2-D Turbulence for Forcing in all Scales

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- 1. Dissipation cut-off estimates via finite time averages
- 2. Cascades of pseudo fluxes
- 3. Dissipation law
- 4. Inertial range

Navier-Stokes eqns w/ per BCs on
$$\Omega = [0, L]^2$$

$$\begin{split} \frac{du}{dt} + \nu Au + B(u, u) &= f(t), \quad u = u(t) \in H \ , \ t \geq t_0 \ , \ u(t_0) = u_0 \ . \\ A &= -\mathcal{P}\Delta, \quad B(u, v) = \mathcal{P}\left((u \cdot \nabla)v\right), \quad \mathcal{P} = \text{Helmholtz-Leray proj.} \\ u &= \sum_{k \in \mathbb{Z}^2} \hat{u}_k(t) e^{i\kappa_0 k \cdot x} \ , \qquad u_{\kappa,\kappa'} = \sum_{\kappa \leq \kappa_0 |k| < \kappa'} \hat{u}_k e^{i\kappa_0 k \cdot x} \\ \kappa_0 &= \frac{2\pi}{L} \ , \quad \hat{u}_0 = 0 \ , \quad \hat{u}_k^* = \hat{u}_{-k} \ . \\ p_\kappa = u_{\kappa_0,\kappa} \ , \qquad q_\kappa = u_{\kappa,\infty} \ , \end{split}$$

Pseudo-flux of Enstrophy

(NSE,
$$Aq_k$$
), $|\cdot| = |\cdot|_{L^2}$, $||\cdot|| = |A^{1/2} \cdot |\dots$
$$\frac{1}{2} \frac{d}{dt} ||q_{\kappa}||^2 + \nu |Aq_{\kappa}|^2 = \frac{1}{\kappa_0^2} \mathfrak{F}_{\kappa} = -(B(u, u), Aq_{\kappa}) + (f, Aq_{\kappa})$$

net rate of exchange of enstrophy from low to high modes

$$\mathfrak{F}_{\kappa} = \mathfrak{E}_{\kappa}^{\to} - \mathfrak{E}_{\kappa}^{\leftarrow} + \kappa_0^2(f, Aq_{\kappa})$$

$$\mathfrak{E}_{\kappa}^{\rightarrow}(u) = -\kappa_0^2(B(p_{\kappa}, p_{\kappa}), Aq_{\kappa}) \qquad \mathfrak{E}_{\kappa}^{\leftarrow}(u) = -\kappa_0^2(B(q_{\kappa}, q_{\kappa}), Ap_{\kappa})$$

enstrophy cascade: $\langle \mathfrak{F}_{\kappa} \rangle \approx \text{ const for } \underline{\kappa}_i \leq \kappa \leq \overline{\kappa}_i$

Original dissipation cut-off estimate for time indep. force

$$\eta = \nu \kappa_0^2 \langle |Au|^2 \rangle_{\infty} , \qquad \kappa_\eta = \left(\frac{\eta}{\nu^3}\right)^{1/6} , \quad G = \frac{|f|}{\nu^2 \kappa_0^2} ,$$

 $\langle \cdot \rangle_{\infty}$ = "infinite time" average (via generalized, H-B limit) = ensemble average over global attractor

Theorem. [Foias-Manley-Temam '93]

$$G^{1/6} \lesssim \frac{\kappa_{\eta}}{\kappa_0} \leq G^{1/3}$$
 (both acheived)

Framework for time dep force

$$\overline{|f|} = \sup_{t \in \mathbb{R}} |f(t)| < \infty \qquad \overline{G} = \frac{|f|}{\nu^2 \kappa_0^2}$$

finite time ave
$$\langle \cdot \rangle = \langle \cdot \rangle_{t_1}^{t_2} = \frac{1}{t_2 - t_1} \int_{t_1}^{t_2} \cdot dt$$

$$t_2 - t_1 \ge T_{\min} \quad \overline{G} \ge \overline{G}_{\min}$$

$$f \in L^{\infty}(t_0, \infty; D(A^{j/2}) \text{ and either } \begin{cases} \dot{f} \in L^{\infty}(t_0, \infty; H) \\ |\langle f \rangle| > 0 \end{cases}$$

