

# 2D Navier-Stokes equations with singular forcing

Peter Constantin

Department of Mathematics  
The University of Chicago

Texas, October 2009

2D  
Navier-Stokes  
equations with  
singular  
forcing

Peter  
Constantin

Introduction

Kinetics:  
NLFP

Embedding in  
Fluid

Results on  
NS+NLFP

Navier-Stokes

**Collaborator: G. Seregin, Oxford**

**Support: NSF**

## In this talk:

- Motivation: Nonlinear Fokker-Planck Equation

## In this talk:

- Motivation: Nonlinear Fokker-Planck Equation
- Navier-Stokes Results

- Configuration space:  $M =$  compact, separable, metric space.  $m \in M =$  corpus = doodad.

- Configuration space:  $M =$  compact, separable, metric space.  $m \in M =$  corpus = doodad.
- Reference measure:  $d\mu -$  Borel Probability on  $M$ .

- Configuration space:  $M =$  compact, separable, metric space.  $m \in M =$  corpus = doodad.
- Reference measure:  $d\mu$  – Borel Probability on  $M$ .
- Corpora state  $f(m)d\mu(m)$  – Probabililty, AC w.r.  $d\mu$ .

- Configuration space:  $M =$  compact, separable, metric space.  $m \in M =$  corpus = doodad.
- Reference measure:  $d\mu$  – Borel Probability on  $M$ .
- Corpora state  $f(m)d\mu(m)$  – Probabililty, AC w.r.  $d\mu$ .
- Interaction kernel  $k : M \times M \rightarrow \mathbb{R}_+$ ,



- Configuration space:  $M =$  compact, separable, metric space.  $m \in M =$  corpus = doodad.
- Reference measure:  $d\mu$  – Borel Probability on  $M$ .
- Corpora state  $f(m)d\mu(m)$  – Probabililty, AC w.r.  $d\mu$ .
- Interaction kernel  $k : M \times M \rightarrow \mathbb{R}_+$ ,
- symmetric:  $k(m, p) = k(p, m)$

- Configuration space:  $M =$  compact, separable, metric space.  $m \in M =$  corpus = doodad.
- Reference measure:  $d\mu$  – Borel Probability on  $M$ .
- Corpora state  $f(m)d\mu(m)$  – Probabililty, AC w.r.  $d\mu$ .
- Interaction kernel  $k : M \times M \rightarrow \mathbb{R}_+$ ,
- symmetric:  $k(m, p) = k(p, m)$
- uniformly bi-Lipschitz:

$$|k(m, n) - k(p, n)| \leq Ld(m, p)$$

- Configuration space:  $M =$  compact, separable, metric space.  $m \in M =$  corpus = doodad.
- Reference measure:  $d\mu$  – Borel Probability on  $M$ .
- Corpora state  $f(m)d\mu(m)$  – Probabililty, AC w.r.  $d\mu$ .
- Interaction kernel  $k : M \times M \rightarrow \mathbb{R}_+$ ,
- symmetric:  $k(m, p) = k(p, m)$
- uniformly bi-Lipschitz:

$$|k(m, n) - k(p, n)| \leq Ld(m, p)$$

- Potential  $U[f](m) = \int_M k(m, p)f(p)d\mu(p)$

- Configuration space:  $M =$  compact, separable, metric space.  $m \in M =$  corpus = doodad.
- Reference measure:  $d\mu$  – Borel Probability on  $M$ .
- Corpora state  $f(m)d\mu(m)$  – Probabililty, AC w.r.  $d\mu$ .
- Interaction kernel  $k : M \times M \rightarrow \mathbb{R}_+$ ,
- symmetric:  $k(m, p) = k(p, m)$
- uniformly bi-Lipschitz:

$$|k(m, n) - k(p, n)| \leq Ld(m, p)$$

- Potential  $U[f](m) = \int_M k(m, p)f(p)d\mu(p)$
- Potential  $U =$  **micro-micro interaction**

- Configuration space:  $M =$  compact, separable, metric space.  $m \in M =$  corpus = doodad.
- Reference measure:  $d\mu$  – Borel Probability on  $M$ .
- Corpora state  $f(m)d\mu(m)$  – Probabililty, AC w.r.  $d\mu$ .
- Interaction kernel  $k : M \times M \rightarrow \mathbb{R}_+$ ,
- symmetric:  $k(m, p) = k(p, m)$
- uniformly bi-Lipschitz:

$$|k(m, n) - k(p, n)| \leq Ld(m, p)$$

- Potential  $U[f](m) = \int_M k(m, p)f(p)d\mu(p)$
- Potential  $U =$  **micro-micro interaction**
- Free Energy

$$\mathcal{E}[f] = \int_M f \log fd\mu + \frac{1}{2} \int_M U[f]fd\mu$$

- Configuration space:  $M =$  compact, separable, metric space.  $m \in M =$  corpus = doodad.
- Reference measure:  $d\mu$  – Borel Probability on  $M$ .
- Corpora state  $f(m)d\mu(m)$  – Probabililty, AC w.r.  $d\mu$ .
- Interaction kernel  $k : M \times M \rightarrow \mathbb{R}_+$ ,
- symmetric:  $k(m, p) = k(p, m)$
- uniformly bi-Lipschitz:

$$|k(m, n) - k(p, n)| \leq Ld(m, p)$$

- Potential  $U[f](m) = \int_M k(m, p)f(p)d\mu(p)$
- Potential  $U =$  **micro-micro interaction**
- Free Energy

$$\mathcal{E}[f] = \int_M f \log f d\mu + \frac{1}{2} \int_M U[f] f d\mu$$

- Minima of Free Energy: Onsager Equation

$$f = Z^{-1} e^{-U[f]}.$$

## Kinetics: NLFP

$M$  compact, connected Riemannian manifold of dimension  $d$

## Kinetics: NLFP

$M$  compact, connected Riemannian manifold of dimension  $d$   
with metric  $g_{\alpha\beta}$ ,



## Kinetics: NLFP

$M$  compact, connected Riemannian manifold of dimension  $d$   
with metric  $g_{\alpha\beta}$ ,  $g = \det(g_{\alpha\beta})$ ,

## Kinetics: NLFP

$M$  compact, connected Riemannian manifold of dimension  $d$   
with metric  $g_{\alpha\beta}$ ,  $g = \det(g_{\alpha\beta})$ ,  $(g^{\alpha\beta}) = (g_{\alpha\beta})^{-1}$ ,

## Kinetics: NLFP

$M$  compact, connected Riemannian manifold of dimension  $d$  with metric  $g_{\alpha\beta}$ ,  $g = \det(g_{\alpha\beta})$ ,  $(g^{\alpha\beta}) = (g_{\alpha\beta})^{-1}$ , volume element  $d\mu = \sqrt{g}d\phi$  in local coordinates  $\phi$ .

