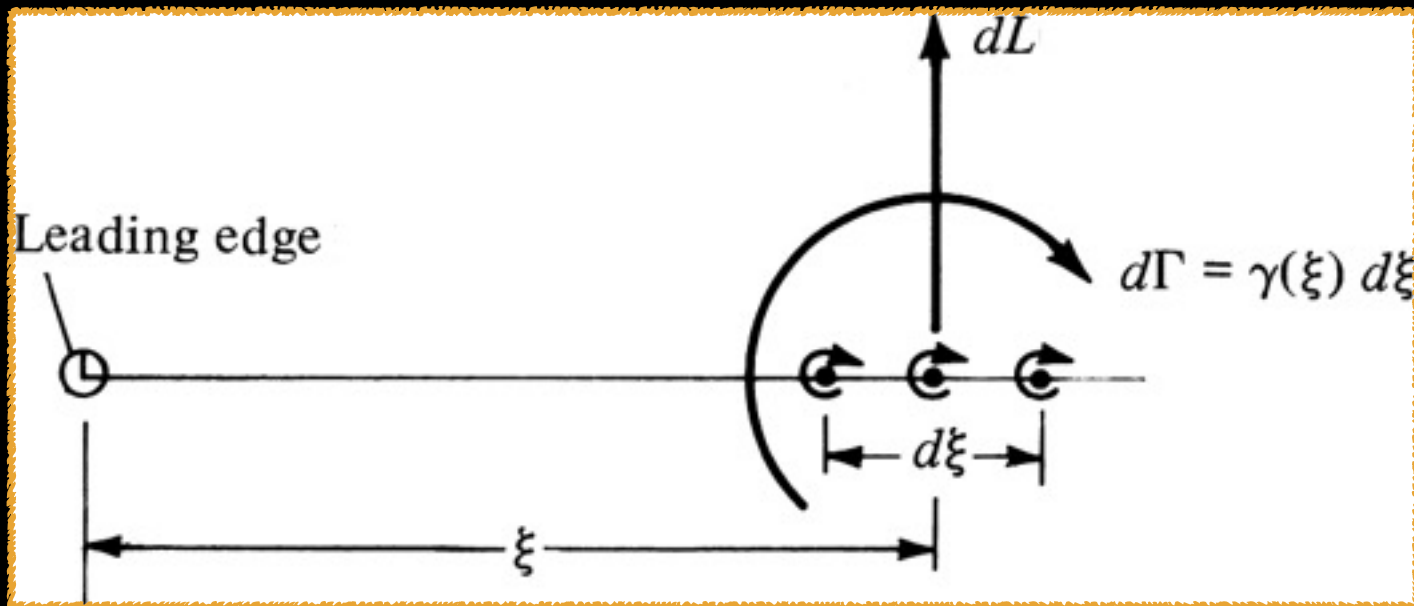
- Lift Coefficient is then:

$$c_l = \frac{L'}{q_\infty S} = 2\pi\alpha$$

Lift Coefficient is proportional to the angle of attack.

The Thin Airfoil Theory

Now, let us calculate moment about the leading edge:



$$dL = \rho_{\infty} V_{\infty} d\Gamma$$

$$dM = -\xi(dL)$$

(about LE)

Total moment about the LE due to the entire vortex sheet is:

$$M'_{LE} = - \int_0^c \xi(dL) = -\rho_{\infty} V_{\infty} \int_0^c \xi \gamma(\xi) d\xi = -q_{\infty} c^2 \frac{\pi \alpha}{2}$$

The Thin Airfoil Theory

- *The moment coefficient about the leading edge:*

$$c_{m,le} = \frac{M'_{LE}}{q_{\infty} S c} = -\frac{\pi \alpha}{2} = -\frac{c_l}{4}$$

- *Moment coefficient about the quarter-chord point is:*

$$c_{m,c/4} = c_{m,le} + \frac{c_l}{4} = 0$$

- *Center of pressure: Moments are zero*
- *Aerodynamic Center: Moments are independent of angle of attack.*
- *For a symmetrical airfoil, the quarter-chord location is both the center of pressure and the aerodynamic center.*

The Cambered Airfoil

$$\int_0^c \frac{\gamma(\xi)d\xi}{(x-\xi)} = V_\infty \left[\alpha - \left(\frac{dz}{dx} \right) \right]$$

- For a cambered airfoil, dz/dx is finite.

$$\frac{1}{2\pi} \int_0^\pi \frac{\gamma(\theta)\sin\theta d\theta}{\cos\theta - \cos\theta_0} = V_\infty \left(\alpha - \frac{dz}{dx} \right)$$

- Solving, we obtain:

$$\gamma(\theta) = 2V_\infty \left(A_0 \frac{1 + \cos\theta}{\sin\theta} + \sum_{n=1}^{\infty} A_n \sin(n\theta) \right)$$

- A_0 and A_n are Fourier coefficients that depend on shape of the camber line and angle of attack.

The Cambered Airfoil

- *The Fourier coefficients are:*

$$A_n = \alpha - \frac{1}{\pi} \int_0^\pi \frac{dz}{dx} d\theta_0$$

$$A_n = \frac{2}{\pi} \int_0^\pi \frac{dz}{dx} \cos(n\theta_0) d\theta_0$$

- *Now, circulation due to the entire vortex sheet from LE to TE is:*

$$\Gamma = \int_0^c \gamma(\xi) d\xi = cV_\infty \left(\pi A_0 + \frac{\pi}{2} A_1 \right)$$

- *The lift coefficient is then given by:*

$$c_l = 2\pi \left[\alpha + \frac{1}{\pi} \int_0^\pi \frac{dz}{dx} (\cos\theta_0 - 1) d\theta_0 \right]$$

The Cambered Airfoil

- Similarly, the moment coefficient about the LE is given by:

$$C_{m,le} = - \left[\frac{C_l}{4} + \frac{\pi}{4} (A_1 - A_2) \right]$$

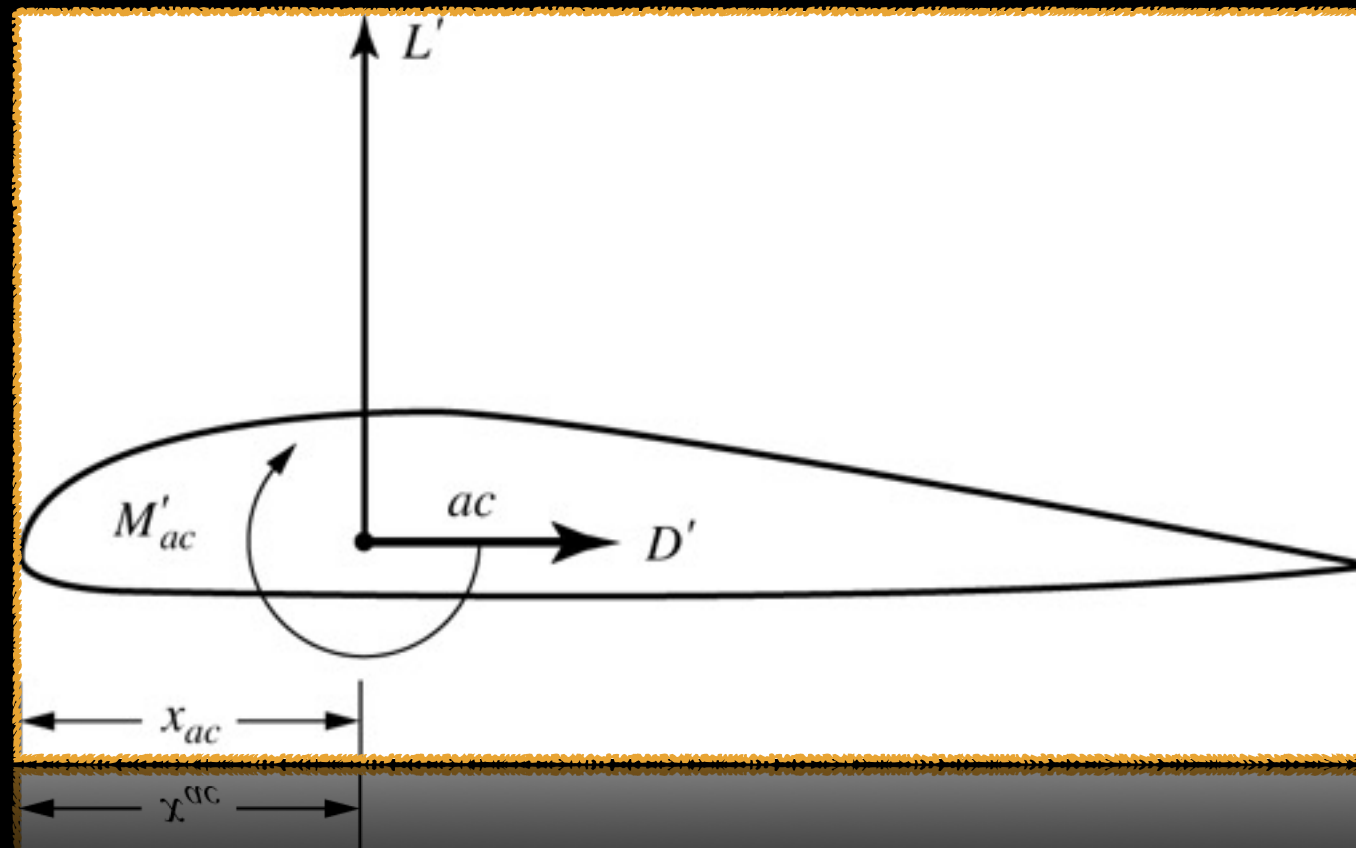
- The moment coefficient about the quarter-chord location is given by:

$$C_{m,c/4} = - \left[\frac{\pi}{4} (A_2 - A_1) \right]$$

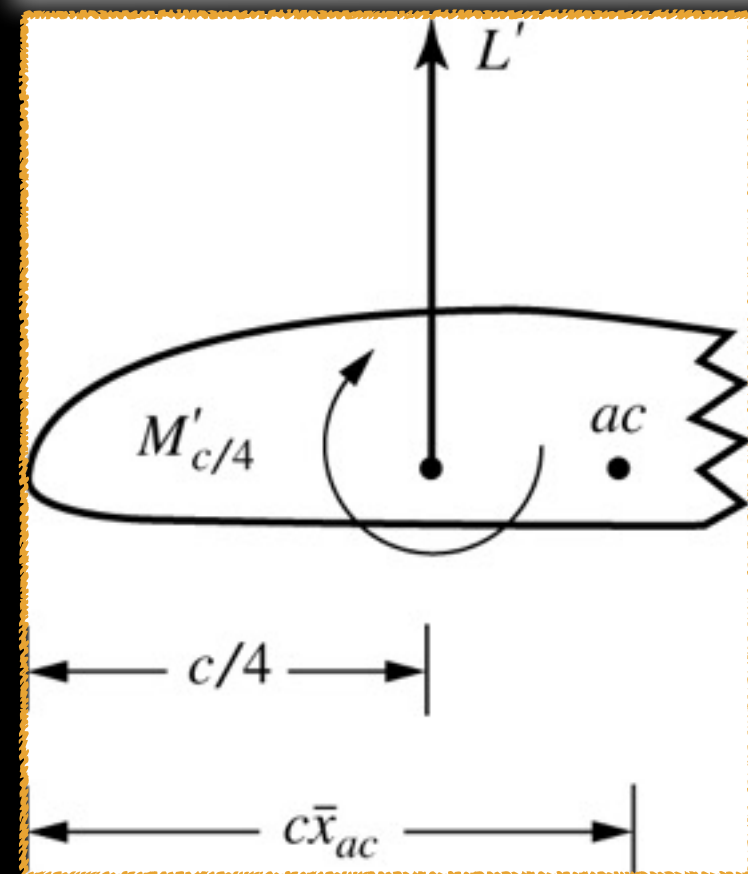
- From the above equation, we can see that:
 - ✱ The quarter-chord is not the center of pressure for a cambered airfoil.
 - ✱ However, since the moment is independent of the angle of attack, the quarter-chord is the theoretical aerodynamic center for a cambered airfoil.

The Aerodynamic Center

- *The aerodynamic center is that point on a body about which the aerodynamically generated moment is independent of the angle of attack.*
- *For most airfoils, it is close to, but not exactly at the quarter-chord location.*
- *So, how do we calculate its location?*



The Aerodynamic Center



- Taking moments about the aerodynamic center, we get:

$$M'_{ac} = L'(c\bar{x}_{ac} - c/4) + M'_{c/4}$$

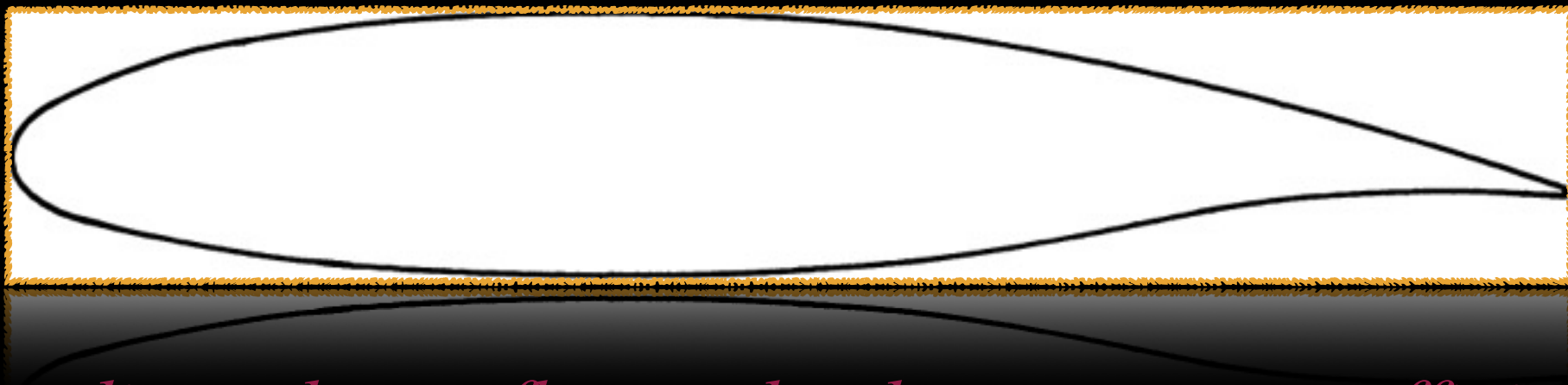
$$\text{or, } c_{m,ac} = c_l(\bar{x}_{ac} - 0.25) + c_{m,c/4}$$

- Rearranging the above equation and recognizing that the slopes of lift and moment coefficients are constants before stall, we have:

$$\bar{x}_{ac} = -\frac{m_0}{a_0} + 0.25 \quad \text{where, } \frac{dc_l}{d\alpha} = a_0; \quad \frac{dc_{m,c/4}}{d\alpha} = m_0$$

Modern Low Speed Airfoils

- *The standard NACA airfoils were based on experimental data in the 1930's and 1940's.*
- *New NASA airfoils were designed using source and vortex panel methods along with numerical prediction of viscous flow behavior. E.g. GA(W)-1 airfoil.*