Two forms of diss cut-off estimates:

$$\left(\frac{\overline{G}}{C_*}\right)^{1/6} \le \frac{\kappa_{\eta}}{\kappa_0} \le (C^*\overline{G})^{1/3}$$

$$\left[\frac{\kappa_{\tau}\kappa_{\sigma}}{\kappa_{0}^{2}}\right]^{1/3} \left(\frac{\overline{G}}{K_{*}}\right)^{1/6} \leq \frac{\kappa_{\eta}}{\kappa_{0}} \leq \left[\frac{\kappa_{0}^{2}}{\kappa_{\tau}\kappa_{\sigma}}\right]^{1/3} (C^{*}\overline{G})^{1/3}$$

where

$$\kappa_{\sigma} = \left(\frac{\langle |Au|^2 \rangle}{\langle ||u||^2 \rangle}\right)^{1/2} \qquad \kappa_{\tau} = \left(\frac{\langle ||u||^2 \rangle}{\langle |u|^2 \rangle}\right)^{1/2}$$

Ave enstrophy lower bound:

$$\langle \|u\|^2 \rangle \ge C(\nu\kappa_0)^2 \overline{G}$$

As in [Foias-J.-Manley-Rosa '02] for $\mathfrak{E}_{\kappa} = \mathfrak{E}_{\kappa}^{\rightarrow} - \mathfrak{E}_{\kappa}^{\leftarrow}$ **Theorem.** Under conditions on f such that

$$\left(\frac{\overline{G}}{C_*}\right)^{1/6} \le \frac{\kappa_\eta}{\kappa_0} ,$$

we have for
$$t_2 - t_1 \ge \frac{2C_*\overline{G}}{\delta\nu\kappa_0^2}$$
 and any $\delta > 0$,
 $1 - \left(\frac{\kappa}{\kappa_\sigma}\right)^2 - \delta \le \frac{\langle \mathfrak{F}_\kappa \rangle}{\eta} \le 1 + \delta$.

Suff cond for direct enstr cascade: $\kappa_0 \ll \kappa_\sigma \implies$

$$\langle \mathfrak{F}_{\kappa}
angle pprox \eta$$
 for $\kappa_0 \leq \kappa \ll \kappa_{\sigma}$

Pseudo-flux of Energy

net rate of exchange of energy from low to high modes

$$\mathfrak{f}_{\kappa} = -(B(u,u),q_{\kappa}) + \kappa_0^2(f,q_{\kappa}) = \mathfrak{e}_{\kappa}^{\rightarrow} - \mathfrak{e}_{\kappa}^{\leftarrow} + \kappa_0^2(f,q_{\kappa})$$

net rate of exchange of energy from high to low modes

$$\mathfrak{g}_{\kappa} = -(B(u,u),p_{\kappa}) + \kappa_0^2(f,p_{\kappa}) = -[\mathfrak{e}_{\kappa}^{\rightarrow} - \mathfrak{e}_{\kappa}^{\leftarrow}] + \kappa_0^2(f,p_{\kappa})$$

where

$$\mathfrak{e}_{\kappa}^{\rightarrow}(u) = -\kappa_0^2(B(p_{\kappa}, p_{\kappa}), q_{\kappa}), \text{ and } \mathfrak{e}_{\kappa}^{\leftarrow}(u) = -\kappa_0^2(B(q_{\kappa}, q_{\kappa}), p_{\kappa}).$$

Theorem. Under the conditions on f such that $\langle ||u||^2 \rangle \geq C(\nu \kappa_0)^2 \overline{G}$

we have for
$$t_2 - t_1 \ge \frac{2\overline{G}}{C\delta\nu\kappa_0^2}$$
 and any $\delta > 0$,

$$1 - \left(\frac{\kappa}{\kappa_{\tau}}\right)^2 - \delta \leq \frac{\langle \mathfrak{f}_{\kappa} \rangle}{\epsilon} \leq 1 + \delta ,$$

and

$$1 - \left(\frac{\kappa_{\sigma}}{\kappa}\right)^2 - \delta \leq \frac{\langle \mathfrak{g}_{\kappa} \rangle}{\epsilon} \leq 1 + \delta \; .$$

direct cascade of pseudo energy flux: $\langle \mathfrak{f}_{\kappa} \rangle \approx \epsilon$ for $\kappa_0 \leq \kappa \ll \kappa_{\tau}$ inverse cascade of pseudo energy flux: $\langle \mathfrak{g}_{\kappa} \rangle \approx \epsilon$ for $\kappa_{\sigma} \ll \kappa$

