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L-C-Z族孤子方程所对应的 完全可积的Bargmann系统*

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摘要 本文获得一个新的有限维对合系, 并由此证明L-C-Z谱问题(2.1)在Bargmann约束: $q = \langle \varphi_1, \varphi_1 \rangle - \langle \varphi_2, \varphi_2 \rangle$, $r = \langle \varphi_1, \varphi_1 \rangle + \langle \varphi_2, \varphi_2 \rangle$ 下被非线性化为一个Liouville完全可积的新的Hamilton系统. 最后, 我们给出L-C-Z族方程解的对合表示.

关键词 L-C-Z族; Bargmann约束; 对合系; 对合解; Hamilton系统; Liouville完全可积.

0 引言

孤子方程,

在孤子理论中, 扩充Liouville完全可积的有限维Hamilton系统是一项十分重要的课题. 其关键在于寻求对合的函数系, 常用的方法是通过发展方程的谱技巧等去求得, 进而产生新的完全可积系统[1, 2]. 最近, 曹策问教授提出了Lax组非线性化产生有限维完全可积系统的思想[3], 并且成功地找到了许多有限维完全可积系统[4]. 本文在文[5]的基础上, 利用“非线性化技术”将L-C-Z族方程的Lax组非线性化; 尔后证明Lax组空间部分的非线性化是一个Liouville完全可积系统(R^{2N} , $d\varphi_1 \wedge d\varphi_2$, $H = -iF$).

$$H = \langle i \wedge \varphi_1, \varphi_2 \rangle - \langle \varphi_1 + \varphi_2, \varphi_1 + \varphi_2 \rangle + \langle \varphi_1, \varphi_2 \rangle - \langle \varphi_1, \varphi_1 \rangle + \langle \varphi_2, \varphi_2 \rangle + \langle \varphi_1, \varphi_2 \rangle^2.$$

同时其时间部分的非线性化将产生N-对合系 $\{F_m\}$. 相容组 $(F_0), (F_m)$ 的对合解由Bargmann约束所决定的映射 f 映到高阶L-C-Z方程的解, 从而说明所有的L-C-Z方程的解均有对合表示.

1 L-C-Z族方程的Lax组非线性化

与L-C-Z族方程相联系的保谱问题为:

$$L\varphi = \xi\varphi, \quad L = L(u) = \frac{1}{i} \begin{pmatrix} r - \partial & q + r \\ r - q & r + \partial \end{pmatrix} \quad (1.1)$$

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$\varphi = (\varphi_1, \varphi_2)^T$; $\partial = \partial / \partial x$; ζ 为 (1.1) 的谱参数; $i = \sqrt{-1}$; 自变量 x 在所论区间 Ω ($\Omega = (-\infty, +\infty)$ 或 $(0, T)$) 内变化; $u(x) \triangleq (q(x), r(x))$ 是 (1.1) 的势函数, 且 $u(x)$ 在无穷远处衰减为零或以 T 为周期.

由 (5), 我们已经知道:

命题 1.1⁽⁵⁾ L-C-Z 族的 Lenard 算子对:

$$K = \begin{pmatrix} -\frac{1}{2}\partial^2 + 2\partial q\partial^{-1}q & \partial r - 2\partial q\partial^{-1}r \\ \partial r + r\partial + 2\partial r\partial^{-1}q + 2qr & -\frac{1}{2}\partial^2 - 2\partial r\partial^{-1}r - q\partial - 2r^2 \\ 0 & \frac{1}{2}\partial \end{pmatrix},$$

$$J = 2i \begin{pmatrix} \frac{1}{2}\partial + q & -r \end{pmatrix}$$

$$\partial\partial^{-1} = \partial^{-1}\partial = 1, \text{ 即 } \partial^{-1} = \frac{1}{2} \left(\int_{-\infty}^x \cdot - \int_x^{+\infty} \cdot \right) \text{ or } \partial^{-1} = \frac{1}{2} \left(\int_0^x \cdot - \int_x^T \cdot \right).$$

具有下述两条性质

1) 若 $\{G_j\}$ 为 Lenard 递推叙列, 则

$$KG_j = JG_{j+1}, (j=0, 1, 2, \dots), G_0 = (r, q)^T. \quad (1.2)$$

2) 若 ζ_j 是 (1.1) 式的特征值, 让 $\tilde{A}_j = \begin{pmatrix} \varphi_{2j}^2 + \varphi_{1j}^2 \\ \varphi_{1j} - \varphi_{2j} \end{pmatrix}$, 那么 \tilde{A}_j 满足线性关系式

$$K\tilde{A}_j = \zeta_j \cdot J\tilde{A}_j \quad (1.3)$$

其中, $\varphi_{1j}, \varphi_{2j}$ 为相应于 ζ_j 的特征函数.