## Kinetics: NLFP

$M$  compact, connected Riemannian manifold of dimension  $d$  with metric  $g_{\alpha\beta}$ ,  $g = \det(g_{\alpha\beta})$ ,  $(g^{\alpha\beta}) = (g_{\alpha\beta})^{-1}$ , volume element  $d\mu = \sqrt{g}d\phi$  in local coordinates  $\phi$ . Generalized Doi-Smoluchowski equation

## Kinetics: NLFP

$M$  compact, connected Riemannian manifold of dimension  $d$  with metric  $g_{\alpha\beta}$ ,  $g = \det(g_{\alpha\beta})$ ,  $(g^{\alpha\beta}) = (g_{\alpha\beta})^{-1}$ , volume element  $d\mu = \sqrt{g}d\phi$  in local coordinates  $\phi$ . Generalized Doi-Smoluchowski equation

$$\partial_t f = \Delta_g f + \operatorname{div}_g(f \nabla_g U)$$

## Kinetics: NLFP

$M$  compact, connected Riemannian manifold of dimension  $d$  with metric  $g_{\alpha\beta}$ ,  $g = \det(g_{\alpha\beta})$ ,  $(g^{\alpha\beta}) = (g_{\alpha\beta})^{-1}$ , volume element  $d\mu = \sqrt{g}d\phi$  in local coordinates  $\phi$ . Generalized Doi-Smoluchowski equation

$$\partial_t f = \Delta_g f + \operatorname{div}_g(f \nabla_g U)$$

$\Delta_g$  Laplace-Beltrami,

## Kinetics: NLFP

$M$  compact, connected Riemannian manifold of dimension  $d$  with metric  $g_{\alpha\beta}$ ,  $g = \det(g_{\alpha\beta})$ ,  $(g^{\alpha\beta}) = (g_{\alpha\beta})^{-1}$ , volume element  $d\mu = \sqrt{g}d\phi$  in local coordinates  $\phi$ . Generalized Doi-Smoluchowski equation

$$\partial_t f = \Delta_g f + \operatorname{div}_g(f \nabla_g U)$$

$\Delta_g$  Laplace-Beltrami,

$$\operatorname{div}_g(f \nabla_g U) = \frac{1}{\sqrt{g}} \partial_\alpha \left( \sqrt{g} g^{\alpha\beta} f \partial_\beta U \right).$$

## Kinetics: NLFP

$M$  compact, connected Riemannian manifold of dimension  $d$  with metric  $g_{\alpha\beta}$ ,  $g = \det(g_{\alpha\beta})$ ,  $(g^{\alpha\beta}) = (g_{\alpha\beta})^{-1}$ , volume element  $d\mu = \sqrt{g}d\phi$  in local coordinates  $\phi$ . Generalized Doi-Smoluchowski equation

$$\partial_t f = \Delta_g f + \operatorname{div}_g(f \nabla_g U)$$

$\Delta_g$  Laplace-Beltrami,

$$\operatorname{div}_g(f \nabla_g U) = \frac{1}{\sqrt{g}} \partial_\alpha \left( \sqrt{g} g^{\alpha\beta} f \partial_\beta U \right).$$

$$U[f](p) = \int_M k(p, q) f(q) d\mu(q)$$



# Dissipative Structure

Free Energy:

2D

Navier-Stokes  
equations with  
singular  
forcing

Peter  
Constantin

Introduction

Kinetics:  
NLFP

Embedding in  
Fluid

Results on  
NS+NLFP

Navier-Stokes

# Dissipative Structure

Free Energy:

$$\mathcal{E}[f] = \int_M \left\{ \log f + \frac{1}{2} U[f] \right\} f d\mu$$

2D

Navier-Stokes  
equations with  
singular  
forcing

Peter  
Constantin

Introduction

Kinetics:  
NLFP

Embedding in  
Fluid

Results on  
NS+NLFP

Navier-Stokes

# Dissipative Structure

Free Energy:

$$\mathcal{E}[f] = \int_M \left\{ \log f + \frac{1}{2} U[f] \right\} f d\mu$$

NLFP:

2D  
Navier-Stokes  
equations with  
singular  
forcing

Peter  
Constantin

Introduction

Kinetics:  
NLFP

Embedding in  
Fluid

Results on  
NS+NLFP

Navier-Stokes

# Dissipative Structure

Free Energy:

$$\mathcal{E}[f] = \int_M \left\{ \log f + \frac{1}{2} U[f] \right\} f d\mu$$

NLFP:

$$\partial_t f = \operatorname{div}_g \cdot \left( f \nabla_g \left( \frac{\delta \mathcal{E}[f]}{\delta f} \right) \right)$$

2D

Navier-Stokes  
equations with  
singular  
forcing

Peter  
Constantin

Introduction

Kinetics:  
NLFP

Embedding in  
Fluid

Results on  
NS+NLFP

Navier-Stokes

# Dissipative Structure

Free Energy:

$$\mathcal{E}[f] = \int_M \left\{ \log f + \frac{1}{2} U[f] \right\} f d\mu$$

NLFP:

$$\partial_t f = \operatorname{div}_g \cdot \left( f \nabla_g \left( \frac{\delta \mathcal{E}[f]}{\delta f} \right) \right)$$

Lyapunov functional:

$$\frac{d}{dt} \mathcal{E}[f] = - \int_M f \left| \nabla_g \left( \frac{\delta \mathcal{E}[f]}{\delta f} \right) \right|^2 d\mu$$

# Dissipative Structure

Free Energy:

$$\mathcal{E}[f] = \int_M \left\{ \log f + \frac{1}{2} U[f] \right\} f d\mu$$

NLFP:

$$\partial_t f = \operatorname{div}_g \cdot \left( f \nabla_g \left( \frac{\delta \mathcal{E}[f]}{\delta f} \right) \right)$$

Lyapunov functional:

$$\frac{d}{dt} \mathcal{E}[f] = - \int_M f \left| \nabla_g \left( \frac{\delta \mathcal{E}[f]}{\delta f} \right) \right|^2 d\mu$$

Dynamics: nontrivial. Multiple steady states, gradient system, finite dimensional attractor.

# Dissipative Structure

2D  
Navier-Stokes  
equations with  
singular  
forcing

Free Energy:

$$\mathcal{E}[f] = \int_M \left\{ \log f + \frac{1}{2} U[f] \right\} f d\mu$$

NLFP:

$$\partial_t f = \operatorname{div}_g \cdot \left( f \nabla_g \left( \frac{\delta \mathcal{E}[f]}{\delta f} \right) \right)$$

Lyapunov functional:

$$\frac{d}{dt} \mathcal{E}[f] = - \int_M f \left| \nabla_g \left( \frac{\delta \mathcal{E}[f]}{\delta f} \right) \right|^2 d\mu$$

Dynamics: nontrivial. Multiple steady states, gradient system, finite dimensional attractor. Inertial Manifolds: Vukadinovic (2008-9).

