

## ▼ Asymptotic solution of SLFM model.

The SLFM model is given by:

$$\tilde{N}(z,t) \frac{\partial^2 Q_F}{\partial z^2} = -A Q_F Q_0$$

$$\tilde{N}(z,t) \frac{\partial^2 Q_0}{\partial z^2} = -r A Q_F Q_0.$$

where  $\tilde{N}(n,t) = \langle N | z=n \rangle_t$ .

### ● Slow chemistry approximation.

As  $A \rightarrow 0$  the equations become:

$$\frac{\partial^2 Q_F}{\partial z^2} = \frac{\partial^2 Q_0}{\partial z^2} = 0.$$

Since  $A$  does not appear in the highest derivative, this is a regular problem  $\leadsto$  the limiting solution  $A \rightarrow 0$  converges to the solution at  $A=0$ .

The solution to these equations is:

$$Q_F = a_F z + b_F$$

$$Q_0 = a_0 z + b_0.$$

Now apply the boundary conditions:

$$Q_F = z$$

$$Q_0 = 1 - z.$$

The SLFM model converges to the slow chemistry approximation.

### ● The fast chemistry case.

First of all, eliminate  $Q_0$ . Use  $\langle z | z=n \rangle = n$  on

$$z = \frac{r Y_F - Y_0 + 1}{r+1}, \text{ for each point in space } \Rightarrow$$

$$\Rightarrow \langle z | z=n \rangle = \frac{r \langle Y_F | z=n \rangle - \langle Y_0 | z=n \rangle + 1}{r+1} = n \Rightarrow$$

$$\Rightarrow \frac{r Q_F - Q_0 + 1}{r+1} = z \Rightarrow \underline{Q_0(z,t) = r Q_F(z,t) + 1 - (r+1)z.}$$

The ODE becomes:  $\tilde{N} \frac{\partial^2 Q_F}{\partial z^2} = -A Q_F (r Q_F + 1 - (r+1)z)$ .

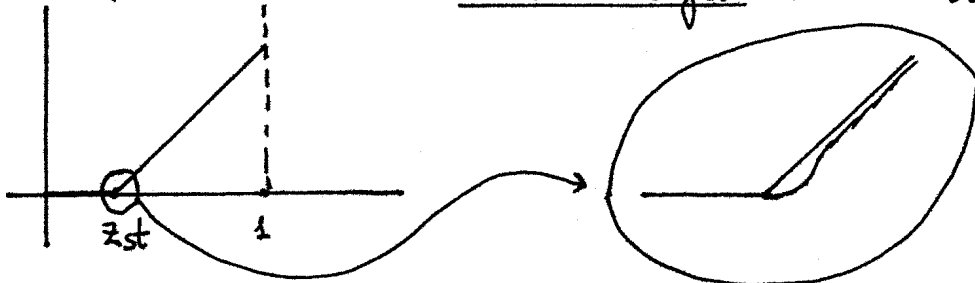
Define  $\epsilon = 1/A$ . Then:

$$\epsilon \tilde{N} \frac{\partial^2 Q_F}{\partial z^2} = -Q_F (r Q_F + 1 - (r+1)z).$$

For  $\epsilon = 0$ :  $Q_F (r Q_F + 1 - (r+1)z) = 0 \leadsto$  equilibrium chemistry approximation:

$$Q_F(z) = \frac{z - z_{st}}{1 - z_{st}} H(z - z_{st}) = Q_F^{eq}(z)$$

For  $\epsilon > 0$  there is a corner layer at  $z = z_{st}$



and  $Q_F(z) = Q_F^{eq}(z) + \delta(\epsilon) y_1(z) + \delta^2(\epsilon) y_2(z) + O(\delta^3)$ .

corner layer terms that smooth the corner at  $z = z_{st}$ .

Physically, the corner layer corresponds to the thin reaction zone and

$\delta(\epsilon)$  = thickness of the reaction zone given reaction rate  $A = 1/\epsilon$ .

To obtain the behaviour in that zone we must map  $z_{st} - \delta < z < z_{st} + \delta \leadsto -(\text{big}) < \xi < +(\text{big})$ .