命题 1.2⁽⁵⁾ 若 $G_j = (G_j^{(1)}, G_j^{(2)})^T$ 是 Lenard 递推叙列, 让

$$V_j = \begin{pmatrix} -\frac{1}{2}\partial G_j^{(2)} + (G_j^{(1)} + 2\partial^{-1}(qG_j^{(1)} - rG_j^{(2)}))\partial \\ -\frac{1}{2}\partial (G_j^{(1)} - G_j^{(2)}) - qG_j^{(1)} + rG_j^{(2)} \\ -\frac{1}{2}\partial (G_j^{(1)} + G_j^{(2)}) - qG_j^{(1)} + rG_j^{(2)} \\ \frac{1}{2}\partial G_j^{(2)} + (G_j^{(1)} + 2\partial^{-1}(qG_j^{(1)} - rG_j^{(2)}))\partial \end{pmatrix},$$

$j = -1, 0, 1, \dots$

$$\text{则 } (W_m, L) = L_*(X_m), m = 0, 1, 2, \dots \quad (1.4)$$

$$\text{此处, } W_m = \sum_{j=0}^m V_{j-1} L^{m-1}; V_{-1} \triangleq \begin{pmatrix} i\partial & 0 \\ 0 & i\partial \end{pmatrix}, G_{-1} \triangleq \begin{pmatrix} 0 \\ 0 \end{pmatrix}. \quad (1.5)$$

L_* 表示 L 的微分映射.

设 ζ_1, \dots, ζ_N 为 (1.1) 的 N 个互不相同的特征值, 那么

$$L \varphi_j = \zeta_j \varphi_j \quad j=1, \dots, N \quad (1.6)$$

浓缩 (1.6) 为向量形式:

$$\begin{cases} \varphi_1' = -i \wedge \varphi_1 + r \varphi_1 + (q+r) \varphi_2 \\ \varphi_2' = (q-r) \varphi_1 - r \varphi_2 + i \wedge \varphi_2 \end{cases} \quad (1.7)$$

其中, $\varphi_k = (\varphi_{k1}, \dots, \varphi_{kN})^T$ $k=1, 2$; $\wedge = \text{diag}(\zeta_1, \dots, \zeta_N)$; $(\varphi_{1j}, \varphi_{2j})^T$ 是相应于特征值 ζ_j 的特征函数, $j=1, 2, \dots, N$.

令

$$G_0 = \sum_{j=1}^N \bar{A}_j = \begin{pmatrix} \langle \varphi_2, \varphi_2 \rangle + \langle \varphi_1, \varphi_1 \rangle \\ -\langle \varphi_2, \varphi_2 \rangle + \langle \varphi_1, \varphi_1 \rangle \end{pmatrix} \quad (1.8)$$

则 (1.8) 产生 Bargmann 约束 ($\langle \cdot, \cdot \rangle$ 表示 R^N 中之标准内积):

$$q = \langle \varphi_1, \varphi_1 \rangle - \langle \varphi_2, \varphi_2 \rangle, \quad r = \langle \varphi_1, \varphi_1 \rangle + \langle \varphi_2, \varphi_2 \rangle \quad (1.9)$$

或 (1.9) 记为 $f: (\varphi_1, \varphi_2)^T \rightarrow u = (q, r)^T$ (1.10)

在 Bargmann 约束 (1.9) 下, (1.7) 被非线性化为一个三次的 Bargmann 系统

$$(B) \quad \begin{cases} \varphi_1' = -i \wedge \varphi_1 + (\langle \varphi_1, \varphi_1 \rangle + \langle \varphi_2, \varphi_2 \rangle) \varphi_1 + 2 \langle \varphi_1, \varphi_1 \rangle \varphi_2 \\ \varphi_2' = i \wedge \varphi_2 - (\langle \varphi_1, \varphi_1 \rangle + \langle \varphi_2, \varphi_2 \rangle) \varphi_2 - 2 \langle \varphi_2, \varphi_2 \rangle \varphi_1 \end{cases} \quad (1.11)$$

引理 1.3 若 φ_1, φ_2 满足 (1.11), 则 $\frac{\partial}{\partial x} \langle \varphi_1, \varphi_2 \rangle = 0$, 即内积 $\langle \varphi_1, \varphi_2 \rangle$ 是依赖于 φ_1, φ_2 的运动常数, 而与自变量 x 无关.

证明 $\partial \langle \varphi_1, \varphi_2 \rangle = \langle \varphi_1', \varphi_2 \rangle + \langle \varphi_1, \varphi_2' \rangle = 0$.

鉴于引理 1.3, 现引入一个约束关系

$$\langle \varphi_1, \varphi_2 \rangle = 0 \quad (1.12)$$

因而, 我们有

定理 1.4 在 Bargmann 约束 (1.9) 及 (1.12) 下, (1.7) 的非线性化 (1.11) 可表为 Hamilton 结构:

$$(H) \quad \begin{cases} \varphi_1' = -\frac{\partial H}{\partial \varphi_2} \\ \varphi_2' = \frac{\partial H}{\partial \varphi_1} \end{cases} \quad (1.13)$$

其中, $H = \langle i \wedge \varphi_1, \varphi_2 \rangle - \langle \varphi_1 + \varphi_2, \varphi_1 + \varphi_2 \rangle \langle \varphi_1, \varphi_2 \rangle - \langle \varphi_1, \varphi_1 \rangle \langle \varphi_2, \varphi_2 \rangle + \langle \varphi_1, \varphi_2 \rangle^2$ (1.14)

称为 (1.13) 的 Hamilton 函数.