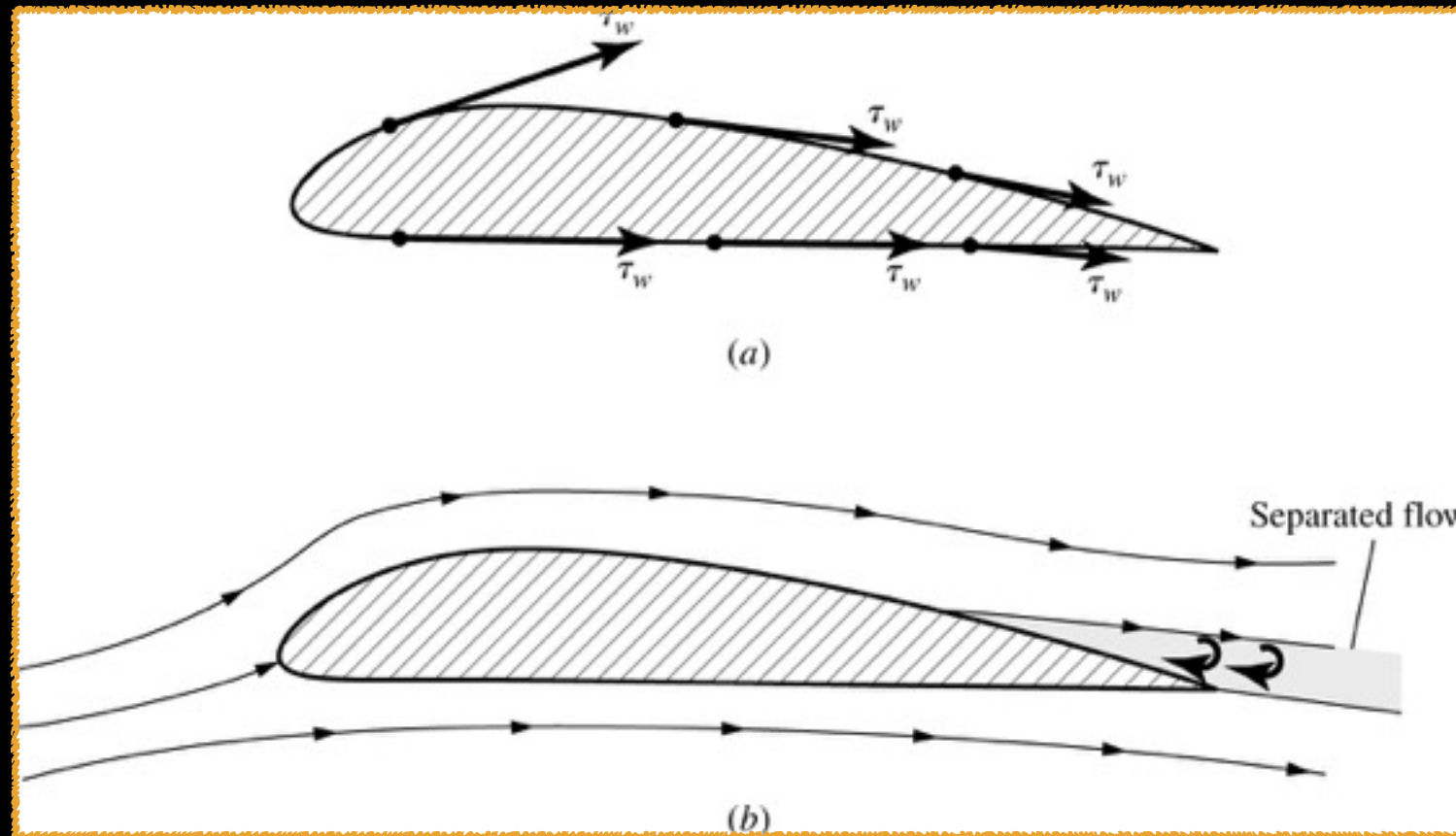


- * *Large leading edge to flatten the the pressure coefficient peak.*
- * *The trailing edge is cusped to increase the camber and loading.*
- * *The design discourages flow separation over the top surface leading to high lift coefficient.*

Viscous Flow: Airfoil Drag

- *Lift: Primarily due to pressure distribution on airfoil surface.*
 - * *Shear stress distribution in the lift direction is generally very small.*
 - * *Lift can therefore be accurately calculated assuming inviscid flow in conjunction with Kutta condition at the TE.*
- *Drag: Predicting drag using an inviscid approach results in zero drag (d'Alembert's paradox).*
 - * *However, when friction is included, this paradox is immediately removed.*

Viscous Flow: Airfoil Drag



● Skin Friction Drag

* Due to shear stress acting on the surface.

● Pressure Drag (form drag)

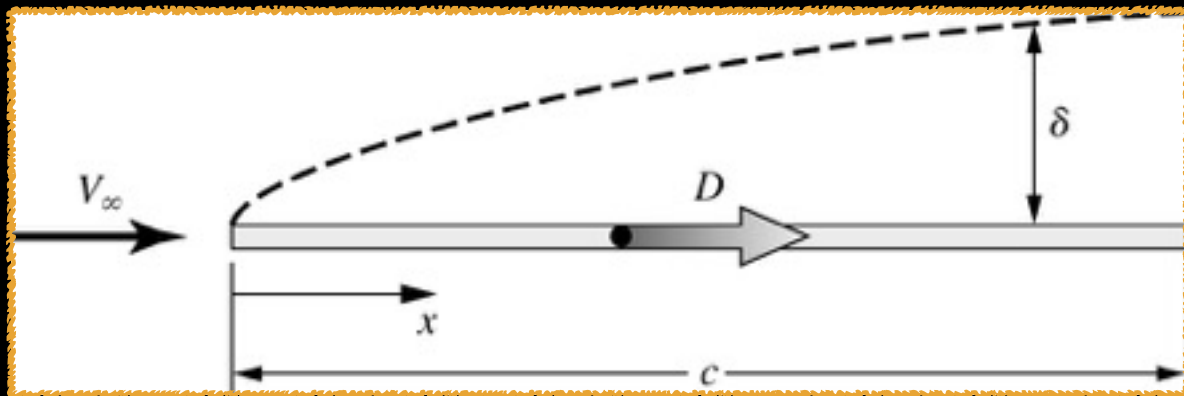
* Due to flow separation.

Skin-Friction Drag: Laminar Flow



Assume that skin-friction for airfoil is same as that for a flat plate.

- *The above assumption becomes more accurate for a thinner airfoil and small angles of attack.*



$$\delta = \frac{5.0x}{\sqrt{Re_x}}$$

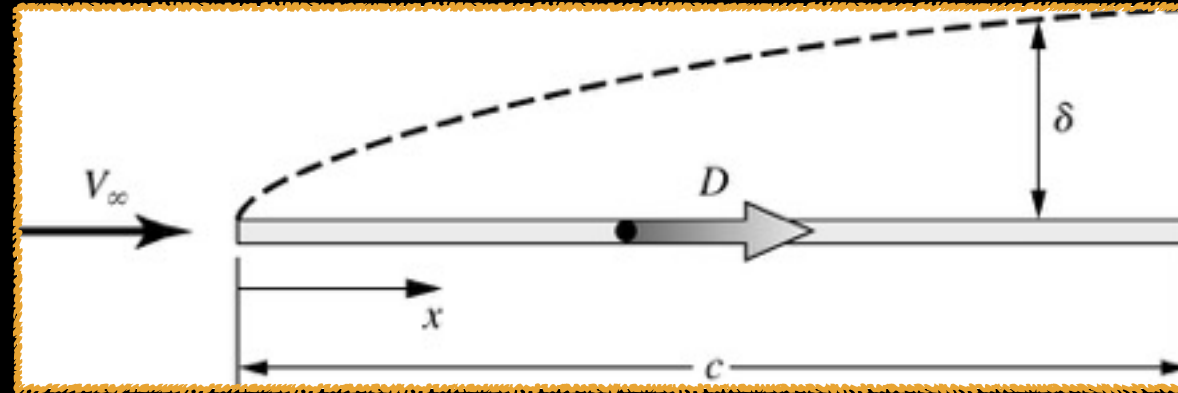
$$Re_x = \frac{\rho_e V_\infty x}{\mu_\infty}$$

- *The total skin-friction drag is given by:*

$$D_f = 2D_{f,top} = 2D_{f,bottom}$$

$$\text{where } c_f \equiv \frac{D_{f,top}}{q_\infty S} = \frac{D_{f,bottom}}{q_\infty S} = \frac{1.328}{\sqrt{Re_c}}$$

Skin-Friction Drag: Turbulent Flow

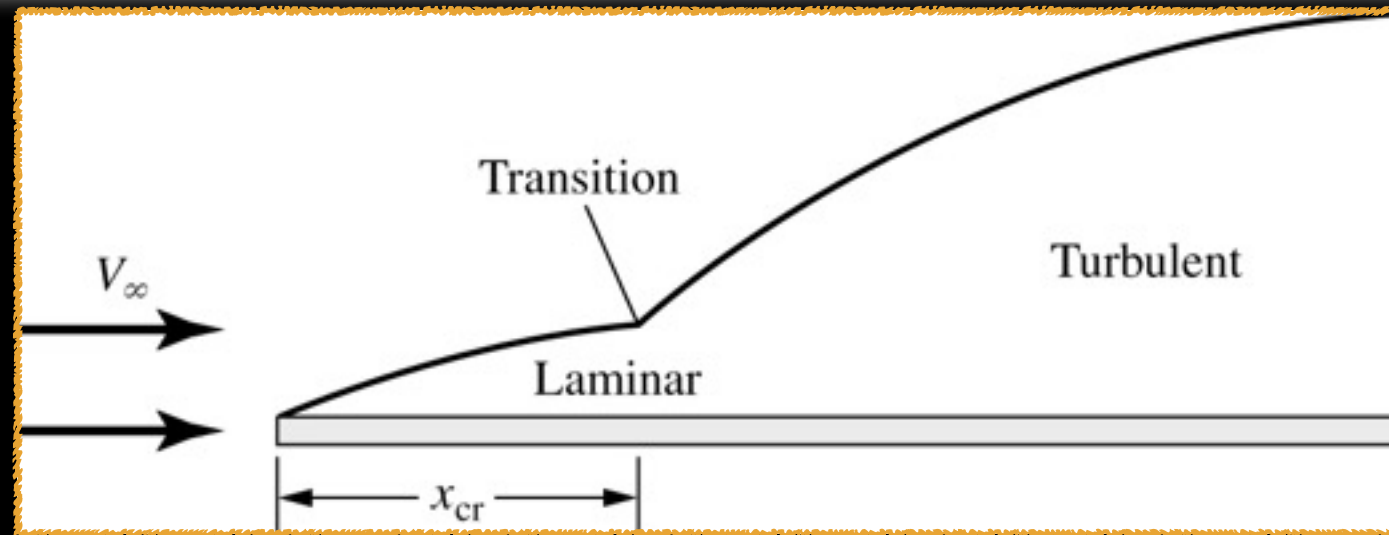


- *In contrast to laminar flow, there are no analytical solution for turbulent flow.*
- *All analyses of turbulent flow are approximate.*