Define  $\xi = \frac{1}{\delta} (z - z_{st})$ .

Then:  $\frac{\partial Q_F}{\partial z} = \frac{\partial Q_F}{\partial \xi} \frac{\partial \xi}{\partial z} = \frac{1}{\delta} \frac{\partial Q_F}{\partial \xi} \Rightarrow \frac{\partial^2 Q_F}{\partial z^2} = \frac{1}{\delta^2} \frac{\partial^2 Q_F}{\partial \xi^2}$

and  $z = z_{st} + \delta \xi$ . Substitute to the ODE:

$$\frac{\epsilon \tilde{N}}{\delta^2(\epsilon)} \frac{\partial^2 Q_F}{\partial \xi^2} = -Q_F [r Q_F + 1 - (r+1)(z_{st} + \delta(\epsilon)\xi)] \quad \text{as } \epsilon \rightarrow 0^+$$

Note that

$$1 - (r+1)z_{st} = 1 - (r+1)\frac{1}{r+1} = 1 - 1 = 0.$$

so the ODE simplifies to:

$$\frac{\epsilon \tilde{N}}{\delta^2(\epsilon)} \frac{\partial^2 Q_F}{\partial \xi^2} = -Q_F [rQ_F - (r+1)\delta(\epsilon)\xi]$$

The solution is:

$$Q_F(\xi; \epsilon) = y_0(\xi) + \delta(\epsilon)y_1(\xi) + O(\delta^2).$$

$$\begin{aligned} \text{For } \epsilon=0 \Rightarrow \delta(\epsilon)=0 \Rightarrow Q_F(\xi; 0) &= y_0(\xi) \\ \text{also } Q_F(\xi; 0) &= Q_F^{\text{eq}}(z_{st} + \delta(\epsilon)\xi) = Q_F^{\text{eq}}(z_{st}) = 0 \} \Rightarrow \\ \Rightarrow y_0(\xi) &= 0. \Rightarrow Q_F(\xi; \epsilon) = \delta(\epsilon)y_1(\xi) + O(\delta^2). \end{aligned}$$

Substitute to the ODE:

$$\frac{\epsilon \tilde{N}}{\delta^2} \frac{\partial^2 Q_F}{\partial \xi^2}$$

$$\frac{\epsilon N}{\delta^2} \frac{\partial^2 (\delta y_1)}{\partial \xi^2} = -\delta y_1 [r\delta y_1 - (r+1)\delta\xi] + O(\delta^3) \Leftrightarrow$$

$$\Leftrightarrow \frac{\epsilon N}{\delta^2} \frac{\partial^2 y_1}{\partial \xi^2} = -y_1 [r\delta y_1 - (r+1)\delta\xi] + O(\delta^2) \Leftrightarrow$$

$$\Leftrightarrow \frac{\epsilon N}{\delta^3} \frac{\partial^2 y_1}{\partial \xi^2} = -ry_1^2 + (r+1)\xi y_1$$

The RHS has finite value as  $\epsilon \rightarrow 0$  therefore so must the LHS:

$$\frac{\partial^2 y_1}{\partial \xi^2} \sim \frac{\delta^3}{\epsilon N}$$

Recall that  $\delta$  is the thickness of reaction zone  $\Rightarrow \partial^2 y_1 / \partial \xi^2$  has order 1 values and:

$$\frac{\delta^3}{\epsilon N} \sim 1 \Leftrightarrow \delta^3 \sim \epsilon N \Leftrightarrow \delta \sim (\epsilon N)^{1/3}.$$

It follows from this result that

$$\text{at } \xi=0 \ (z=z_{st}): Q_F \sim \delta(\epsilon)y_1(0) \sim (\epsilon N_{st})^{1/3} \sim (N_{st}/A)^{1/3}.$$

We define the local Damköhler number from this result:

$$Da_s(t) = \frac{1-z_{st}}{z_{st}} \frac{A}{N_{st}(t)}.$$

Purely for convenience we write:

$$Q_F(z_{st}) \sim \left[ \frac{z_{st}}{1-z_{st}} \frac{N_{st}}{A} \right]^{1/3}$$