Effect of forcing terms on the pseudo fluxes

For time-indep $f = f_{\underline{\kappa},\overline{\kappa}}$, nonlin fluxes $\mathfrak{E}_{\kappa} = \mathfrak{E}_{\kappa}^{\rightarrow} - \mathfrak{E}_{\kappa}^{\leftarrow}$, $\mathfrak{e}_{\kappa} = \mathfrak{e}_{\kappa}^{\rightarrow} - \mathfrak{e}_{\kappa}^{\leftarrow}$

satisfy
$$\langle \mathfrak{E}_{\kappa} \rangle_{\infty}, \langle \mathfrak{e}_{\kappa} \rangle_{\infty} \begin{cases} \geq 0 , & \text{if } \kappa > \overline{\kappa} \\ \leq 0 , & \text{if } \kappa \leq \underline{\kappa} \end{cases}$$

 $\text{Yet } \mathfrak{F}_{\kappa} = \mathfrak{E}_{\kappa} + (f, Aq_{\kappa}), \ \mathfrak{f}_{\kappa} = \mathfrak{e}_{\kappa} + (f, q_{\kappa}) \text{ and } \mathfrak{g}_{\kappa} = -\mathfrak{e}_{\kappa} + (f, p_{\kappa})$

satisfy $\langle \mathfrak{F}_{\kappa} \rangle > 0$ for $\kappa < \kappa_{\sigma}$ and $\langle \mathfrak{f}_{\kappa} \rangle > 0$ for $\kappa < \kappa_{\tau}$

 $\langle \mathfrak{g}_{\kappa} \rangle > 0$ for $\kappa > \kappa_{\sigma}$ For general force conclude: $|\langle \mathfrak{e}_{\kappa} \rangle| < \langle (f, p_{\kappa}) \rangle$ for $\kappa > \kappa_{\sigma}$

Rationale for forcing terms in the pseudo-fluxes

net rate of energy exchange into low modes : $\mathfrak{g}_{\kappa} = -\mathfrak{e}_{\kappa} + (f, p_{\kappa})$

injection of energy into p_{κ} from sources <u>external</u> to p_{κ} minus loss of energy from p_{κ} to q_{κ}

$$\frac{\mathfrak{e}_{\kappa}}{\kappa_0^2} = -(B(u,u),q_{\kappa}) = (B(u,u),p_{\kappa})$$
$$= (B(q_{\kappa},q_{\kappa}),p_{\kappa}) + (B(p_{\kappa},q_{\kappa}),p_{\kappa}) + \underbrace{(B(q_{\kappa},p_{\kappa}),p_{\kappa})}_{0} + \underbrace{(B(p_{\kappa},p_{\kappa}),p_{\kappa})}_{0}$$
$$= (B(q_{\kappa},q_{\kappa}),p_{\kappa}) - (B(p_{\kappa},p_{\kappa}),q_{\kappa})$$

Rationale for forcing terms in the psedo-fluxes

$$\underbrace{\frac{d}{dt}|p_{\kappa}|^{2}}_{\text{net change}} + \underbrace{\nu \|p_{\kappa}\|^{2}}_{\text{internal source}} = \underbrace{-(B(u,u)), p_{\kappa}) + (f, p_{\kappa})}_{\text{external source}} = \frac{\mathfrak{g}_{\kappa}}{\kappa_{0}^{2}}$$

- viscous term has same sign for all κ
- f may have mixing effect similar to nonlinear term

2-D dissipation law $\eta \sim \kappa_0^3 U^3$ where $U = \kappa_0 \langle |u|^2 \rangle^{1/2}$