Peter  
Constantin

Introduction

Kinetics:  
NLFP

Embedding in  
Fluid

Results on  
NS+NLFP

Navier-Stokes

## Embedding in Fluid: Passive

$$\partial_t f + u \cdot \nabla_x f + \operatorname{div}_g(Wf) = \operatorname{div}_g(f \nabla_g(\log f + U[f]))$$



## Embedding in Fluid: Passive

$$\partial_t f + \mathbf{u} \cdot \nabla_x f + \operatorname{div}_g(Wf) = \operatorname{div}_g(f \nabla_g (\log f + U[f]))$$

with

$$W(x, m, t) = \sum_{i,j=1}^n c_{ji}^{\alpha}(m) \frac{\partial u_i}{\partial x^j}(x, t)$$

## Embedding in Fluid: Passive

$$\partial_t f + u \cdot \nabla_x f + \operatorname{div}_g(Wf) = \operatorname{div}_g(f \nabla_g (\log f + U[f]))$$

with

$$W(x, m, t) = \sum_{i,j=1}^n c_{ji}^\alpha(m) \frac{\partial u_i}{\partial x^j}(x, t)$$

$$c_{ji}^\alpha(m) \in T_m^*(M).$$

## Embedding in Fluid: Passive

$$\partial_t f + \mathbf{u} \cdot \nabla_x f + \operatorname{div}_g(Wf) = \operatorname{div}_g(f \nabla_g(\log f + U[f]))$$

with

$$W(x, m, t) = \sum_{i,j=1}^n c_{ji}^{\alpha}(m) \frac{\partial u_i}{\partial x^j}(x, t)$$

$$c_{ji}^{\alpha}(m) \in T_m^*(M).$$

Example, rods in 3D:

## Embedding in Fluid: Passive

$$\partial_t f + \mathbf{u} \cdot \nabla_x f + \operatorname{div}_g(Wf) = \operatorname{div}_g(f \nabla_g(\log f + U[f]))$$

with

$$W(x, m, t) = \sum_{i,j=1}^n c_{ji}^{\alpha}(m) \frac{\partial u_i}{\partial x^j}(x, t)$$

$$c_{ji}^{\alpha}(m) \in T_m^*(M).$$

Example, rods in 3D:  $M = \mathbb{S}^2$ ,

## Embedding in Fluid: Passive

$$\partial_t f + u \cdot \nabla_x f + \operatorname{div}_g(Wf) = \operatorname{div}_g(f \nabla_g (\log f + U[f]))$$

with

$$W(x, m, t) = \sum_{i,j=1}^n c_{ji}^{\alpha}(m) \frac{\partial u_i}{\partial x^j}(x, t)$$

$$c_{ji}(m) \in T_m^*(M).$$

Example, rods in 3D:  $M = \mathbb{S}^2$ ,

$$W(x, m, t) = (\nabla_x u(x, t))m - ((\nabla_x u(x, t))m \cdot m)m.$$

## Embedding in Fluid: Passive

$$\partial_t f + u \cdot \nabla_x f + \operatorname{div}_g(Wf) = \operatorname{div}_g(f \nabla_g(\log f + U[f]))$$

with

$$W(x, m, t) = \sum_{i,j=1}^n c_{ji}^\alpha(m) \frac{\partial u_i}{\partial x^j}(x, t)$$

$$c_{ji}^\alpha(m) \in T_m^*(M).$$

Example, rods in 3D:  $M = \mathbb{S}^2$ ,

$$W(x, m, t) = (\nabla_x u(x, t))m - ((\nabla_x u(x, t))m \cdot m)m.$$

Macro-Micro Effect: from first principles, if scales are separated.

## Active: Navier-Stokes

$$\begin{aligned}\partial_t u + u \cdot \nabla u + \nabla p &= \nu \Delta u + \nabla \cdot \sigma \\ \nabla \cdot u &= 0\end{aligned}$$

2D

Navier-Stokes  
equations with  
singular  
forcing

Peter  
Constantin

Introduction

Kinetics:  
NLFP

Embedding in  
Fluid

Results on  
NS+NLFP

Navier-Stokes

## Active: Navier-Stokes

$$\partial_t u + u \cdot \nabla u + \nabla p = \nu \Delta u + \nabla \cdot \sigma$$

$$\nabla \cdot u = 0$$

$$\sigma = \sigma_{ij}(x, t)$$

**added stress tensor.**



## Active: Navier-Stokes

$$\partial_t u + u \cdot \nabla u + \nabla p = \nu \Delta u + \nabla \cdot \sigma$$

$$\nabla \cdot u = 0$$

$$\sigma = \sigma_{ij}(x, t)$$

**added stress tensor.**

## Micro-Macro Effect

$$\sigma(x) = \int_M \{c \cdot \nabla_g U[f](x, m) - \operatorname{div}_g c\} f(x, m) d\mu(m) \quad *$$

## Active: Navier-Stokes

$$\begin{aligned}\partial_t u + u \cdot \nabla u + \nabla p &= \nu \Delta u + \nabla \cdot \sigma \\ \nabla \cdot u &= 0\end{aligned}$$

$$\sigma = \sigma_{ij}(x, t)$$

**added stress tensor.**

## Micro-Macro Effect

$$\sigma(x) = \int_M \{c \cdot \nabla_g U[f](x, m) - \operatorname{div}_g c\} f(x, m) d\mu(m) \quad *$$

$$f \geq 0, \sup_x \int_M f d\mu(m) < \infty \Rightarrow \sigma \in L^\infty$$

## Theorem

*3DNS + NL Fokker-Planck eqns with \**. Then

$$E(t) = \frac{1}{2} \int |u|^2 dx + \\ + \int \left\{ f \log f + \frac{1}{2} (U[f]) f \right\} dx d\mu.$$

*is nondecreasing on solutions.*

## Theorem

3DNS + NL Fokker-Planck eqns with \*. Then

$$E(t) = \frac{1}{2} \int |u|^2 dx + \int \left\{ f \log f + \frac{1}{2} (U[f]) f \right\} dx d\mu.$$

is nondecreasing on solutions. If  $(u, f)$  is a smooth solution then

$$\frac{dE}{dt} = -\nu \int |\nabla_x u|^2 dx - \int \int_M f |\nabla_g (\log f + U[f])|^2 d\mu dx.$$

## Theorem

*3DNS + NL Fokker-Planck eqns with \* . Then*

$$E(t) = \frac{1}{2} \int |u|^2 dx + \int \left\{ f \log f + \frac{1}{2} (U[f]) f \right\} dx d\mu.$$

*is nondecreasing on solutions. If  $(u, f)$  is a smooth solution then*

$$\frac{dE}{dt} = -\nu \int |\nabla_x u|^2 dx - \int \int_M f |\nabla_g (\log f + U[f])|^2 d\mu dx.$$

*If the smooth solution is time independent, then  $u = 0$  and  $f$  solves the Onsager equation*

$$f = Z^{-1} e^{-U[f]}.$$

## NLFP + 2D time dependent Navier-Stokes

### Theorem

*(C-Seregin)* Let  $q \geq 4$ ,  $\operatorname{div}_x u_0 = 0$ ,  $u_0 \in W^{2,q}(\mathbb{T}^2)$ ,

# NLFP + 2D time dependent Navier-Stokes

## Theorem

*(C-Seregin)* Let  $q \geq 4$ ,  $\operatorname{div}_x u_0 = 0$ ,  $u_0 \in W^{2,q}(\mathbb{T}^2)$ ,  $f_0 > 0$ ,  $f_0(x, m) \in W^{1,q}(\mathbb{T}^2; L^2(M))$  and  $\int_M f_0(x, m) d\mu = 1$ .