$$\delta = \frac{0.37x}{Re_x^{1/5}}$$

$$C_f = \frac{0.074}{Re_c^{1/5}}$$

Skin-Friction Drag: Transition Flow

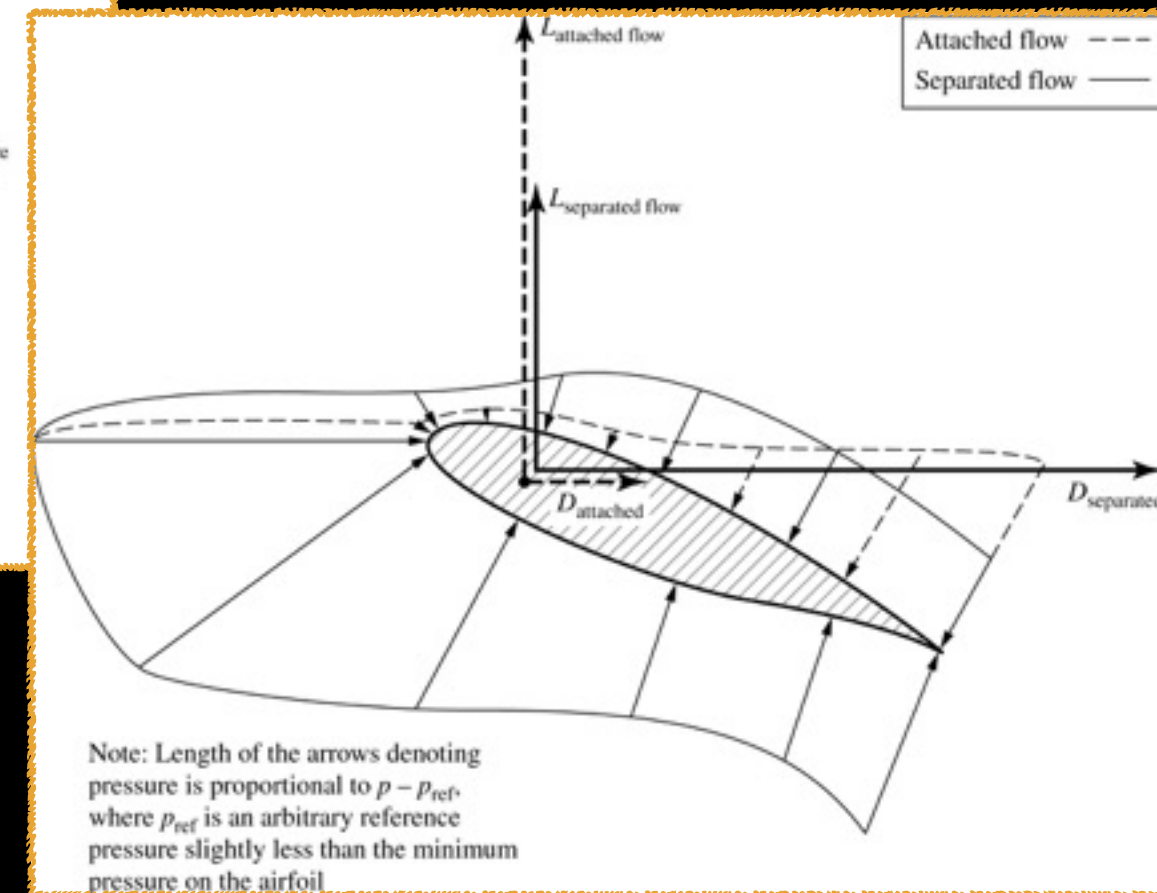
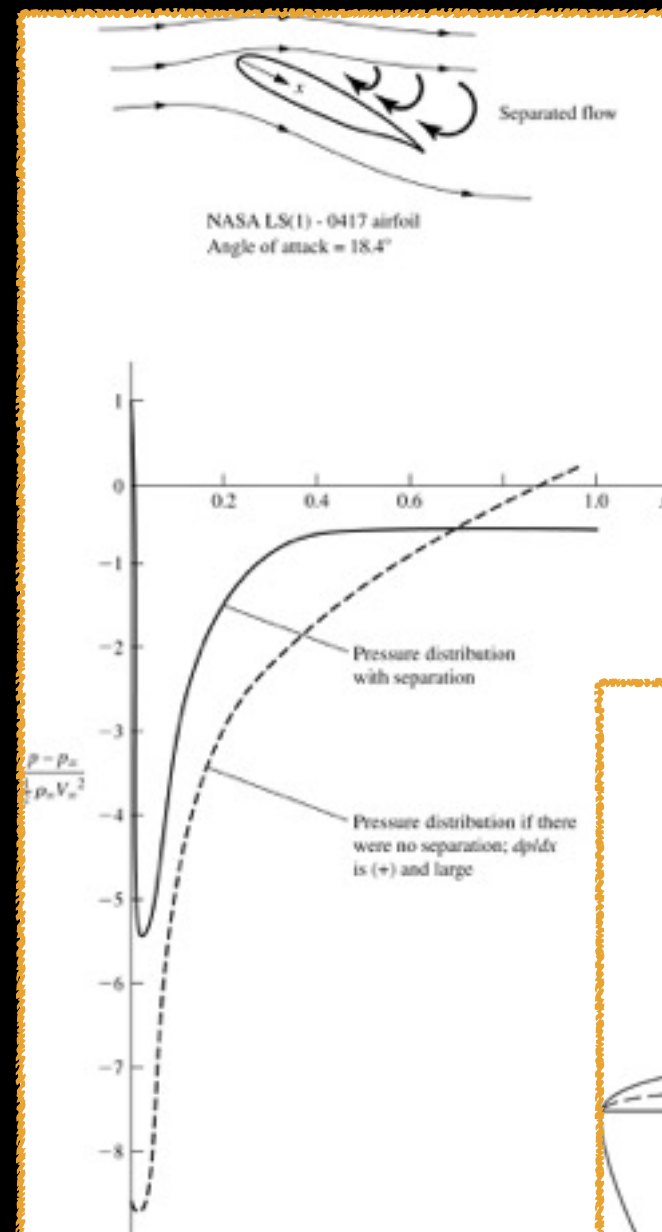
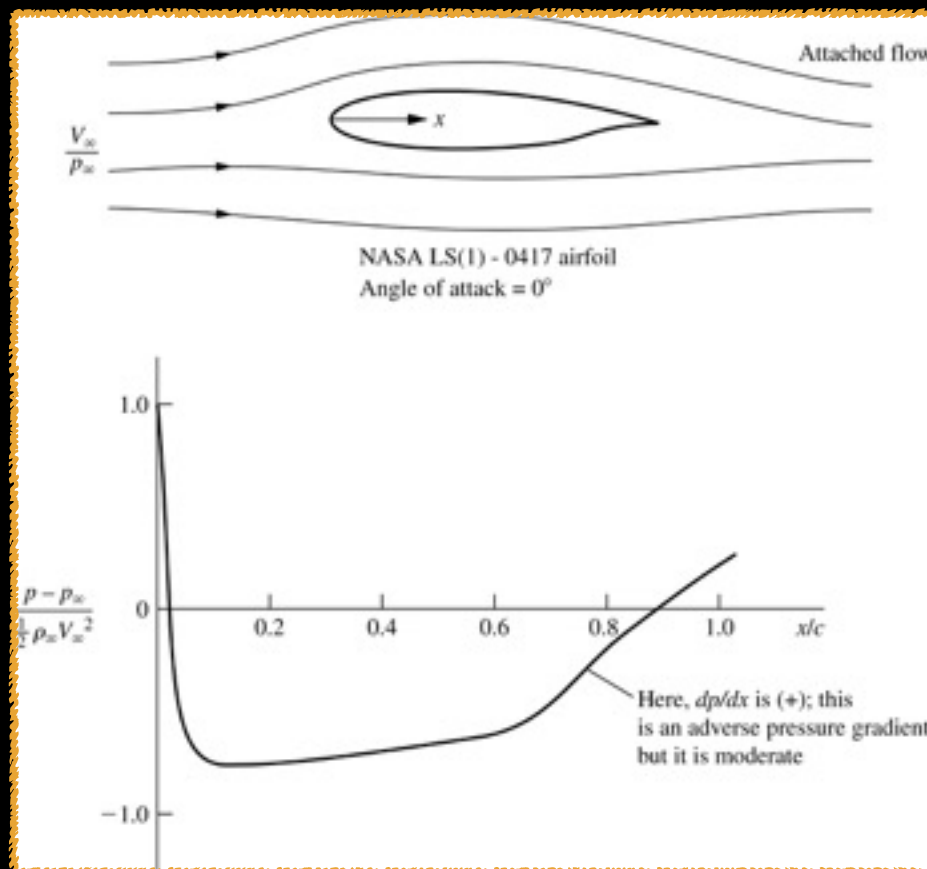