We define the local Damköhler number to be:

$$Da_s(t) = \frac{1-z_{st}}{z_{st}} \frac{A}{N_{st}(t)}$$

Then it follows that

~~$$Q_F(z_{st}) \sim Da_s(t)$$
$$Q_F(z_{st}, t) \sim Da_s(t)^{-1/3}$$~~

$$Q_F(z_{st}, t) \sim Da_s(t)^{-1/3}$$

we also obtain

⑥

▼ The Quasi-steady state condition.

$$\alpha = \frac{1}{A Q_F Q_0} \frac{\partial Q_F}{\partial t} \ll 1.$$

Local Damköhler number  $Da_s(t) = \frac{1-z_{st}}{z_{st}} \frac{A}{N_{st}(t)}$  where  $N_{st}(t) = \tilde{N}(z_{st}, t)$ .

(derivation?).

Check consistency with asymptotic approximation of the quasi-state solution. (dominant balance).

This is given by:  $Q_F(z_{st}, t) = Da_s(t)^{-1/3}$ . (derivation?).

$$\begin{aligned} \frac{\partial Da_s}{\partial t} &= \frac{1-z_{st}}{z_{st}} A \frac{\partial}{\partial t} \left( \frac{1}{N_{st}} \right) = - \frac{1-z_{st}}{z_{st}} A \left( \frac{-1}{N_{st}^2} \right) \frac{\partial N_{st}}{\partial t} = \\ &= - \frac{Da_s(t)}{N_{st}(t)} \frac{\partial N_{st}}{\partial t} \end{aligned}$$

therefore,

$$\begin{aligned} \left. \frac{\partial Q_F}{\partial t} \right|_{z=z_{st}} &= (-1/3) Da_s^{-4/3}(t) \frac{\partial Da_s}{\partial t} = \\ &= (-1/3) Da_s^{-4/3}(t) \left[ - \frac{Da_s(t)}{N_{st}(t)} \frac{\partial N_{st}}{\partial t} \right] \\ &= \frac{1}{3} \frac{Da_s^{-1/3}}{N_{st}} \frac{\partial N_{st}}{\partial t} \end{aligned}$$

Also  $A Q_F Q_0 = r A Da_s^{-2/3}(t)$ . , so we obtain:

$$\begin{aligned} \alpha &= \frac{1}{A Q_F Q_0} \frac{\partial Q_F}{\partial t} = \frac{1}{r A Da_s^{-2/3}(t)} \frac{Da_s^{-4/3}}{N_{st}} \frac{\partial N_{st}}{\partial t} = \\ &= \frac{1}{r A Da_s^{-4/3}(t)} \frac{1}{N_{st}} \frac{\partial N_{st}}{\partial t} = \\ &= \frac{1}{r Da_s^{-4/3}(t)} \frac{N_{st}}{A} \frac{1}{N_{st}^2} \frac{\partial N_{st}}{\partial t} = \left. \begin{aligned} &= \frac{1-z_{st}}{r Da_s^{-4/3}(t)} \frac{1-z_{st}}{z_{st}} \frac{1}{Da_s(t)} \frac{1}{N_{st}^2} \frac{\partial N_{st}}{\partial t} = \\ &= \frac{1-z_{st}}{r z_{st}} \frac{1}{Da_s^{2/3}(t)} \frac{1}{N_{st}^2} \frac{\partial N_{st}}{\partial t} \end{aligned} \right\} \begin{aligned} &Da_s(t) = \frac{1-z_{st}}{z_{st}} \frac{A}{N_{st}(t)} \end{aligned} \end{aligned}$$

and  $z_{st} = \frac{1}{r+1} \Rightarrow \frac{1-z_{st}}{r z_{st}} = 1$  (wow!).

⑦

therefore:  $a = \frac{1}{Da_s^{2/3}(t)} \frac{1}{Nst^2} \frac{\partial Nst}{\partial t} \ll 1$

must be true in order for the quasi-steady equation to be a dominant balance.

Under this condition:

$\tilde{N}(z,t) \frac{\partial^2 Q_F}{\partial z^2} = -A Q_F Q_0$
$\tilde{N}(z,t) \frac{\partial^2 Q_0}{\partial z^2} = -r A Q_F Q_0.$