证明 将 \$H\$ 的表达式分别对 \$\varphi_1, \varphi_2\$ 求导, 然后利用 (1.12) 式即知本定理成立.

定理 1.4 告诉我们, Hamilton 系统 (1.13) 在约束 (1.12) 下恰为 Bargmann 系统 (1.11). 但约束关系式 (1.12) 是本质的 (由引理 1.3).

定理 1.5 设 \$(\varphi_1, \varphi_2)^T\$ 是 (1.11) 的一个解, 那么 \$q = \langle \varphi_1, \varphi_1 \rangle - \langle \varphi_2, \varphi_2 \rangle\$, \$r = \langle \varphi_1, \varphi_1 \rangle + \langle \varphi_2, \varphi_2 \rangle\$ 满足一个定态的 L-C-Z 方程:

$$X_N + \alpha_1 X_{N-1} + \dots + \alpha_N X_0 = 0 \tag{1.15}$$

其中 \$\alpha_1, \dots, \alpha_N\$ 是适当选取的、依赖于 \$\varphi_1, \varphi_2\$ 的运动常数.

证明 让算子 \$J^{-1}K\$ 作用 (1.8) 式的两端, 由命题 1.1, 我们有

$$G_1 + \beta_2 G_{-1} = \sum_{j=1}^N \zeta_j \cdot \tilde{A}_1 = \begin{pmatrix} \langle \varphi_1, \wedge \varphi_1 \rangle + \langle \varphi_2, \wedge \varphi_2 \rangle \\ \langle \varphi_1, \wedge \varphi_1 \rangle - \langle \varphi_2, \wedge \varphi_2 \rangle \end{pmatrix} \tag{1.16}$$

\$J^{-1}K\$ 作用 (1.8) 式 \$k\$ 次后, 得到:

$$\begin{aligned} G_k + \beta_2 G_{k-2} + \dots + \beta_k G_0 + \beta_{k+1} G_{-1} &= \sum_{j=1}^N \zeta_j^k \cdot \tilde{A}_1 \\ &= \begin{pmatrix} \langle \varphi_2, \wedge^k \varphi_2 \rangle + \langle \wedge^k \varphi_1, \varphi_1 \rangle \\ \langle \varphi_1, \wedge^k \varphi_1 \rangle - \langle \wedge^k \varphi_2, \varphi_2 \rangle \end{pmatrix} \end{aligned} \tag{1.17}$$

在 (1.17) 中, \$\beta_2, \dots, \beta_{k+1}\$ 皆为依赖于 \$\varphi_1, \varphi_2\$ 的运动常数; \$G_0, \dots, G_k\$ 为 Lenard 递推叙列, 特别指出: \$G_{-1} = (0, 0)^T\$, \$J^{-1}KG_{-1} = G_0\$.

考察多项式: \$P(z) \triangleq (z - \zeta_1) \dots (z - \zeta_N) = \sum_{k=0}^N p_{N-k} z^k\$ \$p_0 = 1, p_1, \dots, p_N\$

由特征值 \$\zeta_1, \dots, \zeta_N\$ 确定.

让算子 \$J \sum_{k=0}^N p_{N-k} \cdot\$ 作用 (1.17) 式的两端, 经整理即得 (注意 \$p_0 = 1\$):

$$X_N + \alpha_1 X_{N-1} + \dots + \alpha_N X_0 = 0 \tag{1.18}$$

其中, \$\alpha_1, \dots, \alpha_N\$ 为由 \$p_1, p_2, \dots, p_N\$ 与 \$\beta_2, \dots, \beta_N\$ 所决定的且依赖于 \$\varphi_1, \varphi_2\$ 的运动常数.

此外, 由上述证明, 我们还得到

引理 1.6 设 \$(\varphi_1, \varphi_2)^T\$ 是 (1.11) 的一个解: \$\{G_j\}\$ 为 Lenard 递推叙列. 那么存在常数 \$c_2, \dots, c_m\$, 使得:

$$A_m = \begin{pmatrix} A_m^{(1)} \\ A_m^{(2)} \end{pmatrix} \triangleq \begin{pmatrix} \langle \varphi_1, \wedge^m \varphi_1 \rangle + \langle \varphi_2, \wedge^m \varphi_2 \rangle \\ \langle \varphi_1, \wedge^m \varphi_1 \rangle - \langle \varphi_2, \wedge^m \varphi_2 \rangle \end{pmatrix} = G_m + \sum_{s=2}^m c_s G_{m-s} \tag{1.19}$$