- *Flow always starts out as laminar at the leading edge, then becomes unstable and transitions into a turbulent flow.*
- *The value of x where transition takes place is the critical value x_{cr} .*

$$Re_{x_{cr}} = \frac{\rho_\infty V_\infty x_{cr}}{\mu_\infty}$$

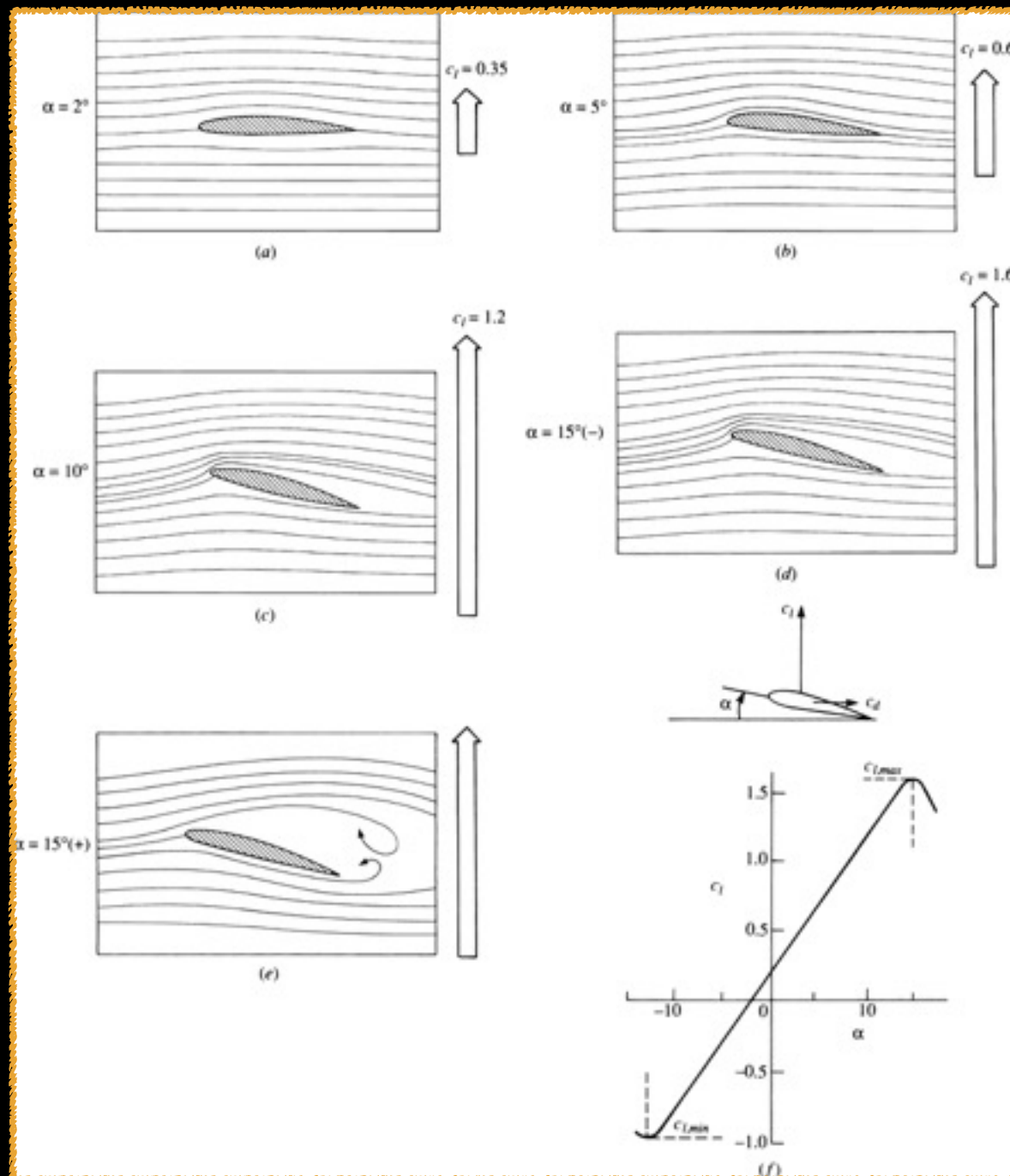
Flow Separation

 *Pressure drag is caused by flow separation.*



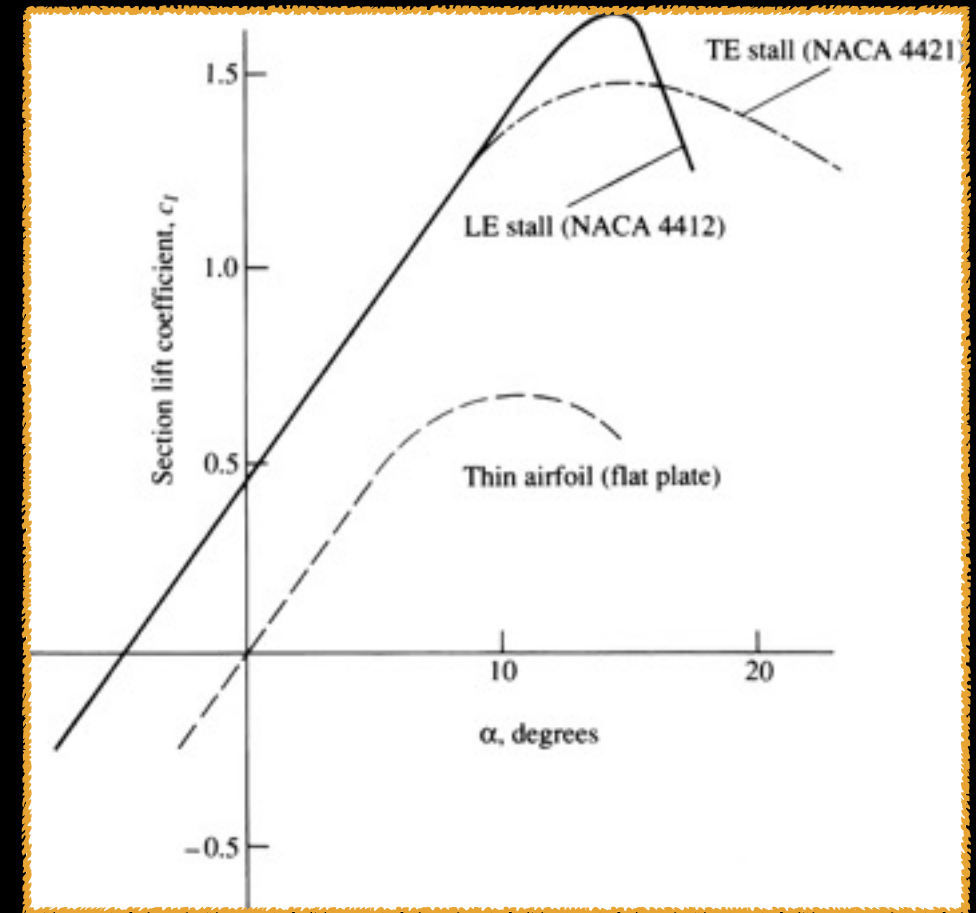
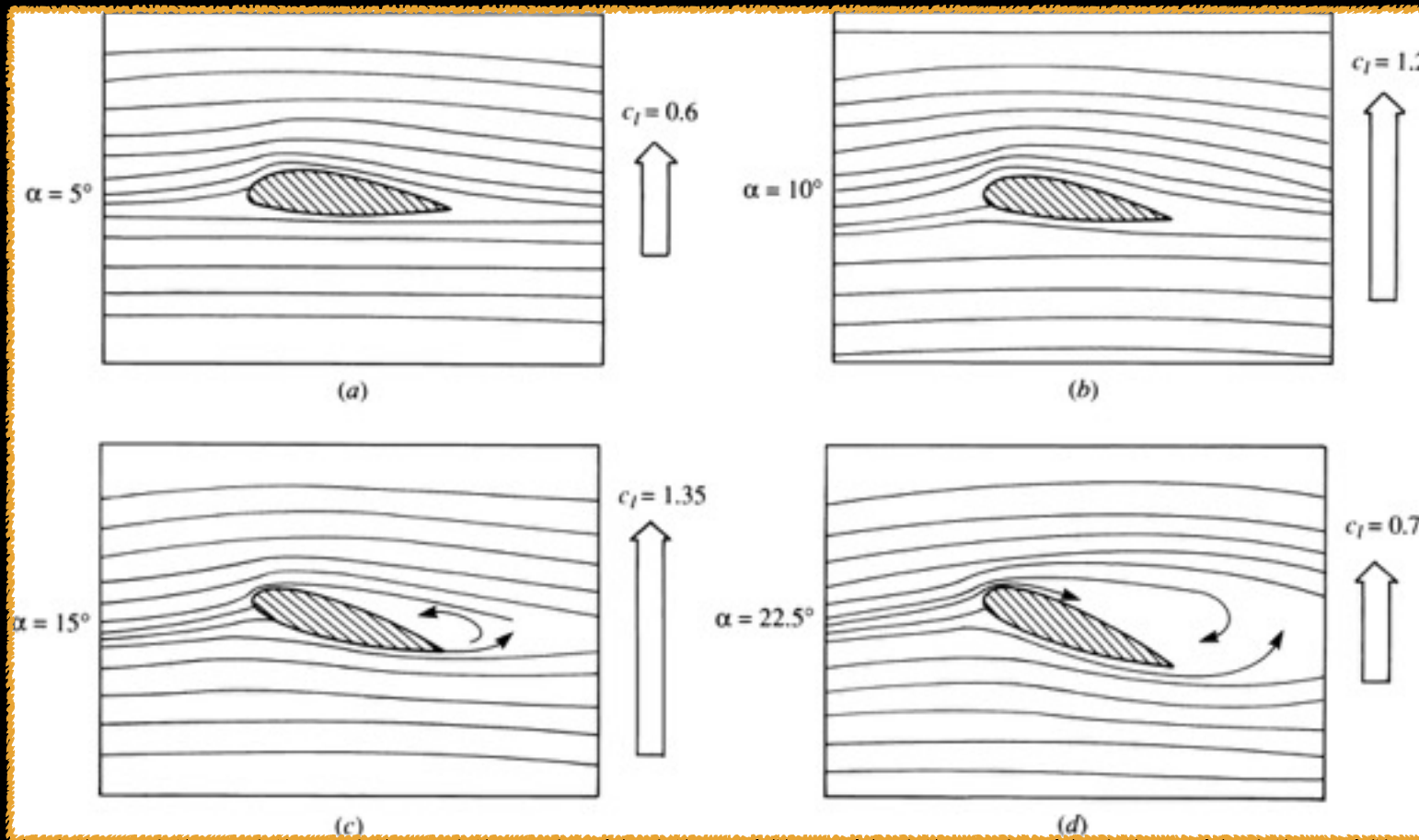
Flow Over An Airfoil - The Real Case