## NLFP + 2D time dependent Navier-Stokes

### Theorem

*(C-Seregin) Let  $q \geq 4$ ,  $\operatorname{div}_x u_0 = 0$ ,  $u_0 \in W^{2,q}(\mathbb{T}^2)$ ,  $f_0 > 0$ ,  $f_0(x, m) \in W^{1,q}(\mathbb{T}^2; L^2(M))$  and  $\int_M f_0(x, m) d\mu = 1$ . Let  $T > 0$  be arbitrary. Let  $p > \frac{2q}{q-2}$ ,  $\alpha > \frac{d}{2} + 1$ .*



## NLFP + 2D time dependent Navier-Stokes

### Theorem

*(C-Seregin) Let  $q \geq 4$ ,  $\operatorname{div}_x u_0 = 0$ ,  $u_0 \in W^{2,q}(\mathbb{T}^2)$ ,  $f_0 > 0$ ,  $f_0(x, m) \in W^{1,q}(\mathbb{T}^2; L^2(M))$  and  $\int_M f_0(x, m) d\mu = 1$ . Let  $T > 0$  be arbitrary. Let  $p > \frac{2q}{q-2}$ ,  $\alpha > \frac{d}{2} + 1$ . There exists a constant  $K = K(T)$  bounded by a double exponential of  $T$ , and a unique solution  $(u, f)$  of the coupled NS-NLFP system such that*

## NLFP + 2D time dependent Navier-Stokes

### Theorem

*(C-Seregin) Let  $q \geq 4$ ,  $\operatorname{div}_x u_0 = 0$ ,  $u_0 \in W^{2,q}(\mathbb{T}^2)$ ,  $f_0 > 0$ ,  $f_0(x, m) \in W^{1,q}(\mathbb{T}^2; L^2(M))$  and  $\int_M f_0(x, m) d\mu = 1$ . Let  $T > 0$  be arbitrary. Let  $p > \frac{2q}{q-2}$ ,  $\alpha > \frac{d}{2} + 1$ . There exists a constant  $K = K(T)$  bounded by a double exponential of  $T$ , and a unique solution  $(u, f)$  of the coupled NS-NLFP system such that*

$$\|\nabla_x \nabla_x u\|_{L^p(0, T; L^q(\mathbb{T}^2))} \leq K,$$

$$\sup_{t \leq T} \|\nabla_x u(\cdot, t)\|_{L^\infty} \leq K,$$

and

$$\sup_{t \leq T} \|f(\cdot, t)\|_{W^{1,q}(\mathbb{T}^2; H^{-\alpha}(M))} \leq K.$$

hold.

## NLFP + 2D time dependent Navier-Stokes

### Theorem

*(C-Seregin) Let  $q \geq 4$ ,  $\operatorname{div}_x u_0 = 0$ ,  $u_0 \in W^{2,q}(\mathbb{T}^2)$ ,  $f_0 > 0$ ,  $f_0(x, m) \in W^{1,q}(\mathbb{T}^2; L^2(M))$  and  $\int_M f_0(x, m) d\mu = 1$ . Let  $T > 0$  be arbitrary. Let  $p > \frac{2q}{q-2}$ ,  $\alpha > \frac{d}{2} + 1$ . There exists a constant  $K = K(T)$  bounded by a double exponential of  $T$ , and a unique solution  $(u, f)$  of the coupled NS-NLFP system such that*

$$\|\nabla_x \nabla_x u\|_{L^p(0, T; L^q(\mathbb{T}^2))} \leq K,$$

$$\sup_{t \leq T} \|\nabla_x u(\cdot, t)\|_{L^\infty} \leq K,$$

and

$$\sup_{t \leq T} \|f(\cdot, t)\|_{W^{1,q}(\mathbb{T}^2; H^{-\alpha}(M))} \leq K.$$

hold.

Previous result, with Masmoudi: no a priori bounds and proof by contradiction, controlled loss of regularity,

2D

Navier-Stokes  
equations with  
singular  
forcing

Peter  
Constantin

Introduction

Kinetics:  
NLFP

Embedding in  
Fluid

Results on  
NS+NLFP

Navier-Stokes

## Related Results

with Fefferman, Titi, Zarnescu: for model with time-averaged velocity

2D

Navier-Stokes  
equations with  
singular  
forcing

Peter  
Constantin

Introduction

Kinetics:  
NLFP

Embedding in  
Fluid

Results on  
NS+NLFP

Navier-Stokes

## Related Results

with Fefferman, Titi, Zarnescu: for model with time-averaged velocity

Masmoudi, Zhang, Zhang: a priori bounds using the method of controlled loss of regularity

2D

Navier-Stokes  
equations with  
singular  
forcing

Peter  
Constantin

Introduction

Kinetics:  
NLFP

Embedding in  
Fluid

Results on  
NS+NLFP

Navier-Stokes

## Related Results

with Fefferman, Titi, Zarnescu: for model with time-averaged velocity

Masmoudi, Zhang, Zhang: a priori bounds using the method of controlled loss of regularity

Lin, Zhang: a corotational model with noncompact target.

# Navier Stokes Equation

$$\begin{aligned}\partial_t u + u \cdot \nabla u + \nabla p &= \nu \Delta u + \nabla \cdot \sigma \\ \nabla \cdot u &= 0\end{aligned}$$

The tensor  $\sigma_{ij}(x, t)$  : driving stress.

2D

Navier-Stokes  
equations with  
singular  
forcing

Peter  
Constantin

Introduction

Kinetics:  
NLFP

Embedding in  
Fluid

Results on  
NS+NLFP

Navier-Stokes

# Navier Stokes Equation

$$\begin{aligned}\partial_t u + u \cdot \nabla u + \nabla p &= \nu \Delta u + \nabla \cdot \sigma \\ \nabla \cdot u &= 0\end{aligned}$$

The tensor  $\sigma_{ij}(x, t)$  : driving stress.

Sufficient for regularity, if  $\sigma$  smooth

$$\int_0^T \|u\|_{L^\infty(dx)}^2 dt < \infty$$



# Navier Stokes Equation

$$\begin{aligned}\partial_t u + u \cdot \nabla u + \nabla p &= \nu \Delta u + \nabla \cdot \sigma \\ \nabla \cdot u &= 0\end{aligned}$$

The tensor  $\sigma_{ij}(x, t)$  : driving stress.