Theorem. [Foias-J.-Manley-Rosa '02] $f = f_{\underline{\kappa},\overline{\kappa}}$, time indep

$$\eta \lesssim \left(\frac{\overline{\kappa}}{\kappa_0}\right)^4 \kappa_0^3 U^3 , \quad \text{where} \quad \langle \cdot \rangle = \langle \cdot \rangle_{\infty}$$

Theorem. *If*

$$\left[\frac{\kappa_{\tau}\kappa_{\sigma}}{\kappa_{0}^{2}}\right]^{1/3} \left(\frac{\overline{G}}{K_{*}}\right)^{1/6} \leq \frac{\kappa_{\eta}}{\kappa_{0}} \leq \left[\frac{\kappa_{0}^{2}}{\kappa_{\tau}\kappa_{\sigma}}\right]^{1/3} (C^{*}\overline{G})^{1/3} ,$$

then

$$\eta \leq C^* K_* \kappa_0^3 U^3$$
, where $\langle \cdot \rangle = \langle \cdot \rangle_{t_1}^{t_2}$

Proof:

$$\begin{bmatrix} \frac{\kappa_0^2}{\kappa_\tau \kappa_\sigma} \end{bmatrix} C^* \overline{G} \ge \left(\frac{\kappa_\eta}{\kappa_0}\right)^3 = \frac{1}{\nu \kappa_0^2} \langle |Au|^2 \rangle^{1/2} = \frac{1}{\nu} \left[\frac{\kappa_\tau \kappa_\sigma}{\kappa_0^2}\right] \langle |u|^2 \rangle^{1/2}$$
$$\implies C^* \overline{G} \ge \frac{1}{\nu} \left[\frac{\kappa_\tau \kappa_\sigma}{\kappa_0^2}\right]^2 \langle |u|^2 \rangle^{1/2}$$
$$C^* \frac{\langle |Au|^2 \rangle}{\nu^2 \kappa_0^4} = C^* \left(\frac{\kappa_\eta}{\kappa_0}\right)^6 \ge \left[\frac{\kappa_\tau \kappa_\sigma}{\kappa_0^2}\right]^2 \frac{C^* \overline{G}}{K_*} \ge \frac{1}{\nu K_*} \left[\frac{\kappa_\tau \kappa_\sigma}{\kappa_0^2}\right]^4 \langle |u|^2 \rangle^{1/2}$$
$$= \frac{1}{\nu \kappa_0^8 K_*} \frac{\langle |Au|^2 \rangle^2}{\langle |u|^2 \rangle^{3/2}}$$

 $\nu \langle |Au|^2 \rangle \le C^* K_* \kappa_0^4 \langle |u|^2 \rangle^{3/2}$

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2-D dissipation law equiv to $\kappa_{\tau}\kappa_{\sigma} \sim \kappa_{0}\kappa_{\eta}$

Theorem. *If*

 $\eta \le C^* K_* \kappa_0^3 U^3 \; ,$

then

$$\kappa_{\tau}\kappa_{\sigma} \leq (C^*K_*)^{1/3}\kappa_0\kappa_\eta \; .$$

Proof: The (one-sided) dissipation law can be rewritten as

$$\nu \langle |Au|^2 \rangle \le C^* K_* \kappa_0^4 \langle |u|^2 \rangle^{3/2} = C^* K_* \kappa_0^4 \kappa_\tau^{-3} \langle ||u||^2 \rangle^{3/2} ,$$

which is equivalent to

$$\langle |Au|^2 \rangle^{2/3} \le (C^* K_*)^{2/3} \nu^{-2/3} \kappa_0^{8/3} \kappa_\tau^{-2} \langle ||u||^2 \rangle$$

which is equivalent to

$$\kappa_{\tau}^{2}\kappa_{\sigma}^{2} = \kappa_{\tau}^{2} \frac{\langle |Au|^{2} \rangle}{\langle ||u||^{2} \rangle} \leq (C^{*}K_{*})^{2/3}\kappa_{0}^{2} \left(\frac{\kappa_{0}^{2} \langle |Au|^{2} \rangle}{\nu^{2}}\right)^{1/3} = (C^{*}K_{*})^{2/3}\kappa_{0}^{2}\kappa_{\eta}^{2} .$$

Inertial range

- 1. A significant amount of enstrophy should be in the inertial range.
- 2. This range should be wide, in particular $\underline{\kappa}_i \ll \overline{\kappa}_i \sim \kappa_\eta$.
- 3. The enstrophy cascade should hold over this range.
- 4. The power law

$$\mathsf{e}_{\kappa,2\kappa} \sim \frac{\eta^{2/3}}{\kappa^2}$$

should hold for all $\kappa \in [\underline{\kappa}_i, \overline{\kappa}_i]$.