- In the real case, *flow separation* occurs over the top surface of the airfoil when the angle of attack exceeds the stall angle.



- Leading-Edge stall*
 - Characteristic of relatively thin airfoils.
 - Thickness-to-chord ratios usually between 10% - 16% of the chord length.

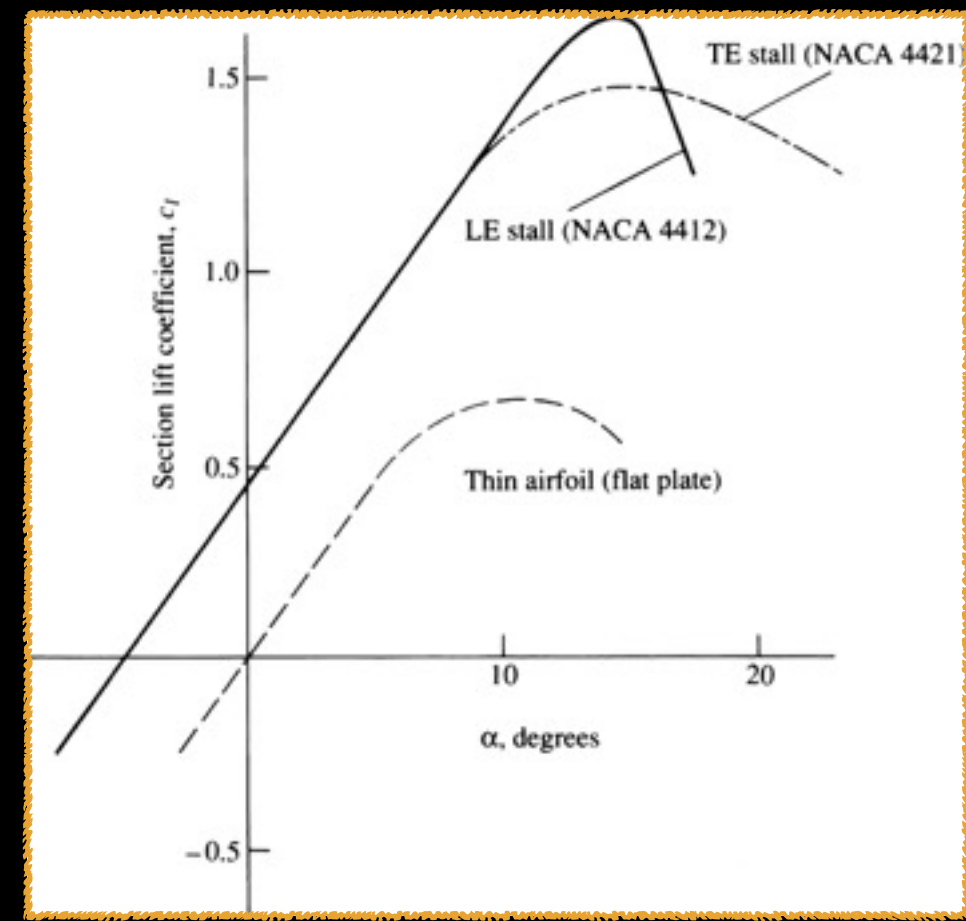
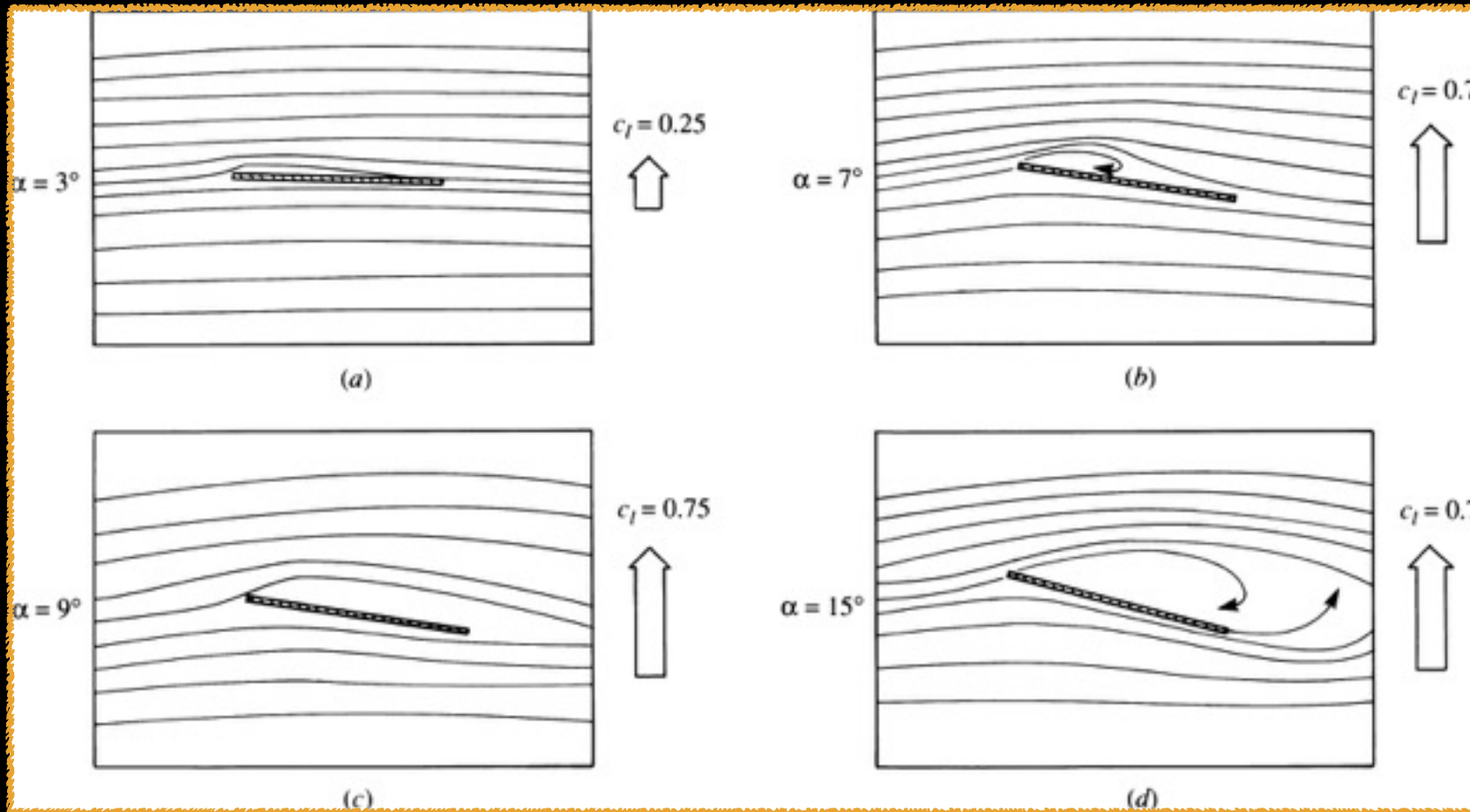
Flow Over An Airfoil - The Real Case