或
$$A_m = \sum_{s=0}^m c_s G_{m-s}, \quad c_0 = 1, \quad c_1 = 0. \tag{1.20}$$

2 一个有限维对合系

在辛空间 $(R^{2N}, d\varphi, \wedge d\varphi_2)$ 中, 两个函数 F, G 的 Poisson 括号定义为:

$$\begin{aligned} (F, G) &= \sum_{j=1}^N \frac{\partial F}{\partial \varphi_{2j}} \cdot \frac{\partial G}{\partial \varphi_{1j}} - \frac{\partial F}{\partial \varphi_{1j}} \cdot \frac{\partial G}{\partial \varphi_{2j}} \\ &= \left\langle \frac{\partial F}{\partial \varphi_2}, \frac{\partial G}{\partial \varphi_1} \right\rangle - \left\langle \frac{\partial F}{\partial \varphi_1}, \frac{\partial G}{\partial \varphi_2} \right\rangle \end{aligned} \quad (2.1)$$

其中, $\varphi_k = (\varphi_{k1}, \dots, \varphi_{kN})^T$ ($k=1, 2$).

F, G 称为对合, 如果 $(F, G) = 0$.

现在, 我们考虑函数 (m 为非负正数):

$$\begin{aligned} F_m &= -\langle \wedge^{m+1} \varphi_1, \varphi_2 \rangle - i \langle \wedge^m (\varphi_1 + \varphi_2), \varphi_1 + \varphi_2 \rangle \langle \varphi_1, \varphi_2 \rangle \\ &\quad - i \sum_{j=0}^m \left| \begin{array}{cc} \langle \varphi_1, \wedge^j \varphi_1 \rangle & \langle \varphi_1, \wedge^j \varphi_2 \rangle \\ \langle \varphi_1, \wedge^{m-j} \varphi_2 \rangle & \langle \varphi_2, \wedge^{m-j} \varphi_2 \rangle \end{array} \right| \end{aligned} \quad (2.2)$$

特别指出: $-iF_0 = H$. 在 (2.2) 中, $\wedge = \text{diag}(\zeta_1, \dots, \zeta_N)$.

引理 2.1 在约束 (1.12) 下, 内积 $\left\langle \frac{\partial F_m}{\partial \varphi_2}, \frac{\partial F_n}{\partial \varphi_1} \right\rangle$ 关于 m, n 对称, 即

$$\left\langle \frac{\partial F_m}{\partial \varphi_2}, \frac{\partial F_n}{\partial \varphi_1} \right\rangle = \left\langle \frac{\partial F_n}{\partial \varphi_2}, \frac{\partial F_m}{\partial \varphi_1} \right\rangle, \quad \forall m, n \quad (2.3)$$

证明 在约束 (1.12) 式下, 我们有:

$$\begin{aligned} -\frac{\partial F_m}{\partial \varphi_2} &= \wedge^{m+1} \varphi_1 + i \langle \wedge^m (\varphi_1 + \varphi_2), \varphi_1 + \varphi_2 \rangle \varphi_1 \\ &\quad + 2i \sum_{j=0}^m (\langle \varphi_1, \wedge^j \varphi_1 \rangle \wedge^{m-j} \varphi_2 - \langle \varphi_1, \wedge^j \varphi_2 \rangle \wedge^{m-j} \varphi_1) \\ -\frac{\partial F_n}{\partial \varphi_1} &= \wedge^{n+1} \varphi_2 + i \langle \wedge^n (\varphi_1 + \varphi_2), \varphi_1 + \varphi_2 \rangle \varphi_2 \\ &\quad + 2i \sum_{j=0}^n (\langle \varphi_2, \wedge^j \varphi_2 \rangle \wedge^{n-j} \varphi_1 - \langle \varphi_1, \wedge^j \varphi_2 \rangle \wedge^{n-j} \varphi_2) \end{aligned}$$

把上二式右端作 R^N 中的标准内积, 只要仔细计算, 就不难发现该内积是若干项关于 m, n 对称的内积之和, 因而本引理成立.

定理 2.2 在约束 (1.12) 下, 以 (2.2) 定义的函数 F_m 是两两对合的, 即

$$(F_k, F_l) = 0, \quad \forall k, l. \quad (2.4)$$

证明 $(F_k, F_l) = \left\langle \frac{\partial F_k}{\partial \varphi_2}, \frac{\partial F_l}{\partial \varphi_1} \right\rangle - \left\langle \frac{\partial F_l}{\partial \varphi_2}, \frac{\partial F_k}{\partial \varphi_1} \right\rangle = 0$.

定理2.3 在约束(1.12)下,由(2.2)定义的Hamilton系统 $(R^{2N}, d\varphi_1 \wedge d\varphi_2, F_m)$ 在Liouville意义下完全可积.特别,在约束(1.12)下,Hamilton系统 $(R^{2N}, d\varphi_1 \wedge d\varphi_2, H = -iF_0)$ 在Liouville意义下完全可积.