Sufficient for regularity, if  $\sigma$  smooth

$$\int_0^T \|u\|_{L^\infty(dx)}^2 dt < \infty$$

Exponent of the amplification factor  
of tracers, key quantity

$$\int_0^T \|\nabla u\|_{L^\infty(dx)} dt < \infty$$

2D

Navier-Stokes  
equations with  
singular  
forcing

Peter  
Constantin

Introduction

Kinetics:  
NLFP

Embedding in  
Fluid

Results on  
NS+NLFP

Navier-Stokes

## 2D, Bounded stress

### Theorem

Let  $\sigma \in L^\infty(dt dx)$ . Let  $u_0 \in L^2(dx)$ . There exists a *unique* weak solution of the forced 2D NS eqns, with

$$u \in L^\infty(dt)(L^2(dx)) \cap L^2(dt)(W^{1,2}(dx))$$

## 2D, Bounded stress

### Theorem

Let  $\sigma \in L^\infty(dt dx)$ . Let  $u_0 \in L^2(dx)$ . There exists a *unique* weak solution of the forced 2D NS eqns, with

$$u \in L^\infty(dt)(L^2(dx)) \cap L^2(dt)(W^{1,2}(dx))$$

Moreover,

$$\int_0^T \|\nabla u\|_{L^q(dx)}^{\frac{q}{q-1}} dt < \infty, \quad \forall q \geq 2$$

## 2D, Bounded stress

### Theorem

Let  $\sigma \in L^\infty(dt dx)$ . Let  $u_0 \in L^2(dx)$ . There exists a *unique* weak solution of the forced 2D NS eqns, with

$$u \in L^\infty(dt)(L^2(dx)) \cap L^2(dt)(W^{1,2}(dx))$$

Moreover,

$$\int_0^T \|\nabla u\|_{L^q(dx)}^{\frac{q}{q-1}} dt < \infty, \quad \forall q \geq 2$$

$$\int_0^T \|u\|_{L^\infty(dx)}^p dt < \infty, \quad \forall p < 2.$$

## Navier-Stokes with nearly singular forces

$$\partial_t u + u \cdot \nabla_x u - \nu \Delta_x u + \nabla_x p = \operatorname{div}_x \sigma, \quad \nabla_x \cdot u = 0$$

### Theorem

Let  $u$  be a solution of the 2D Navier-Stokes system with divergence-free initial data  $u_0 \in W^{1,2}(\mathbb{R}^2) \cap W^{1,r}(\mathbb{R}^2)$ . Let  $T > 0$  and let the forces  $\nabla \cdot \sigma$  obey

$$\begin{aligned} \sigma &\in L^1(0, T; L^\infty(\mathbb{R}^2)) \cap L^2(0, T; L^2(\mathbb{R}^2)) \\ \nabla \cdot \sigma &\in L^1(0, T; L^r(\mathbb{R}^2)) \cap L^2(0, T; L^2(\mathbb{R}^2)) \end{aligned}$$

with  $r > 2$ .

## Navier-Stokes with nearly singular forces

$$\partial_t u + u \cdot \nabla_x u - \nu \Delta_x u + \nabla_x p = \operatorname{div}_x \sigma, \quad \nabla_x \cdot u = 0$$

### Theorem

Let  $u$  be a solution of the 2D Navier-Stokes system with divergence-free initial data  $u_0 \in W^{1,2}(\mathbb{R}^2) \cap W^{1,r}(\mathbb{R}^2)$ . Let  $T > 0$  and let the forces  $\nabla \cdot \sigma$  obey

$$\begin{aligned} \sigma &\in L^1(0, T; L^\infty(\mathbb{R}^2)) \cap L^2(0, T; L^2(\mathbb{R}^2)) \\ \nabla \cdot \sigma &\in L^1(0, T; L^r(\mathbb{R}^2)) \cap L^2(0, T; L^2(\mathbb{R}^2)) \end{aligned}$$

with  $r > 2$ . Let

$$\|\sigma\|_{L^\infty} \sim K, \quad \|\nabla \cdot \sigma\|_{L^r} \sim B$$

Then

Peter  
Constantin

$$\int_0^T \|\nabla u(t)\|_{L^\infty} dt \leq K \log_*(B)$$

Introduction

Kinetics:  
NLFP

Embedding in  
Fluid

Results on  
NS+NLFP

Navier-Stokes

Then

$$\int_0^T \|\nabla u(t)\|_{L^\infty} dt \leq K \log_*(B)$$

and also

$$\frac{1}{M} \sum_{q=1}^M \int_0^T \|\Delta_q \nabla u(t)\|_{L^\infty} dt \leq K$$



Then

$$\int_0^T \|\nabla u(t)\|_{L^\infty} dt \leq K \log_*(B)$$

and also

$$\frac{1}{M} \sum_{q=1}^M \int_0^T \|\Delta_q \nabla u(t)\|_{L^\infty} dt \leq K$$

with  $K$  depending on  $T$ , norms of  $\sigma$  and the initial velocity, but not on gradients of  $\sigma$  nor  $M$ , and  $B$  depending on norms of the spatial gradients of  $\sigma$ .

Then

$$\int_0^T \|\nabla u(t)\|_{L^\infty} dt \leq K \log_*(B)$$

and also

$$\frac{1}{M} \sum_{q=1}^M \int_0^T \|\Delta_q \nabla u(t)\|_{L^\infty} dt \leq K$$

with  $K$  depending on  $T$ , norms of  $\sigma$  and the initial velocity, but not on gradients of  $\sigma$  nor  $M$ , and  $B$  depending on norms of the spatial gradients of  $\sigma$ .

$$u = \sum_{q=-1}^{\infty} \Delta_q(u)$$

Littlewood-Paley decomposition

## Hölder regularity of weak solution

$\Omega_1 \Subset \Omega$  domains in  $\mathbb{R}^2$ ,  $0 < T_1 < T$ .

## Hölder regularity of weak solution

$\Omega_1 \Subset \Omega$  domains in  $\mathbb{R}^2$ ,  $0 < T_1 < T$ .

$$Q = \Omega \times (-T, 0), \quad Q_1 = \Omega_1 \times (-T_1, 0).$$

## Hölder regularity of weak solution

$\Omega_1 \Subset \Omega$  domains in  $\mathbb{R}^2$ ,  $0 < T_1 < T$ .

$$Q = \Omega \times (-T, 0), \quad Q_1 = \Omega_1 \times (-T_1, 0).$$

Parabolic balls:

## Hölder regularity of weak solution

$\Omega_1 \Subset \Omega$  domains in  $\mathbb{R}^2$ ,  $0 < T_1 < T$ .

$$Q = \Omega \times (-T, 0), \quad Q_1 = \Omega_1 \times (-T_1, 0).$$

Parabolic balls:  $Q(z_0, R) = B(x_0, R) \times (t_0 - R^2, t_0)$ , where  $z_0 = (x_0, t_0)$ ,  $x_0 \in \mathbb{R}^2$ ,  $t_0 \in \mathbb{R}$ .