Logarithmic correction

As in [Rose-Sulem'78], [Ohkitani'89]

$$\frac{\eta^{2/3}}{\kappa^2} \prec \mathsf{e}_{\kappa,2\kappa} \prec \frac{\eta^{2/3}}{\kappa^2} \,,$$

where

 $a \prec b$ when $a \leq C(\log(s\overline{G}))^{\alpha}b$ for some $\alpha \in \mathbb{R}$, and large enough \overline{G}

and where C and s are shape factors, with a similar convention for \succ .

Theorem. For any κ such that

$$\frac{\eta^{2/3}}{\kappa^2} \prec \mathsf{e}_{\kappa,2\kappa}$$

we have $\kappa \prec \kappa_{\eta}$, and consequently

 $\overline{\kappa}_i \prec \kappa_\eta \; .$

Theorem. *If*

$$\kappa_\eta \prec \kappa_\sigma$$
,

then

$$e_{\kappa,2\kappa} \prec \frac{\eta^{2/3}}{\kappa^2}$$
, for all $\kappa \ge \kappa_0$.

Moreover, for any κ such that

$$\frac{\eta^{2/3}}{\kappa^2} \prec \mathsf{e}_{\kappa,2\kappa} \prec \frac{\eta^{2/3}}{\kappa^2} \,,$$

we have $\kappa_{\tau} \prec \kappa$, and consequently

$$\kappa_{\tau} \prec \underline{\kappa}_i , \qquad \kappa_{\tau} \prec \kappa_0$$

Theorem. *If*

$$\mathsf{e}_{\kappa,2\kappa} \prec \frac{\eta^{2/3}}{\kappa^2}$$

holds for

$$\kappa_1 = \left(1 - \frac{C}{\log \overline{G}}\right)^{-1/2} \kappa_\tau \le \kappa \le \kappa_2 = \left(1 - \frac{C}{\log \overline{G}}\right)^{1/2} \kappa_\eta$$

then for all \overline{G} large enough,

 $\kappa_\eta \prec \kappa_\sigma$

and

$$\frac{\langle \|q_{\kappa_1}\|^2 \rangle}{\langle \|p_{\kappa_1}\|^2 \rangle} \ge \frac{1}{\log \overline{G}} \ .$$

Theorem. If $t_2 - t_1 \ge T_{\min}$ and $\overline{G} \ge \overline{G}_{\min}$, then

$$\left(\frac{\overline{G}}{C_*}\right)^{1/6} \le \frac{\kappa_{\eta}}{\kappa_0} \le (C^*\overline{G})^{1/3}$$

force	T_{\min}	\overline{G}_{\min}	C_*	C^*
$f, \dot{f} \in L^{\infty}(t_0, \infty; H)$	$\frac{c}{\nu\kappa_0^2\Gamma^2}$	$c(\Gamma_1^2 + 1)$	$\frac{c}{\Gamma^2}$	С
$f \in L^{\infty}(t_0, \infty; H), \langle f \rangle > 0$	$\frac{c}{\nu\kappa_0^2\Gamma_0^2}$	$\frac{c}{\Gamma_0^2}$	$\frac{c}{\Gamma_0^2}$	С

$$\Gamma_1 = \sup\left(\frac{\langle |\dot{f}|^2 \rangle}{\nu^2 \kappa_0^4 \langle |f|^2 \rangle}\right)^{1/2} \quad \Gamma = \inf\frac{\langle |f|^2 \rangle^{1/2}}{|f|} \quad \Gamma_0 = \inf\frac{|\langle f \rangle|}{|f|}$$

inf are over all t_1, t_2 s.t. $t_2 - t_1 \ge 1/(\nu \kappa_0^2)$.