至此,已说明Bargmann系统(1.11)的完全可积性,且其对合系为 $\{F_m\}$ (对合是指在(1.12)下成立)

3 L-C-Z族方程的对合解

考虑 F_m 一流的正则方程(在(1.12)下):

$$(F_m): \begin{cases} \frac{\partial \varphi_1}{\partial t_m} = -\frac{\partial F_m}{\partial \varphi_2} = \wedge^{m+1} \varphi_1 + i \langle \wedge^m (\varphi_1 + \varphi_2), \varphi_1 + \varphi_2 \rangle \varphi_1 + \\ \quad 2i \sum_{j=0}^m (\langle \varphi_1, \wedge^j \varphi_1 \rangle \wedge^{m-j} \varphi_2 - \langle \varphi_1, \wedge^j \varphi_2 \rangle \wedge^{m-j} \varphi_1) \\ \frac{\partial \varphi_2}{\partial t_m} = \frac{\partial F_m}{\partial \varphi_1} = -\wedge^{m+1} \varphi_2 - i \langle \wedge^m (\varphi_1 + \varphi_2), \varphi_1 + \varphi_2 \rangle \varphi_2 - \\ \quad 2i \sum_{j=0}^m (\langle \varphi_2, \wedge^j \varphi_2 \rangle \wedge^{m-j} \varphi_1 - \langle \varphi_1, \wedge^j \varphi_2 \rangle \wedge^{m-j} \varphi_2) \end{cases} \quad (3.1)$$

当 $m=0$ 时,令 $t_0 = x$,那么 $(-iF_0)$ 就是(1.11)

(3.1)简记为

$$(F_m): \frac{\partial}{\partial t_m} \begin{pmatrix} \varphi_1 \\ \varphi_2 \end{pmatrix} = \begin{pmatrix} -\partial F_m / \partial \varphi_2 \\ \partial F_m / \partial \varphi_1 \end{pmatrix} = I \nabla F_m, \quad I = \begin{pmatrix} 0 & -I_N \\ I_N & 0 \end{pmatrix} \quad (3.2)$$

$\nabla F_m = (\partial F_m / \partial \varphi_1, \partial F_m / \partial \varphi_2)^T$, I_N 为 R^N 上的 $N \times N$ 单位矩阵.以 $g_m^{t_m}$ 表示(3.2)初值问题的解算子,那么(3.2)的解可以表为:

$$\begin{pmatrix} \varphi_1(t_m) \\ \varphi_2(t_m) \end{pmatrix} = g_m^{t_m} \begin{pmatrix} \varphi_1(0) \\ \varphi_2(0) \end{pmatrix} \quad (3.3)$$

既然任二 F_k, F_l 在(1.12)下对合, $(F_k, F_l) = 0$,因此我们有:

命题3.1^[6] 1) 任二正则系统 $(F_k), (F_l)$ 相容; 2) Hamilton相流 $g_k^{t_k}, g_l^{t_l}$ 可换.上述两条均在(1.12)下成立.

以 $x = t_0, t_m$ 分别表示系统 $(-iF_0) = (H), (F_m)$ 的流变量,定义:

$$\begin{pmatrix} \varphi_1(x, t_m) \\ \varphi_2(x, t_m) \end{pmatrix} = g_0^x g_m^{t_m} \begin{pmatrix} \varphi_1(0, 0) \\ \varphi_2(0, 0) \end{pmatrix} \quad (3.4)$$

由于流 $g_0^x, g_m^{t_m}$ 在(1.12)下可换,所以 $\varphi_1(x, t_m), \varphi_2(x, t_m)$ 是 (x, t_m) 的二元光滑函数. $(\varphi_1(x, t_m), \varphi_2(x, t_m))^T$ 称为在(1.12)下的相容系统 $(F_0), (F_m)$ 的对合解.

定理3.2 设 $(\varphi_1(x, t_m), \varphi_2(x, t_m))^T$ 是在 (1.12) 下的相容系统 $(H) = (-iF_0), (F_m)$ 的一个对合解, 那么

1) 流方程 $(-iF_0) = (H), (F_m)$ 在 Bargmann 约束 (1.9) 及约束 (1.12) 下, 可以分别被约化为高阶 L-C-Z 方程 Lax 组的空间部分和时间部分 $(u = (q, r)^T)$ 作为它们的位势):

$$L(u) \varphi = \Lambda_0 \varphi \quad (\text{空间部分}) \quad (3.5)$$

$$\frac{\partial \varphi}{\partial t_m} = (W_m + c_1 W_{m-1} + \dots + c_m W_0) \varphi \quad (\text{时间部分}) \quad (3.6)$$

其中, $\Lambda_0 = \begin{pmatrix} \Lambda & 0 \\ 0 & \Lambda \end{pmatrix}_{2N \times 2N}$; $\varphi = (\varphi_1, \varphi_2)^T$, $\varphi_k = (\varphi_{k1}, \dots, \varphi_{kN})^T$, $k = 1$

2; $L(u) = \frac{1}{i} \begin{pmatrix} (r - \partial) I_N & (q + r) I_N \\ (r - q) I_N & (r + \partial) I_N \end{pmatrix}$; $c_j (j = 1, 2, \dots, m)$ 为依赖于 φ_1, φ_2 的运动常数 (与自变量 x 无关); $W_k (k = 0, \dots, m)$ 如 (1.5) 所示.