## Hölder regularity of weak solution

$\Omega_1 \Subset \Omega$  domains in  $\mathbb{R}^2$ ,  $0 < T_1 < T$ .

$$Q = \Omega \times (-T, 0), \quad Q_1 = \Omega_1 \times (-T_1, 0).$$

Parabolic balls:  $Q(z_0, R) = B(x_0, R) \times (t_0 - R^2, t_0)$ , where  $z_0 = (x_0, t_0)$ ,  $x_0 \in \mathbb{R}^2$ ,  $t_0 \in \mathbb{R}$ . Means:

## Hölder regularity of weak solution

$\Omega_1 \Subset \Omega$  domains in  $\mathbb{R}^2$ ,  $0 < T_1 < T$ .

$$Q = \Omega \times (-T, 0), \quad Q_1 = \Omega_1 \times (-T_1, 0).$$

Parabolic balls:  $Q(z_0, R) = B(x_0, R) \times (t_0 - R^2, t_0)$ , where  $z_0 = (x_0, t_0)$ ,  $x_0 \in \mathbb{R}^2$ ,  $t_0 \in \mathbb{R}$ . Means:

$$(f)_{z_0, R} = \frac{1}{|Q(z_0, R)|} \int_{Q(z_0, R)} f(z) dz,$$
$$[p]_{x_0, R} = \frac{1}{|B(x_0, R)|} \int_{B(x_0, R)} p(x) dx.$$



## Spaces

For  $0 < \gamma < 1$ ,

## Spaces

For  $0 < \gamma < 1$ , space  $M_{2,\gamma}(Q)$ ,

## Spaces

For  $0 < \gamma < 1$ , space  $M_{2,\gamma}(Q)$ , seminorm

$$\|\sigma\|_{M_{2,\gamma}(Q)} = \sup_{Q(z_0,R) \subset Q} R^{1-\gamma} \left( \frac{1}{|Q(z_0,R)|} \int_{Q(z_0,R)} |\sigma(z) - (\sigma)_{z_0,R}|^2 dz \right)^{\frac{1}{2}} < \infty.$$

## Spaces

For  $0 < \gamma < 1$ , space  $M_{2,\gamma}(Q)$ , seminorm

$$\|\sigma\|_{M_{2,\gamma}(Q)} = \sup_{Q(z_0,R) \subset Q} R^{1-\gamma} \left( \frac{1}{|Q(z_0,R)|} \int_{Q(z_0,R)} |\sigma(z) - (\sigma)_{z_0,R}|^2 dz \right)^{\frac{1}{2}} < \infty.$$

Note

$$L^r(Q) \subset M_{2,\gamma}(Q)$$

for  $r \geq \frac{4}{1-\gamma}$ .

## Theorem

*(C-Seregin) Let*

$$u \in L^4(Q; \mathbb{R}^2), \quad p \in L^2(Q), \quad \sigma \in M_{2,\gamma}(Q; \mathbb{M}^{2 \times 2})$$

*with  $0 \leq \gamma < 1$ , satisfying the Navier-Stokes equations*

$$\partial_t u + u \cdot \nabla u - \nu \Delta u = -\nabla p + \operatorname{div} \sigma, \quad \operatorname{div} u = 0$$

*in  $Q$  in the sense of distributions.*

## Theorem

*(C-Seregin) Let*

$$u \in L^4(Q; \mathbb{R}^2), \quad p \in L^2(Q), \quad \sigma \in M_{2,\gamma}(Q; \mathbb{M}^{2 \times 2})$$

*with  $0 \leq \gamma < 1$ , satisfying the Navier-Stokes equations*

$$\partial_t u + u \cdot \nabla u - \nu \Delta u = -\nabla p + \operatorname{div} \sigma, \quad \operatorname{div} u = 0$$

*in  $Q$  in the sense of distributions. Then*

$$u \in C^\gamma(\overline{Q_1})$$

*if  $0 < \gamma < 1$  and*

$$u \in BMO(Q_1)$$

*if  $\gamma = 0$ .*

## Additional results

Let  $H$  and  $V$  be the  $L^2$  and  $H^1$  spaces of divergence-free functions.

2D

Navier-Stokes  
equations with  
singular  
forcing

Peter  
Constantin

Introduction

Kinetics:  
NLFP

Embedding in  
Fluid

Results on  
NS+NLFP

Navier-Stokes

## Additional results

Let  $H$  and  $V$  be the  $L^2$  and  $H^1$  spaces of divergence-free functions.

### Theorem

*(C-S) Let  $u \in L^\infty(0, T; H) \cap L^2(0, T; V)$ ,  $p \in L^2(0, T; L^2(\mathbb{T}^2))$  be a solution of the initial value problem*

$$\partial_t u + u \cdot \nabla u - \nu \Delta u + \nabla p = \operatorname{div} \sigma, \quad \operatorname{div} u = 0,$$

$$u(\cdot, 0) = u_0(\cdot) \in H,$$

*where  $\sigma \in L^r(\mathbb{T}^2 \times (0, T); \mathbb{M}^{2 \times 2})$  with  $r \geq 4$ .*



## Additional results

Let  $H$  and  $V$  be the  $L^2$  and  $H^1$  spaces of divergence-free functions.

### Theorem

*(C-S) Let  $u \in L^\infty(0, T; H) \cap L^2(0, T; V)$ ,  $p \in L^2(0, T; L^2(\mathbb{T}^2))$  be a solution of the initial value problem*

$$\partial_t u + u \cdot \nabla u - \nu \Delta u + \nabla p = \operatorname{div} \sigma, \quad \operatorname{div} u = 0,$$

$$u(\cdot, 0) = u_0(\cdot) \in H,$$

*where  $\sigma \in L^r(\mathbb{T}^2 \times (0, T); \mathbb{M}^{2 \times 2})$  with  $r \geq 4$ . Then, given  $s > 0$ , there exists a constant  $C_s$  depending only on  $s, \nu$ , the norm of  $u_0$  in  $H$ , the norm of  $\sigma$  in  $L^r(\mathbb{T}^2)$ , such that*

$$\|u\|_{L^\infty(\mathbb{T}^2 \times (s, T))} \leq C_s.$$

*Moreover, the function  $u$  is Hölder continuous in  $\mathbb{T}^2 \times [s, T]$  with exponent  $\gamma = 1 - \frac{4}{r}$ .*

2D

Navier-Stokes  
equations with  
singular  
forcing

Peter  
Constantin

Introduction

Kinetics:  
NLFP

Embedding in  
Fluid

Results on  
NS+NLFP

Navier-Stokes

## Idea of proof of Hölder continuity for NS

Local iterative estimates for  $L^4$  space-time integrals of the velocity, in the spirit of De Giorgi, Campanato.