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sup,

Theorem. If $f \in L^{\infty}(t_0, \infty; X)$, $t_2 - t_1 \ge T_{\min}$ and $\overline{G} \ge \overline{G}_{\min}$, then

$$\left[\frac{\kappa_{\tau}\kappa_{\sigma}}{\kappa_{0}^{2}}\right]^{1/3} \left(\frac{\overline{G}}{K_{*}}\right)^{1/6} \leq \frac{\kappa_{\eta}}{\kappa_{0}} \leq \left[\frac{\kappa_{0}^{2}}{\kappa_{\tau}\kappa_{\sigma}}\right]^{1/3} (C^{*}\overline{G})^{1/3}$$

force	T_{\min}	\overline{G}_{\min}	K_*	C^*
$X = D(A), \ \dot{f} \in L^{\infty}(t_0, \infty; H)$	$\frac{\overline{G}^{1/2}\log^{1/4}\overline{G}}{\nu\kappa_0^2}$	$\frac{c(\psi_2^2 + \Gamma_1^2)}{\varphi_2 \log^{1/2} \overline{G}}$	$\frac{c\psi_2\log^{1/2}\overline{G}}{\Gamma}$	$c\psi_2$
$X = D(A^{3/2}), \ \langle f \rangle > 0$	$\frac{c\overline{G}^{1/2}}{\nu\kappa_0^2}$	$\frac{c\theta_2}{\Gamma_0 \left[\log\frac{e\theta_3}{\theta_2}\right]^{1/2}}$	$\frac{c\theta_2 \left[\log \frac{e\theta_3}{\theta_2}\right]^{1/2}}{\Gamma_0}$	$c\psi_2$
$(\lambda i/2 c 2)$	$\frac{1}{ \lambda_i/2 }$	Aj	/2/c	

$$\psi_j^2 = \sup \frac{\langle |A^{j/2}f|^2 \rangle}{\kappa_0^{2j} \langle |f|^2 \rangle} \qquad \varphi_j = \frac{|A^{j/2}f|}{\kappa_0^{2j}\overline{|f|}} \qquad \theta_j = \sup \frac{|A^{j/2} \langle f \rangle|}{\kappa_0^j |\langle f \rangle|}$$

Sharpening in diss cut-off est that results if $\kappa_{\tau}\kappa_{\sigma} \sim \kappa_{0}\kappa_{\eta}$

Corollary. If $t_2 - t_1 \ge T_{\min}$, $\overline{G} \ge \overline{G}_{\min}$, and

$$rac{\kappa_{ au}\kappa_{\sigma}}{\kappa_0^2}\sim rac{\kappa_{\eta}}{\kappa_0} \ ,$$

then

$$\left(\frac{\overline{G}}{K_*}\right)^{1/4} \le \frac{\kappa_{\eta}}{\kappa_0} \le (C^*\overline{G})^{1/4} .$$

Summary

•
$$\left(\frac{\kappa_{\tau}\kappa_{\sigma}}{\kappa_{0}^{2}}\right)^{1/3}\overline{G}^{1/6} \lesssim \kappa_{\eta}/\kappa_{0} \lesssim \left(\frac{\kappa_{0}^{2}}{\kappa_{\tau}\kappa_{\sigma}}\right)^{1/3}\overline{G}^{1/3}$$

- Tme dep forcing in all scales. Shape factors fixed as $\overline{G} \to \infty$
- Lower bound on $\kappa_{\eta}/\kappa_0 \implies$ direct enstrophy cascade if $\kappa_{\sigma} \gg \kappa_0$
- Energy power law $\implies \kappa_{\sigma} \sim \kappa_{\eta}$ (up to log)
- 2-D diss law $\eta \sim \kappa_0^3 U^3$ equiv to $\kappa_\tau \kappa_\sigma \sim \kappa_0 \kappa_\eta$

•
$$\kappa_{\tau}\kappa_{\sigma} \sim \kappa_{0}\kappa_{\eta} \implies$$
 sharp bound $\overline{G}^{1/4} \sim \kappa_{\eta}/\kappa_{0}$