$$2) u(x, t_m) = \begin{pmatrix} q(x, t_m) \\ r(x, t_m) \end{pmatrix} = \begin{pmatrix} \langle \varphi_1, \varphi_1 \rangle - \langle \varphi_2, \varphi_2 \rangle \\ \langle \varphi_1, \varphi_1 \rangle + \langle \varphi_2, \varphi_2 \rangle \end{pmatrix} \text{ 满足 } m+1 \text{ 阶的}$$

$$\text{L-C-Z 方程: } \frac{\partial u}{\partial t_m} = X_m + c_1 X_{m-1} + \dots + c_m X_0 \quad (3.7)$$

这里, $c_j (j = 1, 2, \dots, m)$ 为不依赖于 x 的常数, 其意义同 (3.6) 中的一样; $X_k (k = 0, \dots, m)$ 为 L-C-Z 向量场^[5].

证明 在 $q(x, t_m) = \langle \varphi_1, \varphi_1 \rangle - \langle \varphi_2, \varphi_2 \rangle$, $r(x, t_m) = \langle \varphi_1, \varphi_1 \rangle + \langle \varphi_2, \varphi_2 \rangle$ 及 $\langle \varphi_1, \varphi_2 \rangle = 0$ 下, 流方程 (H) 为 (3.5) 是显然的.

由引理 1.6, 存在常数 c_2, \dots, c_m , 使得:

$$\begin{aligned} A_m &= \begin{pmatrix} A_m^{(1)} \\ A_m^{(2)} \end{pmatrix} \triangleq \begin{pmatrix} \langle \varphi_1, \Lambda^m \varphi_1 \rangle + \langle \varphi_2, \Lambda^m \varphi_2 \rangle \\ \langle \varphi_1, \Lambda^m \varphi_1 \rangle - \langle \varphi_2, \Lambda^m \varphi_2 \rangle \end{pmatrix} \\ &= \sum_{s=0}^m c_s G_{m-s}, \quad c_0 = 1, c_1 = 0, \forall m \end{aligned} \quad (3.8)$$

规定 $A_{-1} = 0$, $\partial^{-1}(q \cdot 0 - r \cdot 0) = \frac{i}{2}$, 经一系列仔细计算, 我们发现

$$\begin{aligned} & \sum_{j=0}^m -\frac{1}{2} A_{j-1}^{(2)} \Lambda^{m-1} \varphi_1 + (A_{j-1}^{(1)} + 2 \partial^{-1}(q A_{j-1}^{(1)} - r A_{j-1}^{(2)})) \Lambda^{m-1} \varphi_1 \\ & \quad + (-\frac{1}{2} (A_{j-1}^{(1)} + A_{j-1}^{(2)})_x - q A_{j-1}^{(1)} + r A_{j-1}^{(2)}) \Lambda^{m-1} \varphi_2 \\ & = \Lambda^{m+1} \varphi_1 + i \langle \Lambda^m (\varphi_1 + \varphi_2), \varphi_1 + \varphi_2 \rangle \varphi_1 \\ & \quad + 2 i \sum_{j=0}^m (\langle \varphi_1, \Lambda^j \varphi_1 \rangle \Lambda^{m-1} \varphi_2 - \langle \varphi_1, \Lambda^j \varphi_2 \rangle \Lambda^{m-1} \varphi_1) \end{aligned}$$

$$\begin{aligned}
 &= -\frac{\partial F_m}{\partial \varphi_2} \\
 &\quad \sum_{j=0}^m \left\{ -\frac{1}{2}(A_{j-1}^{(1)} - A_{j-1}^{(2)})_x - qA_{j-1}^{(1)} + \right. \\
 &\quad \quad \left. + (A_{j-1}^{(1)} + 2\partial^{-1}(qA_{j-1}^{(1)} - rA_{j-1}^{(2)})) \right\} \wedge^{m-j} \varphi_1 + \frac{1}{2} A_{j-1}^{(2)} \wedge^{m-j} \varphi_2 \\
 &= -\wedge^{m+1} \varphi_2 - i \langle \wedge^m (\varphi_1 + \varphi_2) \rangle \varphi_1 + \varphi_2 - 2i \sum_{j=1}^m (\langle \varphi_2, \wedge^j \varphi_2 \rangle \wedge^{m-j} \varphi_1 - \\
 &\quad - \langle \varphi_1, \wedge^j \varphi_2 \rangle \wedge^{m-j} \varphi_2) \\
 &= \frac{\partial F_m}{\partial \varphi_1} .
 \end{aligned}$$

在上述运算过程中, 用到等式: $\partial \langle \varphi_1, \wedge^{j-1} \varphi_2 \rangle = qA_{j-1}^{(1)} - rA_{j-1}^{(2)}$.