## Idea of proof of Hölder continuity for NS

Local iterative estimates for  $L^4$  space-time integrals of the velocity, in the spirit of De Giorgi, Campanato. Pressure

$$p = R_i R_j (\sigma_{ij} - u_i u_j),$$

The iteration relates integrals on smaller parabolic cubes to integrals on larger ones.

## Idea of proof of Hölder continuity for NS

Local iterative estimates for  $L^4$  space-time integrals of the velocity, in the spirit of De Giorgi, Campanato. Pressure

$$p = R_i R_j (\sigma_{ij} - u_i u_j),$$

The iteration relates integrals on smaller parabolic cubes to integrals on larger ones. For the iterative procedure to succeed, the modulus of absolute continuity of the map

$$\Omega \subset \{\mathbb{T}^2 \times (0, T)\} \mapsto \int_{\Omega} |u(x, t)|^4 dx dt,$$

needs to be controlled uniformly a priori, to guarantee that such an integral is arbitrarily small, if the parabolic Lebesgue measure of  $\Omega$  is small enough.

$$\Phi(u; z_0, \varrho) = \left( \int_{Q(z_0, \varrho)} |u - (u)_{z_0, \varrho}|^4 dz \right)^{\frac{1}{2}},$$

$$\Phi(u; z_0, \varrho) = \left( \int_{Q(z_0, \varrho)} |u - (u)_{z_0, \varrho}|^4 dz \right)^{\frac{1}{2}},$$

$$\Psi(u; z_0, \varrho) = \left( \int_{Q(z_0, \varrho)} |u|^4 dz \right)^{\frac{1}{2}},$$

$$\Phi(u; z_0, \varrho) = \left( \int_{Q(z_0, \varrho)} |u - (u)_{z_0, \varrho}|^4 dz \right)^{\frac{1}{2}},$$

$$\Psi(u; z_0, \varrho) = \left( \int_{Q(z_0, \varrho)} |u|^4 dz \right)^{\frac{1}{2}},$$

$$D(p; z_0, \varrho) = \int_{Q(z_0, \varrho)} |p - [p]_{x_0, \varrho}|^2 dz.$$

$$\Phi(u; z_0, \varrho) = \left( \int_{Q(z_0, \varrho)} |u - (u)_{z_0, \varrho}|^4 dz \right)^{\frac{1}{2}},$$

$$\Psi(u; z_0, \varrho) = \left( \int_{Q(z_0, \varrho)} |u|^4 dz \right)^{\frac{1}{2}},$$

$$D(p; z_0, \varrho) = \int_{Q(z_0, \varrho)} |p - [p]_{x_0, \varrho}|^2 dz.$$

## Lemma

Let the function  $v \in L^4(Q(z_0, R))$  satisfy the heat equation

$$\partial_t v - \Delta v = 0$$

in  $Q(z_0, R)$ .



$$\Phi(u; z_0, \varrho) = \left( \int_{Q(z_0, \varrho)} |u - (u)_{z_0, \varrho}|^4 dz \right)^{\frac{1}{2}},$$

$$\Psi(u; z_0, \varrho) = \left( \int_{Q(z_0, \varrho)} |u|^4 dz \right)^{\frac{1}{2}},$$

$$D(p; z_0, \varrho) = \int_{Q(z_0, \varrho)} |p - [p]_{x_0, \varrho}|^2 dz.$$

## Lemma

Let the function  $v \in L^4(Q(z_0, R))$  satisfy the heat equation

$$\partial_t v - \Delta v = 0$$

in  $Q(z_0, R)$ . Then

$$\Phi(v; z_0, \varrho) \leq c \left( \frac{\varrho}{R} \right)^4 \Phi(v; z_0, R)$$

for all  $0 < \varrho \leq R$ .

## Lemma

Given  $G \in L^2(Q(z_0, R); \mathbb{M}^{2 \times 2})$ , there exists a unique function

$$w \in C([t_0 - R^2, t_0]; L^2(B(x_0, R); \mathbb{R}^2)) \cap \\ L^2([t_0 - R^2, t_0]; W^{1,2}(B(x_0, R); \mathbb{R}^2))$$

## Lemma

Given  $G \in L^2(Q(z_0, R); \mathbb{M}^{2 \times 2})$ , there exists a unique function

$$w \in C([t_0 - R^2, t_0]; L^2(B(x_0, R); \mathbb{R}^2)) \cap \\ L^2([t_0 - R^2, t_0]; W^{1,2}(B(x_0, R); \mathbb{R}^2))$$

such that

$$\partial_t w - \Delta w = -\operatorname{div} G$$

in  $Q(z_0, R)$

## Lemma

Given  $G \in L^2(Q(z_0, R); \mathbb{M}^{2 \times 2})$ , there exists a unique function

$$w \in C([t_0 - R^2, t_0]; L^2(B(x_0, R); \mathbb{R}^2)) \cap \\ L^2([t_0 - R^2, t_0]; W^{1,2}(B(x_0, R); \mathbb{R}^2))$$

such that

$$\partial_t w - \Delta w = -\operatorname{div} G$$

in  $Q(z_0, R)$  and

$$w = 0$$

on the parabolic boundary of  $Q(z_0, R)$ .

## Lemma

Given  $G \in L^2(Q(z_0, R); \mathbb{M}^{2 \times 2})$ , there exists a unique function

$$w \in C([t_0 - R^2, t_0]; L^2(B(x_0, R); \mathbb{R}^2)) \cap \\ L^2([t_0 - R^2, t_0]; W^{1,2}(B(x_0, R); \mathbb{R}^2))$$

such that

$$\partial_t w - \Delta w = -\operatorname{div} G$$

in  $Q(z_0, R)$  and

$$w = 0$$

on the parabolic boundary of  $Q(z_0, R)$ . Moreover, the function  $w$  satisfies the estimates:

$$|w|_{2, Q(z_0, R)}^2 \equiv \sup_{t_0 - R^2 < t < t_0} \|w(\cdot, t)\|_{2, B(x_0, R)}^2 + \|\nabla w\|_{2, Q(z_0, R)}^2$$

$$\leq 2\|G\|_{2, Q(z_0, R)}^2,$$

$$\Phi(w; z_0, R) \leq c|w|_{2, Q(z_0, R)}^2.$$

## Lemma

*For solutions of NSE, we have*

$$\begin{aligned} \Phi(u; z_0, \varrho) \leq c \left\{ \left[ \left( \frac{\varrho}{R} \right)^4 + \Psi(u; z_0, R) \right] \Phi(u; z_0, R) + \right. \\ \left. + D(p; z_0, R) + MR^{2+2\gamma} \right\} \end{aligned}$$

## Lemma

*For solutions of NSE, we have*

$$\begin{aligned} \Phi(u; z_0, \varrho) \leq c \left\{ \left[ \left( \frac{\varrho}{R} \right)^4 + \Psi(u; z_0, R) \right] \Phi(u; z_0, R) + \right. \\ \left. + D(p; z_0, R) + MR^{2+2\gamma} \right\} \end{aligned}$$

*whenever  $Q(z_0, R) \subset Q$  and  $0 < \varrho \leq R$ .*

## Lemma

For solutions of NSE, we have

$$\begin{aligned} \Phi(u; z_0, \varrho) \leq c \left\{ \left[ \left( \frac{\varrho}{R} \right)^4 + \Psi(u; z_0, R) \right] \Phi(u; z_0, R) + \right. \\ \left. + D(p; z_0, R) + MR^{2+2\gamma} \right\} \end{aligned}$$

whenever  $Q(z_0, R) \subset Q$  and  $0 < \varrho \leq R$ . Here,  
 $M = \|\sigma\|_{M_{2,\gamma}(Q)}^2$ .