Trailing-Edge stall

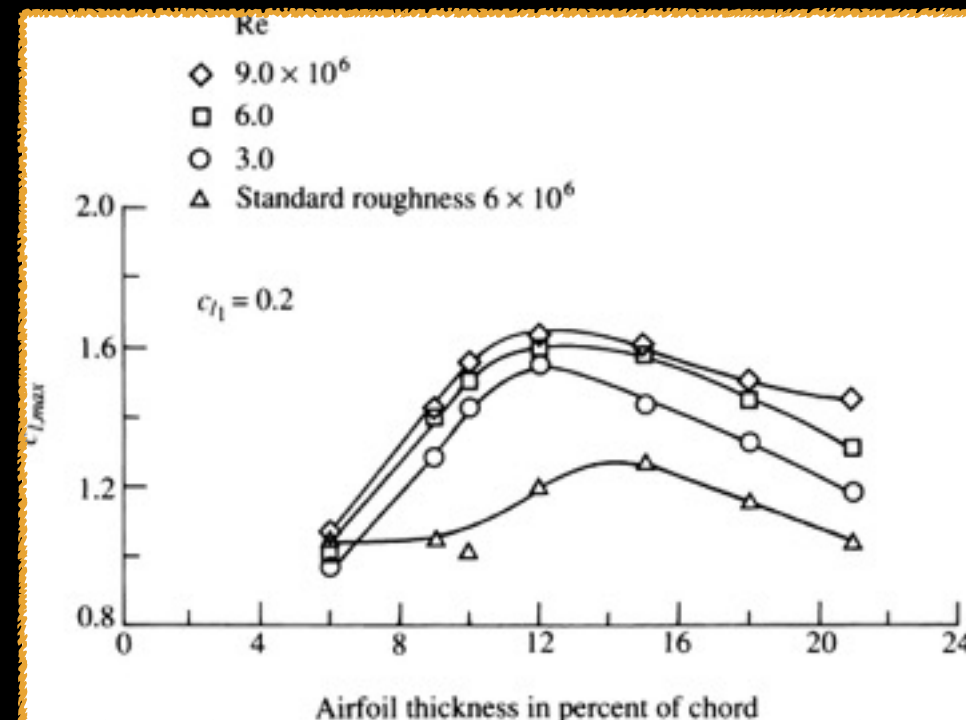
- ✱ *Characteristic of thicker airfoils.*
- ✱ *Progressive and gradual movement of separation from TE to LE as angle of attack is increased.*

Flow Over An Airfoil - The Real Case



Thin-Airfoil Stall (flat plate)

- ✱ This type of stall is associated with the extreme thinness of the airfoil.
- ✱ The thickness is about 2% of the chord length.



Other Airfoil Aerodynamics

Two figures of merit that are primarily used to judge the quality of a given airfoil are:

- * L/D ratio

- * Maximum lift coefficient.

$$V_{stall} = \sqrt{\frac{2W}{\rho_{\infty} S C_{L,max}}}$$

Tremendous incentive exists to increase the maximum lift coefficient.

- * Lower stalling speeds or higher payload capacity.

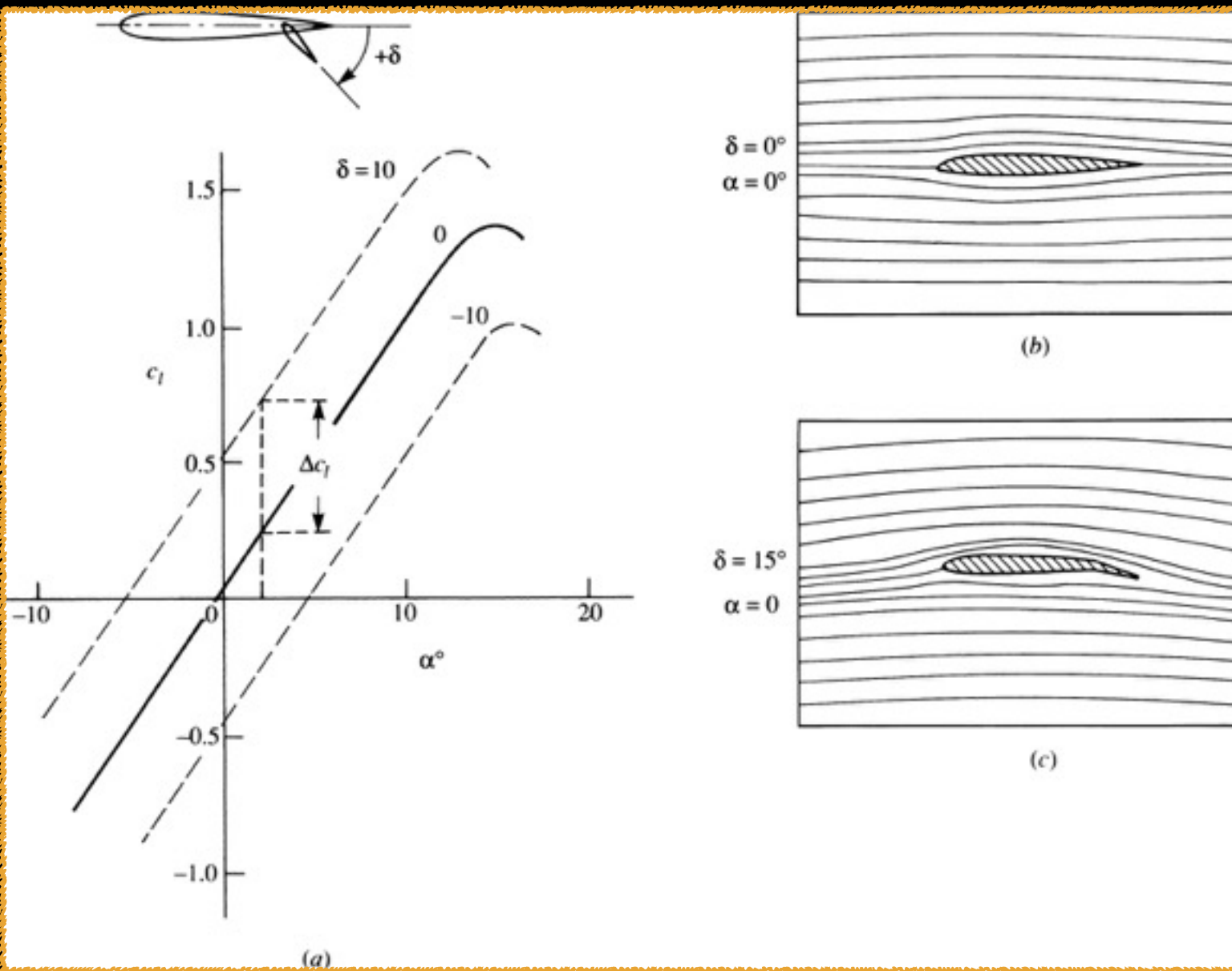
- * Maneuverability of an airplane depends on high value of $C_{L,max}$.

For an airfoil at a given Re , $C_{L,max}$ depends primarily on its shape.