将 (3.8) 代入上二式, 得到

$$\begin{aligned}
 \frac{\partial \varphi_1}{\partial t_m} &= -\frac{\partial F_m}{\partial \varphi_2} \\
 &= \sum_{j=0}^m \sum_{s=0}^j c_s \left\{ -\frac{1}{2} G_{j-1-s}^{(2)} \wedge^{m-j} \varphi_1 + (G_{j-1-s}^{(1)} + 2\partial^{-1}(qG_{j-1-s}^{(1)} - \right. \\
 &\quad \left. - rG_{j-1-s}^{(2)})) \wedge^{m-j} \varphi_{1,x} + \left(-\frac{1}{2}(G_{j-1-s}^{(1)} + G_{j-1-s}^{(2)})_x \right. \right. \\
 &\quad \left. \left. - qG_{j-1-s}^{(1)} + rG_{j-1-s}^{(2)} \right) \wedge^{m-j} \varphi_2 \right\} \\
 &= \sum_{s=0}^m c_s \sum_{k=0}^{m-s} \left\{ -\frac{1}{2} G_{k-1}^{(2)} \wedge^{m-s-k} \varphi_1 + (G_{k-1}^{(1)} + 2\partial^{-1}(qG_{k-1}^{(1)} - \right. \\
 &\quad \left. - rG_{k-1}^{(2)})) \wedge^{m-s-k} \varphi_{1,x} + \left(-\frac{1}{2}(G_{k-1}^{(1)} + G_{k-1}^{(2)})_x - \right. \right. \\
 &\quad \left. \left. - qG_{k-1}^{(1)} + rG_{k-1}^{(2)} \right) \wedge^{m-s-k} \varphi_2 \right\} \quad (3.9)
 \end{aligned}$$

$$\begin{aligned}
 \frac{\partial \varphi_2}{\partial t_m} &= \frac{\partial F_m}{\partial \varphi_1} \\
 &= \sum_{j=0}^m \sum_{s=0}^j c_s \left\{ -\frac{1}{2}(G_{j-1-s}^{(1)} - G_{j-1-s}^{(2)})_x - qG_{j-1-s}^{(1)} + rG_{j-1-s}^{(2)} \right\} \wedge^{m-j} \varphi_1 \\
 &\quad + \frac{1}{2} G_{j-1-s}^{(2)} \wedge^{m-j} \varphi_2 + (G_{j-1-s}^{(1)} + 2\partial^{-1}(qG_{j-1-s}^{(1)} - \\
 &\quad - rG_{j-1-s}^{(2)})) \wedge^{m-j} \varphi_{2,x} \} \\
 &= \sum_{s=0}^m c_s \sum_{k=0}^{m-s} \left\{ \left(-\frac{1}{2}(G_{k-1}^{(1)} - G_{k-1}^{(2)})_x - qG_{k-1}^{(1)} + rG_{k-1}^{(2)} \right) \wedge^{m-s-k} \varphi_1 \right.
 \end{aligned}$$

$$+\frac{1}{2}G_{k-1,x}^{(2)} \wedge^{m-s-k} \varphi_2 + (G_{k-1}^{(1)} + 2\partial^{-1}(qG_{k-1}^{(1)} - rG_{k-1}^{(2)})) \wedge^{m-s-k} \varphi_{2,x} \} \quad (3.10)$$

综合 (3.9) 与 (3.10) 两式, 并结合 (1.5) 式, 立得:

$$\begin{aligned} \frac{\partial \varphi}{\partial t_m} &= \sum_{s=0}^m c_s \sum_{k=0}^{m-s} V_{k-1} \cdot \wedge_0^{m-s-k} \varphi \\ &= \sum_{s=0}^m c_s \sum_{k=0}^{m-s} V_{k-1} \cdot L^{m-s-k} \varphi \\ &= \sum_{s=0}^m c_s W_{m-s} \varphi = (W_m + c_1 W_{m-1} + \dots + c_m W_0) \varphi \end{aligned}$$

其中, $W_{m-s} = \sum_{k=0}^{m-s} V_{k-1} \cdot L^{m-s-k}$; $L=L(u) = \frac{1}{i} \begin{pmatrix} (r-\partial)I_N & (q+r)I_N \\ (r-q)I_N & (r-q)I_N \end{pmatrix}$, V_{k-1}

由下式给出:

$$V_{k-1} = \begin{pmatrix} (-\frac{1}{2}G_{k-1,x}^{(2)} + (G_{k-1}^{(1)} + 2\partial^{-1}(qG_{k-1}^{(1)} - rG_{k-1}^{(2)}))\partial)I_N \\ (-\frac{1}{2}(G_{k-1}^{(1)} - G_{k-1}^{(2)})_x - qG_{k-1}^{(1)} + rG_{k-1}^{(2)})I_N \\ (-\frac{1}{2}(G_{k-1}^{(1)} + G_{k-1}^{(2)})_x - qG_{k-1}^{(1)} + rG_{k-1}^{(2)})I_N \\ (-\frac{1}{2}G_{k-1,x}^{(2)} + (G_{k-1}^{(1)} + 2\partial^{-1}(qG_{k-1}^{(1)} - rG_{k-1}^{(2)}))\partial)I_N \end{pmatrix} \quad K=0, 1, \dots, m-s$$

$G_{k-1} = (G_{k-1}^{(1)}, G_{k-1}^{(2)})^T$ 是 Lenard 递推叙列.