## Lemma

$$D(p; z_0, \varrho) \leq c \left[ \left( \frac{\varrho}{R} \right)^4 D(p; z_0, R) + \Psi(u; z_0, R) \Phi(u; z_0, R) + MR^{2+2\gamma} \right]$$



## Lemma

For solutions of NSE, we have

$$\Phi(u; z_0, \varrho) \leq c \left\{ \left[ \left( \frac{\varrho}{R} \right)^4 + \Psi(u; z_0, R) \right] \Phi(u; z_0, R) + D(p; z_0, R) + MR^{2+2\gamma} \right\}$$

whenever  $Q(z_0, R) \subset Q$  and  $0 < \varrho \leq R$ . Here,  
 $M = \|\sigma\|_{M_{2,\gamma}(Q)}^2$ .

## Lemma

$$D(p; z_0, \varrho) \leq c \left[ \left( \frac{\varrho}{R} \right)^4 D(p; z_0, R) + \Psi(u; z_0, R) \Phi(u; z_0, R) + MR^{2+2\gamma} \right]$$

whenever  $Q(z_0, R) \subset Q$  and  $0 < \varrho \leq R$ .

The following result is used to control the modulus of absolute continuity.

## Theorem

*Let  $u \in L^\infty(0, T; H) \cap L^2(0, T; V)$  be a solution of the 2D Navier-Stokes equations with initial data  $u_0 \in H \cap L^r(\mathbb{T}^2)$  and  $\sigma \in L^r(\mathbb{T}^2 \times (0, T); \mathbb{M}^{2 \times 2})$  with  $r \geq 4$ . There exists a constant  $K$  depending only on the norm  $\|\sigma\|_{L^r(\mathbb{T}^2 \times (0, T))}$ ,  $\nu$ ,  $T$  and the norm of  $u_0$  in  $H \cap L^r(\mathbb{T}^2)$  such that*

$$\sup_{0 \leq t \leq T} \|u(\cdot, t)\|_{L^r(\mathbb{T}^2)} \leq K.$$

## Generalized Ladyzhenskaya inequalities

The previous result is based on the inequality

$$\|f\|_{L^{2r}(\mathbb{R}^2)}^2 \leq C_r \|f\|_{L^r(\mathbb{R}^2)} \|\nabla f\|_{L^2(\mathbb{R}^2)}$$

## Generalized Ladyzhenskaya inequalities

The previous result is based on the inequality

$$\|f\|_{L^{2r}(\mathbb{R}^2)}^2 \leq C_r \|f\|_{L^r(\mathbb{R}^2)} \|\nabla f\|_{L^2(\mathbb{R}^2)}$$

that holds for all  $f \in L^r(\mathbb{R}^2)$  with  $\nabla f \in L^2(\mathbb{R}^2)$ ,

## Generalized Ladyzhenskaya inequalities

The previous result is based on the inequality

$$\|f\|_{L^{2r}(\mathbb{R}^2)}^2 \leq C_r \|f\|_{L^r(\mathbb{R}^2)} \|\nabla f\|_{L^2(\mathbb{R}^2)}$$

that holds for all  $f \in L^r(\mathbb{R}^2)$  with  $\nabla f \in L^2(\mathbb{R}^2)$ , generalizing the  $r = 2$  inequality due to Ladyzhenskaya.

## Generalized Ladyzhenskaya inequalities

The previous result is based on the inequality

$$\|f\|_{L^{2r}(\mathbb{R}^2)}^2 \leq C_r \|f\|_{L^r(\mathbb{R}^2)} \|\nabla f\|_{L^2(\mathbb{R}^2)}$$

that holds for all  $f \in L^r(\mathbb{R}^2)$  with  $\nabla f \in L^2(\mathbb{R}^2)$ , generalizing the  $r = 2$  inequality due to Ladyzhenskaya. In turn, the inequality above is the  $n = 2$  particular case of

$$\|f\|_{L^{2r}(\mathbb{R}^n)}^2 \leq C \|f\|_{L^r(\mathbb{R}^n)} \|\nabla f\|_{B_2^{0,n}(\mathbb{R}^n)}$$

## Generalized Ladyzhenskaya inequalities

The previous result is based on the inequality

$$\|f\|_{L^{2r}(\mathbb{R}^2)}^2 \leq C_r \|f\|_{L^r(\mathbb{R}^2)} \|\nabla f\|_{L^2(\mathbb{R}^2)}$$

that holds for all  $f \in L^r(\mathbb{R}^2)$  with  $\nabla f \in L^2(\mathbb{R}^2)$ , generalizing the  $r = 2$  inequality due to Ladyzhenskaya. In turn, the inequality above is the  $n = 2$  particular case of

$$\|f\|_{L^{2r}(\mathbb{R}^n)}^2 \leq C \|f\|_{L^r(\mathbb{R}^n)} \|\nabla f\|_{B_2^{0,n}(\mathbb{R}^n)}$$

valid for all  $r \geq \frac{n}{2}$ , with  $B_2^{0,n}$  the Besov space with norm

$$\|f\|_{B_q^{s,p}(\mathbb{R}^n)} = \left[ \sum_{j=-\infty}^{\infty} \lambda_j^{qs} \|\Delta_j f\|_{L^p(\mathbb{R}^n)}^q \right]^{\frac{1}{q}}$$

## References

- P. Constantin, Smoluchowski Navier-Stokes systems, Contemporary Mathematics **429** G-Q Chen, E. Hsu, M. Pinsky editors, AMS, Providence (2007), 85-109
- P. Constantin, N. Masmoudi, Global well-posedness for a Smoluchowski equation coupled with Navier-Stokes equations in 2D, Commun. Math. Phys. **278** (2008), 179-191.
- P. Constantin, G. Seregin, Hölder Continuity of Solutions of 2D Navier-Stokes Equations with Singular Forcing, to appear
- P. Constantin, G. Seregin, Global regularity of solutions of coupled Navier-Stokes equations and nonlinear Fokker Planck equations, to appear
- P. Constantin, The Onsager equation for corpora, J. Comp. Theor. Nanoscience, 2009, to appear.
- P. Constantin, A. Zlatos, On the high intensity limit of interacting corpora, CMS, to appear.