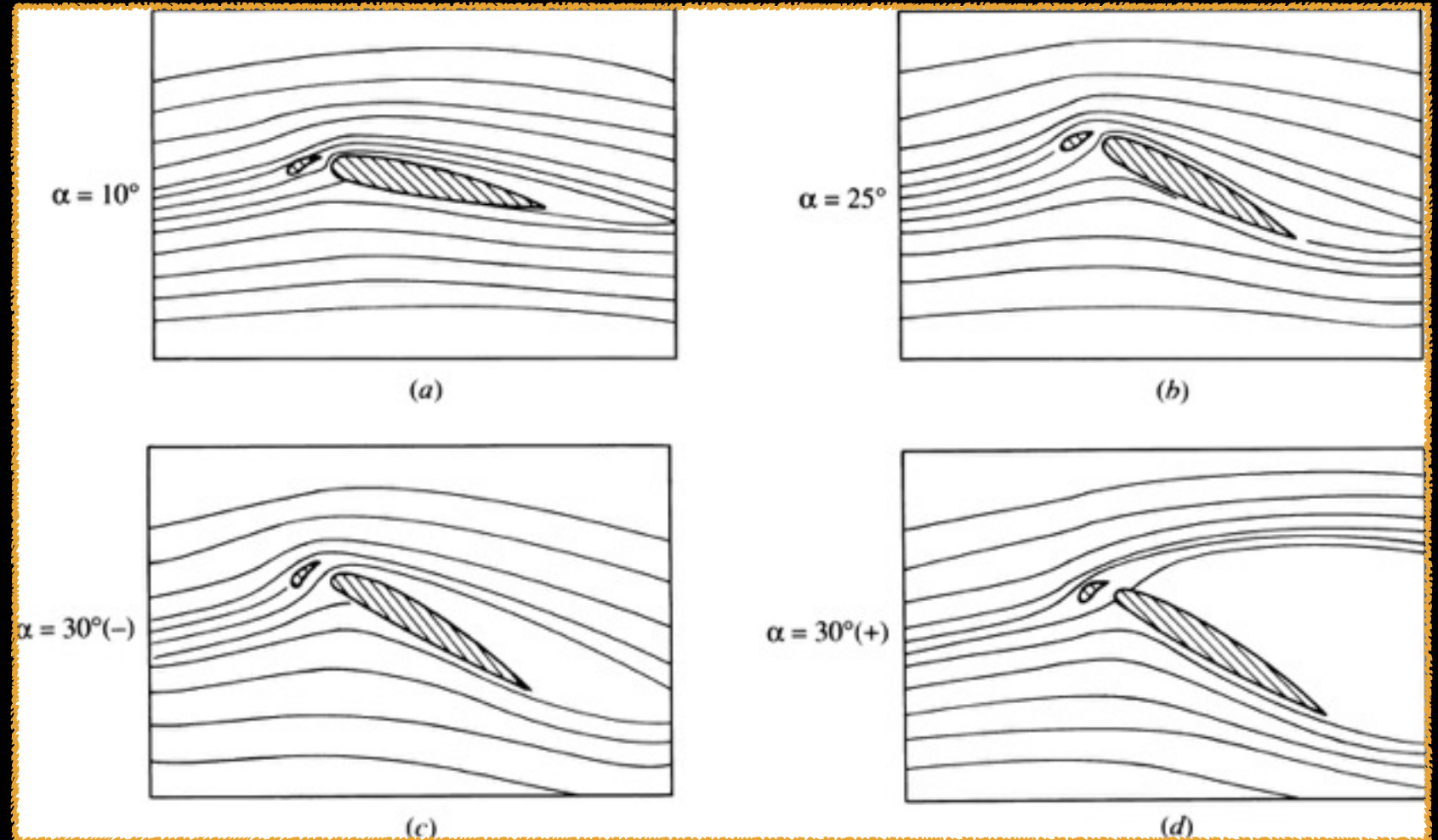
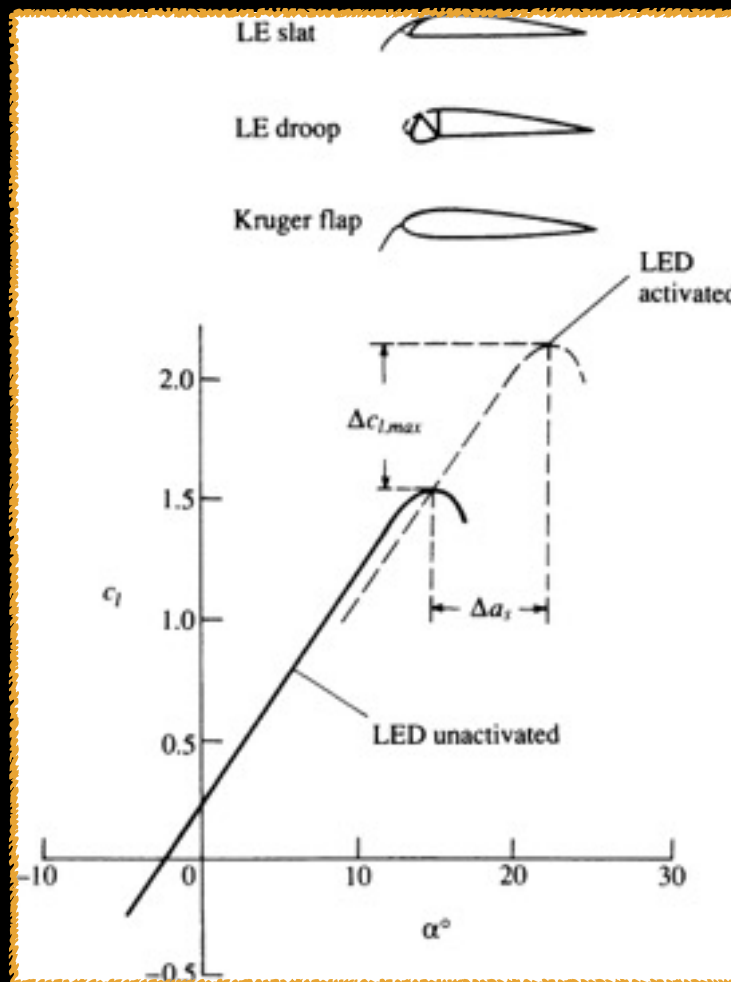
- * To increase $C_{L,max}$ further, special measures have to be carried out.

- * Measures include use of flaps, and/or LE slats - high lift devices.

High Lift Devices - TE Flaps



High Lift Devices - LE Slats

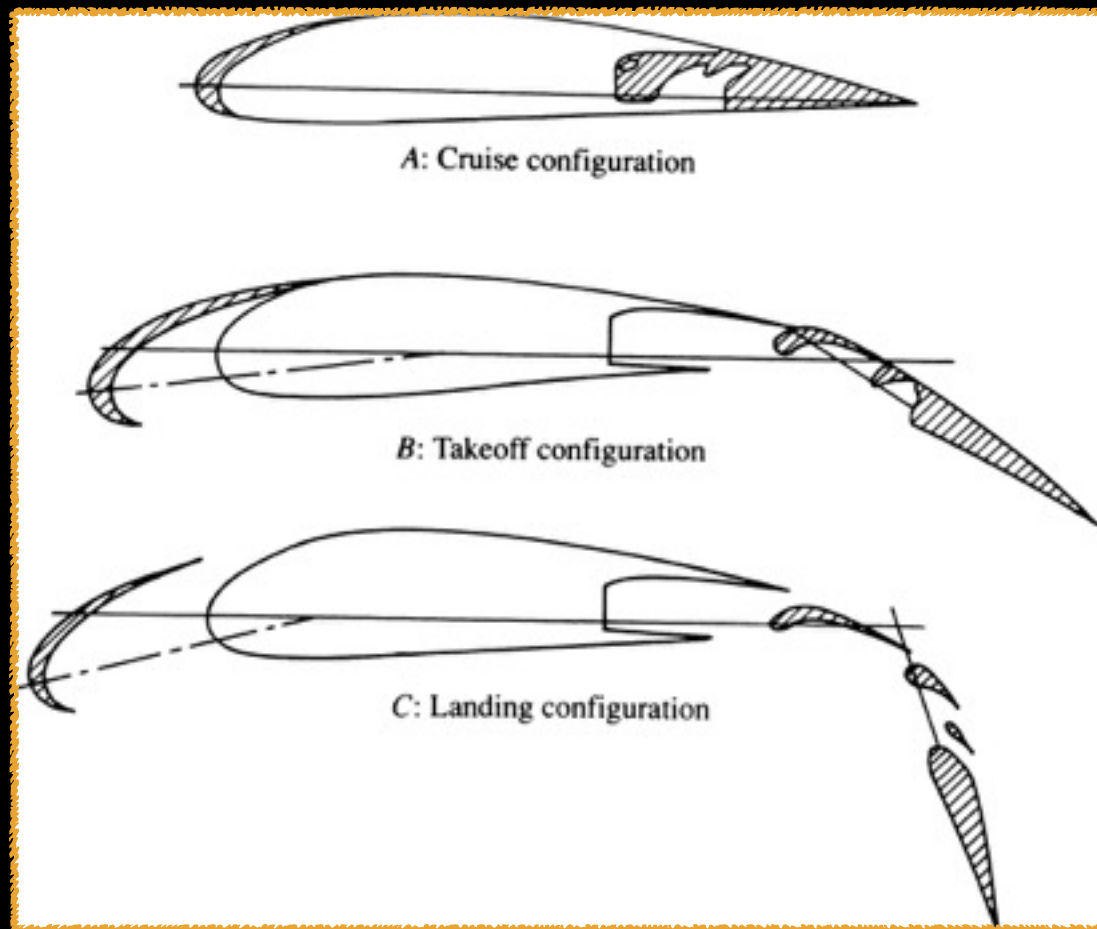


The adverse pressure gradient on the top surface is mitigated delaying flow separation.

✳ Stall angle and maximum lift coefficient increased.

There is no change in the zero-lift angle, but the lift curve is extended to a higher stalling angle of attack.

High Lift Devices



- *The main flow over the top surface of the airfoil is essentially separated.*
- *The local flow through the gaps in the multi element flap is locally attached to the top surface of the flap.*
- *Because of this locally attached flow, the lift coefficient is still quit high, around 4.5.*