$$\begin{aligned} \frac{\partial q}{\partial t_m} &= 2 \langle \varphi_1, \frac{\partial \varphi_1}{\partial t_m} \rangle - 2 \langle \varphi_2, \frac{\partial \varphi_2}{\partial t_m} \rangle = -2 \langle \varphi_1, \frac{\partial F_m}{\partial \varphi_2} \rangle - 2 \langle \varphi_2, \frac{\partial F_m}{\partial \varphi_1} \rangle \\ &= 2 \langle \varphi_1, \wedge^{m+1} \varphi_1 \rangle + 2i \langle \wedge^m (\varphi_1 + \varphi_2), \varphi_1 + \varphi_2 \rangle (\langle \varphi_1, \varphi_1 \rangle \\ &\quad + \langle \varphi_2, \varphi_2 \rangle) + 2 \langle \varphi_2, \wedge^{m+1} \varphi_2 \rangle \\ &= i(\langle \varphi_1, \wedge^m \varphi_1 \rangle - \langle \varphi_2, \wedge^m \varphi_2 \rangle)_x = iA_{m,x}^{(2)} = 2i \cdot \frac{1}{2} A_{m,x}^{(2)} \quad (3.11) \end{aligned}$$

$$\begin{aligned} \frac{\partial r}{\partial t_m} &= 2 \langle \varphi_1, \frac{\partial \varphi_1}{\partial t_m} \rangle + 2 \langle \varphi_2, \frac{\partial \varphi_2}{\partial t_m} \rangle = -2 \langle \varphi_1, \frac{\partial F_m}{\partial \varphi_2} \rangle + 2 \langle \varphi_2, \frac{\partial F_m}{\partial \varphi_1} \rangle \\ &= 2 \langle \varphi_1, \wedge^{m+1} \varphi_1 \rangle + 2i \langle \wedge^m (\varphi_1 + \varphi_2), \varphi_1 + \varphi_2 \rangle (\langle \varphi_1, \varphi_1 \rangle \\ &\quad - \langle \varphi_2, \varphi_2 \rangle) - 2 \langle \varphi_2, \wedge^{m+1} \varphi_2 \rangle \end{aligned}$$

$$= i(\langle \varphi_1 + \varphi_2, \wedge^m(\varphi_1 + \varphi_2) \rangle)_x = 2i \left(\frac{1}{2} A_m^{(1)} x + q A_m^{(1)} - r A_m^{(2)} \right) \quad (3.12)$$

由 (3.8)、(3.11)、(3.12) 及 L-C-Z 向量场 $X_k = JG_k$, 我们有

$$\begin{aligned} \frac{\partial u}{\partial t_m} &= \frac{\partial}{\partial t_m} \begin{pmatrix} q(x, t_m) \\ r(x, t_m) \end{pmatrix} = 2i \left(\frac{0}{2} \partial + q \quad \frac{1}{2} \partial \quad -r \right) \begin{pmatrix} A_m^{(1)} \\ A_m^{(2)} \end{pmatrix} \\ &= JA_m = J \left(\sum_{s=0}^m c_s G_{m-s} \right) \\ &= X_m + c_1 X_{m-1} + \cdots + c_m X_0 \end{aligned}$$

附注 3.3 L-C-Z 向量场

$$X_0 = i \begin{pmatrix} q_x \\ r_x \end{pmatrix}; X_1 = \begin{pmatrix} -\frac{1}{2} r_{xx} + r_x q + q_x r \\ -\frac{1}{2} q_{xx} - q q_x + 3 r_x \end{pmatrix}$$

当 $m=1$ 时, 按定理 3.2, 我们可求得二阶非线性 L-C-Z 方程: $q_t = -\frac{1}{2} r_{xx} +$

$r_x q + q_x r$, $r_t = -\frac{1}{2} q_{xx} - q q_x + 3 r_x$ 的对合解: $q = \langle \varphi_1, \varphi_1 \rangle - \langle \varphi_2, \varphi_2 \rangle$,
 $r = \langle \varphi_1, \varphi_1 \rangle + \langle \varphi_2, \varphi_2 \rangle$, 其中, $(\varphi_1, \varphi_2)^T$ 是在 (1.12) 下的相容系统 (H)、
 (F_1) 的对合解.

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Completely Integrable Bargmann System Associated With the L-C-Z Soliton Hierarchy

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ABSTRACT In this paper, a new involutive system is obtained and under the Bargmann constraint: $q = \langle \varphi_1, \varphi_1 \rangle - \langle \varphi_2, \varphi_2 \rangle$, $r = \langle \varphi_1, \varphi_1 \rangle + \langle \varphi_2, \varphi_2 \rangle$, the L-C-Z spectral problem (2.1) is nonlinearized as a new completely integrable Hamiltonian system in Liouville's sense. Finally, we present the involutive solutions of L-C-Z hierarchy of soliton equations.

KEY WORDS L-C-Z hierarchy, Bargmann constraint, Involutive system, Involutive solution, Hamiltonian system, Complete integrability in Liouville's sense.